Disorder and critical phenomena in pinning models

Giambattista Giacomin

Université Paris Diderot and Laboratoire Probabilités et Modèles Aléatoires (LPMA)

January 25, 2012

This talk will be about disordered pinning models

- definitions: pure and disordered static models
- role of IID disorder: Harris criterion
- disorder smoothing,...

This talk will be about disordered pinning models

- definitions: pure and disordered static models
- role of IID disorder: Harris criterion
- disorder smoothing,...

- correlated disorder [Poisat]
- dynamical pinning [Caputo, Lacoin, Martinelli, Simenhaus, Toninelli]

This talk will be about disordered pinning models

- definitions: pure and disordered static models
- role of IID disorder: Harris criterion
- disorder smoothing,...

- correlated disorder [Poisat]
- dynamical pinning [Caputo, Lacoin, Martinelli, Simenhaus, Toninelli]
- new LD approach [Bolthausen, Cheliotis, den Hollander, Opoku]

This talk will be about disordered pinning models

- definitions: pure and disordered static models
- role of IID disorder: Harris criterion
- disorder smoothing,...

- correlated disorder [Poisat]
- dynamical pinning [Caputo, Lacoin, Martinelli, Simenhaus, Toninelli]
- new LD approach [Bolthausen, Cheliotis, den Hollander, Opoku]
- hierarchical models [...]

This talk will be about disordered pinning models

- definitions: pure and disordered static models
- role of IID disorder: Harris criterion
- disorder smoothing,...

- correlated disorder [Poisat]
- dynamical pinning [Caputo, Lacoin, Martinelli, Simenhaus, Toninelli]
- new LD approach [Bolthausen, Cheliotis, den Hollander, Opoku]
- hierarchical models [...]

"The Example": (1 + d)-directed walk models

Symmetric Random Walk $\{S_n\}_n$ with increments in $\{-1, 0, +1\}$ (d = 1)



Model ($\beta \ge 0$, $h \in \mathbb{R}$):

$$\frac{\mathrm{d}\mathbf{P}_{N,\omega}}{\mathrm{d}\mathbf{P}}(S) = \frac{1}{Z_{N,\omega}} \exp\left(\sum_{n=1}^{N} \left(\beta\omega_n + h\right) \delta_n\right)$$

with $\delta_n = \mathbf{1}_{S_n=0}$. The disorder ω is a IID sequence $\mathcal{N}(0,1)$ of law \mathbb{P} .

"The Example": (1 + d)-directed walk models

Symmetric Random Walk $\{S_n\}_n$ with increments in $\{-1, 0, +1\}$ (d = 1)



Model ($\beta \geq 0$, $h \in \mathbb{R}$):

$$\frac{\mathrm{d}\mathbf{P}_{N,\omega}}{\mathrm{d}\mathbf{P}}(S) = \frac{1}{Z_{N,\omega}} \exp\left(\sum_{n=1}^{N} \left(\beta\omega_n + h\right) \delta_n\right)$$

with $\delta_n = \mathbf{1}_{S_n=0}$. The disorder ω is a IID sequence $\mathcal{N}(0,1)$ of law \mathbb{P} .

"The Example" again: *d*-dim. heteropolymer

A polymer chain made up of *charged* monomers, interacting with a potential near a point in space



Many other physical, biological,... systems:

- Two dimensional interfaces near a (rough) wall
- Of course: DNA denaturation (Poland-Scheraga)
- Flux-lines in super-conductors

"The Example" again: *d*-dim. heteropolymer

A polymer chain made up of *charged* monomers, interacting with a potential near a point in space



Many other physical, biological,... systems:

- Two dimensional interfaces near a (rough) wall
- Of course: DNA denaturation (Poland-Scheraga)
- Flux-lines in super-conductors

"The Example" again: *d*-dim. heteropolymer

A polymer chain made up of *charged* monomers, interacting with a potential near a point in space



Many other physical, biological,... systems:

- Two dimensional interfaces near a (rough) wall
- Of course: DNA denaturation (Poland-Scheraga)
- Flux-lines in super-conductors
- many more, but also: exactly solvable if $\beta = 0$

'Rethinking 'The Example"

Symmetric Random Walk $\{S_n\}_n$ with increments in $\{-1, 0, +1\}$



Model ($\beta \geq 0$, $h \in \mathbb{R}$):

$$\frac{\mathrm{d}\mathbf{P}_{N,\omega}}{\mathrm{d}\mathbf{P}}(S) = \frac{1}{Z_{N,\omega}} \exp\left(\sum_{n=1}^{N} \left(\beta\omega_n + h\right) \delta_n\right)$$

with $\delta_n = \mathbf{1}_{S_n=0}$.

'Rethinking 'The Example"

Symmetric Random Walk $\{S_n\}_n$ with increments in $\{-1, 0, +1\}$



Model ($\beta \geq 0$, $h \in \mathbb{R}$):

$$\frac{\mathrm{d}\mathbf{P}_{N,\omega}}{\mathrm{d}\mathbf{P}}(\tau) = \frac{1}{Z_{N,\omega}} \exp\left(\sum_{n=1}^{N} \left(\beta\omega_n + h\right) \delta_n\right)$$

with $\delta_n = \mathbf{1}_{n \in \tau}$.

Basic building block: a Discrete Renewal Process

Basic building block: a Discrete Renewal Process

• $\tau = \{\tau_0, \tau_1, \tau_2, ...\}$ discrete renewal sequence (that is, $\tau_0 = 0$ and $\{\tau_j - \tau_{j-1}\}_{j \in \mathbb{N}}$ is IID), of law **P**, s. t.

$$K(n) = \mathbf{P}(\tau_1 = n) \sim C_K / n^{1+\alpha}, \ (C_K > 0),$$

or even $\mathbf{P}(au_1 = n) \sim L(n)/n^{1+lpha}$, $L(\cdot)$ a slowly varying function, and

$$\sum_{n\in\mathbb{N}}K(n)\leq 1.$$

Basic building block: a Discrete Renewal Process

② If $\sum_{n} K(n) < 1$ ⇒ renewal on $\mathbb{N} \cup \{\infty\}$, with $K(\infty) = 1 - \sum_{n} K(n)$ (terminating renewal), otherwise the renewal is persistent.

Obs.:
$$\alpha = 1/2$$
 for both $d = 1$ and 3, but $\sum_n K(n) < 1$ if $d = 3$

Basic building block: a Discrete Renewal Process

• $\tau = \{\tau_0, \tau_1, \tau_2, ...\}$ discrete renewal sequence (that is, $\tau_0 = 0$ and $\{\tau_j - \tau_{j-1}\}_{j \in \mathbb{N}}$ is IID), of law **P**, s. t.

$$K(n) = \mathbf{P}(\tau_1 = n) \sim C_K / n^{1+\alpha}, \ (C_K > 0),$$

or even $\mathbf{P}(au_1 = n) \sim L(n)/n^{1+lpha}$, $L(\cdot)$ a slowly varying function, and

$$\sum_{n\in\mathbb{N}}K(n)\leq 1.$$

② If $\sum_{n} K(n) < 1$ ⇒ renewal on $\mathbb{N} \cup \{\infty\}$, with $K(\infty) = 1 - \sum_{n} K(n)$ (terminating renewal), otherwise the renewal is persistent.



The Poland-Scheraga model



- The two thick lines are the DNA strands. They may be paired, gaining thus energetic contributions that depend on whether the base pair is A-T or G-C.
- There are then sections of unpaired bases (the *loops*) to which an entropy is associated: loops correspond to inter-arrival of length n ≥ 2.

The Poland-Scheraga model



- The two thick lines are the DNA strands. They may be paired, gaining thus energetic contributions that depend on whether the base pair is A-T or G-C.
- There are then sections of unpaired bases (the *loops*) to which an entropy is associated: loops correspond to inter-arrival of length n ≥ 2.
- Entropy: $C\mu^n/n^{1+\alpha}$, $\alpha \approx 1.15$ and μ drops out when $N \to \infty$.

The Poland-Scheraga model



- The two thick lines are the DNA strands. They may be paired, gaining thus energetic contributions that depend on whether the base pair is A-T or G-C.
- There are then sections of unpaired bases (the *loops*) to which an entropy is associated: loops correspond to inter-arrival of length n ≥ 2.
- Entropy: $C\mu^n/n^{1+\alpha}$, $\alpha \approx 1.15$ and μ drops out when $N \to \infty$.

Annealed, pure and homogeneous models

Observe that

$$\mathbb{E}Z^{c}_{N,\omega,\beta,h} = \mathbf{E}\left[\exp\left(\sum_{n=1}^{N}((\beta^{2}/2) + h)\delta_{n}\right)\delta_{N}\right]$$

which is the partition function of a homogeneous model.

Annealed, pure and homogeneous models

Observe that

$$\mathbb{E}Z^{c}_{N,\omega,\beta,h} = \mathbf{E}\left[\exp\left(\sum_{n=1}^{N}((\beta^{2}/2) + h)\delta_{n}\right)\delta_{N}\right]$$

which is the partition function of a homogeneous model. So:

• The annealed (or *pure*) model is just a homogeneous model with pinning potential $h + \beta^2/2$;

Annealed, pure and homogeneous models

Observe that

$$\mathbb{E}Z^{c}_{N,\omega,\beta,h} = \mathbf{E}\left[\exp\left(\sum_{n=1}^{N}((\beta^{2}/2) + h)\delta_{n}\right)\delta_{N}\right]$$

which is the partition function of a homogeneous model.

So:

- The annealed (or *pure*) model is just a homogeneous model with pinning potential $h + \beta^2/2$;
- Homogeneous pinning models are *exactly solvable* exhibiting a surprisingly wide spectrum of behaviors (when α varies).

Theorem (Existence of the free energy and self-averaging). The limit

$$\lim_{N\to\infty}\frac{1}{N}\log Z_{N,\omega}=:\mathrm{F}(\beta,h)$$

exists $\mathbb{P}(d\omega)$ -a.s. and in $L^1(\mathbb{P})$, $F(\beta, h)$ is non-random and $F(\beta, h) \ge 0$. Moreover $F(\cdot, \cdot)$ is convex, $F(\beta, \cdot)$ is non-decreasing and $F(0, h) \le F(\beta, h) \le F(0, h + \beta^2/2) =: F^a(\beta, h)$.

Proof of $F(\beta, h) \leq F(0, h + \beta^2/2)$: Jensen's inequality

 $\mathbb{E}\log Z_{N,\omega} \leq \log \mathbb{E} Z_{N,\omega}.$

Theorem (Existence of the free energy and self-averaging). The limit

$$\lim_{N\to\infty}\frac{1}{N}\log Z_{N,\omega}=:\mathrm{F}(\beta,h)$$

exists $\mathbb{P}(d\omega)$ -a.s. and in $L^1(\mathbb{P})$, $F(\beta, h)$ is non-random and $F(\beta, h) \ge 0$. Moreover $F(\cdot, \cdot)$ is convex, $F(\beta, \cdot)$ is non-decreasing and $F(0, h) \le F(\beta, h) \le F(0, h + \beta^2/2) =: F^a(\beta, h)$.

Proof of $F(\beta, h) \leq F(0, h + \beta^2/2)$: Jensen's inequality

 $\mathbb{E} \log Z_{N,\omega} \, \leq \, \log \mathbb{E} Z_{N,\omega}.$

Two fundamental facts

• Dichotomy: $F(\beta, h) > 0$ (localization) versus $F(\beta, h) = 0$ (delocalization)

Theorem (Existence of the free energy and self-averaging). The limit

$$\lim_{N\to\infty}\frac{1}{N}\log Z_{N,\omega}=:\mathrm{F}(\beta,h)$$

exists $\mathbb{P}(d\omega)$ -a.s. and in $L^1(\mathbb{P})$, $F(\beta, h)$ is non-random and $F(\beta, h) \ge 0$. Moreover $F(\cdot, \cdot)$ is convex, $F(\beta, \cdot)$ is non-decreasing and $F(0, h) \le F(\beta, h) \le F(0, h + \beta^2/2) =: F^a(\beta, h)$.

Proof of $F(\beta, h) \leq F(0, h + \beta^2/2)$: Jensen's inequality

 $\mathbb{E} \log Z_{N,\omega} \, \leq \, \log \mathbb{E} Z_{N,\omega}.$

Two fundamental facts

Dichotomy:
 F(β, h) > 0 (localization) versus F(β, h) = 0 (delocalization)

• Key ingredient: super-additivity of $\mathbb{E} \log Z_{N,\omega}^c$, in particular $F(\beta, h) > 0 \iff$ there exists N such that $\mathbb{E} \log Z_{N,\omega}^c > 0$

Theorem (Existence of the free energy and self-averaging). The limit

$$\lim_{N\to\infty}\frac{1}{N}\log Z_{N,\omega}=:\mathrm{F}(\beta,h)$$

exists $\mathbb{P}(d\omega)$ -a.s. and in $L^1(\mathbb{P})$, $F(\beta, h)$ is non-random and $F(\beta, h) \ge 0$. Moreover $F(\cdot, \cdot)$ is convex, $F(\beta, \cdot)$ is non-decreasing and $F(0, h) \le F(\beta, h) \le F(0, h + \beta^2/2) =: F^a(\beta, h)$.

Proof of $F(\beta, h) \leq F(0, h + \beta^2/2)$: Jensen's inequality

 $\mathbb{E} \log Z_{N,\omega} \, \leq \, \log \mathbb{E} Z_{N,\omega}.$

Two fundamental facts

- Dichotomy:
 F(β, h) > 0 (localization) versus F(β, h) = 0 (delocalization)
- Key ingredient: super-additivity of $\mathbb{E} \log Z_{N,\omega}^c$, in particular $F(\beta, h) > 0 \iff$ there exists N such that $\mathbb{E} \log Z_{N,\omega}^c > 0$

Localization and Delocalization

Proof of $F(\beta, h) \ge 0$: $F(\beta, h) = \lim_{N} \frac{1}{N} \mathbb{E} \log \mathbf{E} \left[\exp \left(\sum_{n=1}^{N} (\beta \omega_n + h) \delta_n \right) \right]$ $\ge \liminf_{N} \frac{1}{N} \mathbb{E} \log \mathbf{E} \left[\exp \left(\sum_{n=1}^{N} (\beta \omega_n + h) \delta_n \right); \tau_1 > N \right]$ $= \lim_{N} \frac{1}{N} \log \mathbf{P} (\tau_1 > N) = 0.$

Localization and Delocalization

Proof of $F(\beta, h) \ge 0$:

$$F(\beta, h) = \lim_{N} \frac{1}{N} \mathbb{E} \log \mathbf{E} \left[\exp \left(\sum_{n=1}^{N} (\beta \omega_n + h) \delta_n \right) \right]$$

$$\geq \liminf_{N} \frac{1}{N} \mathbb{E} \log \mathbf{E} \left[\exp \left(\sum_{n=1}^{N} (\beta \omega_n + h) \delta_n \right); \tau_1 > N \right]$$

$$= \lim_{N} \frac{1}{N} \log \mathbf{P} (\tau_1 > N) = 0.$$

$$F(\beta, h) \qquad Delocalized \qquad Localized \qquad h$$

G.G. (Paris Diderot and LPMA)

Localization and Delocalization

Proof of $F(\beta, h) \ge 0$:

$$F(\beta, h) = \lim_{N} \frac{1}{N} \mathbb{E} \log \mathbf{E} \left[\exp \left(\sum_{n=1}^{N} (\beta \omega_n + h) \delta_n \right) \right]$$

$$\geq \liminf_{N} \frac{1}{N} \mathbb{E} \log \mathbf{E} \left[\exp \left(\sum_{n=1}^{N} (\beta \omega_n + h) \delta_n \right); \tau_1 > N \right]$$

$$= \lim_{N} \frac{1}{N} \log \mathbf{P} (\tau_1 > N) = 0.$$

G.G. (Paris Diderot and LPMA)

The phase diagram

Critical point/curve: $h_c(\beta) := \max\{h : F(\beta, h) = 0\}$



The phase diagram

Critical point/curve: $h_c(\beta) := \max\{h : F(\beta, h) = 0\}$



Plenty of questions, but above all:

- Can one compute or estimate $h_c(\beta)$?
- Critical behavior? $F(\beta, h) \overset{h \searrow h_c(\beta)}{\sim} const.(h h_c(\beta))^{\nu_q}$

The pure model is solvable and it displays any kind of critical behavior!

The pure model is solvable and it displays any kind of critical behavior! $h_c(0) = -\log \sum_n K(n) \geq 0$

The pure model is solvable and it displays any kind of critical behavior! $h_c(0) = -\log \sum_n K(n) (\ge 0)$ and if $h > h_c(0)$ the free energy is given by

$$\sum_{n} \mathcal{K}(n) \exp(-F(0,h)n) = \exp(-h),$$

which directly yields

$$F(0,h) \stackrel{h\searrow h_c(0)}{\sim} const.(h-h_c(0))^{\nu_a}$$

with

$$u_a = egin{cases} 1/lpha & ext{ for } lpha \in (0,1) \ 1 & ext{ for } lpha > 1 \end{cases}$$

The pure model is solvable and it displays any kind of critical behavior! $h_c(0) = -\log \sum_n K(n) (\geq 0)$ and if $h > h_c(0)$ the free energy is given by

$$\sum_{n} K(n) \exp(-F(0,h)n) = \exp(-h),$$

which directly yields

$$F(0,h) \stackrel{h\searrow h_c(0)}{\sim} const.(h-h_c(0))^{\nu_a}$$

with

$$u_{a} = egin{cases} 1/lpha & ext{ for } lpha \in (0,1) \ 1 & ext{ for } lpha > 1 \end{cases}$$

...M. Fisher '84. But: Erdos, Pollard, Feller, Garsia, Lamperti... (40's...)

General principles to deal with disorder(?)

Recall the main questions:

- Compute or estimate $h_c(\beta)$
- Critical behavior? $F(\beta, h) \overset{h \searrow h_c(\beta)}{\sim} const.(h h_c(\beta))^{\nu_q}$

General principles to deal with disorder(?)

Recall the main questions:

• Compute or estimate $h_c(\beta)$

• Critical behavior? F
$$(eta,h) \stackrel{h\searrow h_c(eta)}{\sim} const.(h-h_c(eta))^{
u_q}$$

Harris Criterion (A. B. Harris 1974)

Knowing the critical behavior of the pure system one can decide whether (at small disorder) the critical behavior of pure and disordered systems coincide (the disorder is irrelevant) or differ (the disorder is relevant).

General principles to deal with disorder(?)

Recall the main questions:

• Compute or estimate $h_c(\beta)$

• Critical behavior? F
$$(eta,h) \stackrel{h\searrow h_c(eta)}{\sim} const.(h-h_c(eta))^{
u_q}$$

Harris Criterion (A. B. Harris 1974)

Knowing the critical behavior of the pure system one can decide whether (at small disorder) the critical behavior of pure and disordered systems coincide (the disorder is irrelevant) or differ (the disorder is relevant).

HC for pinning models [Forgacs et al. (1986), Derrida et al. (1992)]:

•
$$h_c(\beta) = h_c^a(\beta)$$
 and $\nu_q = \nu_a$ for β small if $\alpha < 1/2$

• $h_c(\beta) \neq h_c^a(\beta)$ and (probably) $\nu_q \neq \nu_a$ for $\beta > 0$ and $\alpha > 1/2$.

Harris criterion for pinning models: rigorous results

• If $0 \le \alpha < 1/2$ disorder is irrelevant if β is not too large: there exists $\beta_0 \in (0,\infty]$ such that for $\beta < \beta_0$ we have

$$h_c(\beta) = h_c^a(\beta)$$

and $\nu_q = \nu_a$, that is

$$\lim_{h \searrow h_c(\beta)} \frac{\log F(\beta, h)}{\log(h - h_c(\beta))} = 1/\alpha$$

[Alexander 08], [Toninelli 08], [GT09], [Lacoin 10]

Harris criterion for pinning models: rigorous results

• If $0 \le \alpha < 1/2$ disorder is irrelevant if β is not too large: there exists $\beta_0 \in (0,\infty]$ such that for $\beta < \beta_0$ we have

$$h_c(\beta) = h_c^a(\beta)$$

and $\nu_q = \nu_a$, that is

$$\lim_{h \searrow h_c(\beta)} \frac{\log F(\beta, h)}{\log(h - h_c(\beta))} = 1/\alpha$$

[Alexander 08], [Toninelli 08], [GT09],[Lacoin 10] • If $\alpha > 1/2$ disorder is relevant for every $\beta > 0$: $h_c(\beta) > h_c^a(\beta)$ and

$$u_q \ge 2 > \nu_a = 1/lpha \quad \text{smoothing!}$$

Moreover

$$h_c(\beta) - h_c^{\mathfrak{d}}(\beta) pprox egin{cases} eta^{2lpha/(2lpha-1)} & ext{if } lpha \in (1/2,1) \ eta^2 & ext{if } lpha > 1 \end{cases}$$

[GT 06], [Derrida GLT 09], [A Zygouras 10]

 Marginal irrelevance of weak disorder: Forgacs, Luck, Nieuwenhuizen, Orland (1986, 1 + 1-dim. wetting)

• Marginal relevance of disorder:

Derrida, Hakim, Vannimenus (1992, 1+1-dim. wetting)

- Marginal irrelevance of weak disorder: Forgacs, Luck, Nieuwenhuizen, Orland (1986, 1 + 1-dim. wetting) Grosberg, Shakhnovich (1986, pinning of a heteropolymer in d = 3) Gangardt, Nechaev (2008, 1 + 1-dim. wetting)
- Marginal relevance of disorder:

Derrida, Hakim, Vannimenus (1992, 1 + 1-dim. wetting) Bhattacharjee, Mukherji (1993, 1 + 1-dim. wetting) Tang, Chaté (2001, 1 + 1-dim. pinning) Stepanow, Chudnovskiy (2002, 1 + 1-dim. pinning)

• Marginal irrelevance of weak disorder:

Forgacs, Luck, Nieuwenhuizen, Orland (1986, 1 + 1-dim. wetting) Grosberg, Shakhnovich (1986, pinning of a heteropolymer in d = 3) Gangardt, Nechaev (2008, 1 + 1-dim. wetting)

• Marginal relevance of disorder:

Derrida, Hakim, Vannimenus (1992, 1 + 1-dim. wetting) Bhattacharjee, Mukherji (1993, 1 + 1-dim. wetting) Tang, Chaté (2001, 1 + 1-dim. pinning) Stepanow, Chudnovskiy (2002, 1 + 1-dim. pinning)

Claim in case of disorder relevance:

critical point shift of $pprox \exp(-c/eta^2)$ (more vague about $u_q)$

• Marginal irrelevance of weak disorder:

Forgacs, Luck, Nieuwenhuizen, Orland (1986, 1 + 1-dim. wetting) Grosberg, Shakhnovich (1986, pinning of a heteropolymer in d = 3) Gangardt, Nechaev (2008, 1 + 1-dim. wetting)

• Marginal relevance of disorder:

Derrida, Hakim, Vannimenus (1992, 1 + 1-dim. wetting) Bhattacharjee, Mukherji (1993, 1 + 1-dim. wetting) Tang, Chaté (2001, 1 + 1-dim. pinning) Stepanow, Chudnovskiy (2002, 1 + 1-dim. pinning)

Claim in case of disorder relevance: critical point shift of $pprox \exp(-c/\beta^2)$ (more vague about u_q)

Rigorous: $\nu_q \ge 2$ (poor...)

$$c_{arepsilon}\exp(-c/eta^{2+arepsilon})\leq h_{c}(eta)-h_{c}^{a}(eta)\leq c\exp(-c/eta^{2})$$

for $\beta \leq \beta_{\rm 0}$ [A08,T08,GLT10,GLT11]

Smoothing inequality

Theorem [GT06, CMP and PRL]

Under assumptions on the disorder, for every $\beta>0$ there exists ${\cal C}_\beta$ such that for every h

$$\operatorname{F}(\beta,h) \leq \alpha C_{\beta} \left(h - h_{c}(\beta)\right)^{2}.$$

Possibly more transparent when written as

$$\mathsf{0} \leq \mathrm{F}(eta, h) - \mathrm{F}(eta, h_{m{c}}(eta)) \leq lpha \mathcal{C}_{eta} \left(h - h_{m{c}}(eta)
ight)^2$$

(the result is non trivial only for $h < h_c(\lambda)$). Rephrasing: $F(\beta, \cdot)$ is $C^{1,1}$ at $h_c(\beta) \implies$ the transition is at least of second order (almost third...)

Smoothing inequality

Theorem [GT06, CMP and PRL]

Under assumptions on the disorder, for every $\beta > 0$ there exists C_{β} such that for every h

$$\operatorname{F}(\beta,h) \leq \alpha C_{\beta} \left(h - h_{c}(\beta)\right)^{2}.$$

Possibly more transparent when written as

$$\mathsf{0} \leq \operatorname{F}(eta, h) - \operatorname{F}(eta, h_c(eta)) \leq lpha \mathcal{C}_eta \left(h - h_c(eta)
ight)^2$$

(the result is non trivial only for $h < h_c(\lambda)$). Rephrasing: $F(\beta, \cdot)$ is $C^{1,1}$ at $h_c(\beta) \Longrightarrow$ the transition is at least of second order (almost third...)

What assumptions?

The disorder is IID and the law of ω_1 either has a strictly positive density (with a finite entropy condition wrt Gaussian) or it has compact support. Generalizes to non IID: e.g. stationary Gaussian with summable covariance. [Berger] G.G. (Paris Diderot and LPMA) Paris, IHP 16 / 19



Consider blocks of length ℓ (large, but finite) and choose N to guaranty that with large probability there is at least a (good!) block in which $\log Z_{\ell, \theta^{i\ell}\omega, \beta, h_c(\beta)}^c$ is larger than $\ell \frac{1}{2} F(\beta, h_c(\beta) + \Delta), \Delta > 0$.



Consider blocks of length ℓ (large, but finite) and choose N to guaranty that with large probability there is at least a (good!) block in which $\log Z^{c}_{\ell,\theta^{i\ell}\omega,\beta,h_{c}(\beta)}$ is larger than $\ell \frac{1}{2} F(\beta, h_{c}(\beta) + \Delta), \Delta > 0$. Fact:

$$\mathbb{P}\left(\log Z^{c}_{\ell,\omega,\beta,h_{c}(\beta)} > \ell \frac{1}{2} \mathrm{F}(\beta,h_{c}(\beta)+\Delta)\right) \underset{\sim}{>} \exp\left(-\frac{1}{2}\ell \frac{\Delta^{2}}{\beta^{2}}\right)$$

because to have $\log Z_{\ell,\omega,\beta,h_c(\beta)}^c > \ell \frac{1}{2} F(\beta,h_c(\beta) + \Delta)$ it suffices that the environment looks like $(\omega_1 + \Delta/\beta, \dots, \omega_\ell + \Delta/\beta)$.



Consider blocks of length ℓ (large, but finite) and choose N to guaranty that with large probability there is at least a (good!) block in which $\log Z^{c}_{\ell,\theta^{i\ell}\omega,\beta,h_{c}(\beta)}$ is larger than $\ell \frac{1}{2} F(\beta,h_{c}(\beta)+\Delta), \Delta > 0$. Fact:

$$\mathbb{P}\left(\log Z^{c}_{\ell,\omega,\beta,h_{c}(\beta)} > \ell \frac{1}{2} F(\beta,h_{c}(\beta)+\Delta)\right) \gtrsim \exp\left(-\frac{1}{2} \ell \frac{\Delta^{2}}{\beta^{2}}\right)$$

because to have $\log Z_{\ell,\omega,\beta,h_c(\beta)}^c > \ell \frac{1}{2} F(\beta,h_c(\beta) + \Delta)$ it suffices that the environment looks like $(\omega_1 + \Delta/\beta, \dots, \omega_\ell + \Delta/\beta)$. This requires $N = O\left(\ell \exp\left(\frac{1}{2}\ell\frac{\Delta^2}{\beta^2}\right)\right)$.



Consider blocks of length ℓ (large, but finite) and choose N to guaranty that with large probability there is at least a (good!) block in which $\log Z^{c}_{\ell,\theta^{i\ell}\omega,\beta,h_{c}(\beta)}$ is larger than $\ell \frac{1}{2} F(\beta,h_{c}(\beta)+\Delta), \Delta > 0$. Fact:

$$\mathbb{P}\left(\log Z^{\boldsymbol{c}}_{\ell,\omega,\beta,\boldsymbol{h_c}(\beta)} > \ell \frac{1}{2} \mathrm{F}(\beta,\boldsymbol{h_c}(\beta) + \Delta)\right) \gtrsim \exp\left(-\frac{1}{2}\ell \frac{\Delta^2}{\beta^2}\right)$$

because to have $\log Z_{\ell,\omega,\beta,h_c(\beta)}^c > \ell \frac{1}{2} F(\beta, h_c(\beta) + \Delta)$ it suffices that the environment looks like $(\omega_1 + \Delta/\beta, \dots, \omega_\ell + \Delta/\beta)$. This requires $N = O\left(\ell \exp\left(\frac{1}{2}\ell\frac{\Delta^2}{\beta^2}\right)\right)$. Then make a lower bound on $Z_{N,\omega,\beta,h_c(\beta)}^c$ by considering only the τ

trajectories visiting the only the first good block





By super-additivity:

$$\frac{1}{N}\mathbb{E}\log Z^{c}_{N,\omega,\beta,h_{c}(\beta)} \leq F(\beta,h_{c}(\beta)) = 0$$



By super-additivity:

$$\frac{1}{N} \mathbb{E} \log Z_{N,\omega,\beta,h_{c}(\beta)}^{c} \leq F(\beta,h_{c}(\beta)) = 0$$

But (recall $N = O\left(\ell \exp\left(\frac{1}{2}\ell\frac{\Delta^{2}}{\beta^{2}}\right)\right)$
 $0 \geq \mathbb{E} \log Z_{N,\omega,\beta,h_{c}(\beta)}^{c} \geq \ell \frac{1}{2}F(\beta,h_{c}(\beta)+\Delta) - C\log N$
 $\geq \ell\left(\frac{1}{2}F(\beta,h_{c}(\beta)+\Delta) - \frac{C}{2}\frac{\Delta^{2}}{\beta^{2}}\right) + O(\log \ell)$

and the non-positivity of the blue term is the smoothing inequality.

Other approaches to smoothing:

- The inequality ν ≥ 2/d in [Chayes×2, Fisher, Spencer 86] is about correlation functions and it is valid under *complex assumptions* (conditional result), verified for the Ising model and for the *quenched averaged* correlation length. It is very unclear what this approach yields for pinning models, and above all for α > 1.
- [Aizenman, Wehr 91] uses a substantially different mechanism: typical fluctuations (CLT: [AW]) versus atypical deviations (Large Deviations: [GT])

Other approaches to smoothing:

- The inequality ν ≥ 2/d in [Chayes×2, Fisher, Spencer 86] is about correlation functions and it is valid under *complex assumptions* (conditional result), verified for the Ising model and for the *quenched averaged* correlation length. It is very unclear what this approach yields for pinning models, and above all for α > 1.
- [Aizenman, Wehr 91] uses a substantially different mechanism: typical fluctuations (CLT: [AW]) versus atypical deviations (Large Deviations: [GT])

Other approaches to smoothing:

- The inequality ν ≥ 2/d in [Chayes×2, Fisher, Spencer 86] is about correlation functions and it is valid under *complex assumptions* (conditional result), verified for the Ising model and for the *quenched averaged* correlation length. It is very unclear what this approach yields for pinning models, and above all for α > 1.
- [Aizenman, Wehr 91] uses a substantially different mechanism: typical fluctuations (CLT: [AW]) versus atypical deviations (Large Deviations: [GT])

Important issue

For $\alpha \ge 1/2$, is $\nu_q > 2?$

Other approaches to smoothing:

- The inequality ν ≥ 2/d in [Chayes×2, Fisher, Spencer 86] is about correlation functions and it is valid under *complex assumptions* (conditional result), verified for the Ising model and for the *quenched averaged* correlation length. It is very unclear what this approach yields for pinning models, and above all for α > 1.
- [Aizenman, Wehr 91] uses a substantially different mechanism: typical fluctuations (CLT: [AW]) versus atypical deviations (Large Deviations: [GT])

Important issue

For $\alpha \ge 1/2$, is $\nu_q > 2$? This is certainly expected, and observed numerically in [Coluzzi, Yeramian 07] for $\alpha = 1.15$, along with the role of atypical deviations.