

Extremes and Order Statistics of $1d$ Branching Brownian Motion

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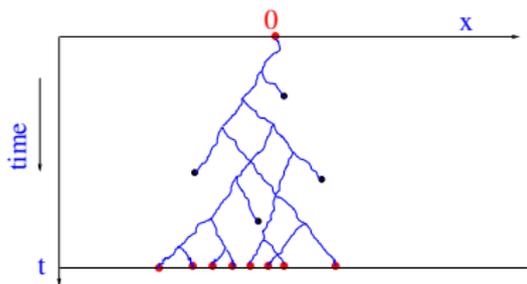
Collaborators: S. N. Majumdar (Orsay) & K. Ramola (Brandeis)

Refs: [Phys. Rev. Lett. 112, 210602 \(2014\)](#)

[Longer version: Chaos, Solitons and Fractals 74, 79 \(2015\)](#)

[Span distribution: Phys. Rev. E 91, 042131 \(2015\)](#)

Branching Brownian Motion with Death



b , d and $D \Rightarrow 3$ parameters of the model

Dynamics: In a small time interval dt , a particle

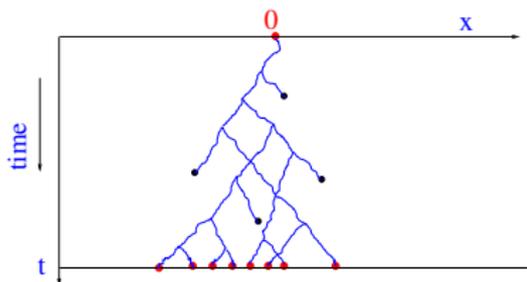
- branches into 2 with proba. $b dt$
- dies with proba. $d dt$
- diffuses with proba. $1 - (b + d) dt$

$$x(t + dt) = x(t) + \eta(t) dt$$

where $\langle \eta(t) \rangle = 0$

$$\langle \eta(t) \eta(t') \rangle = 2D \delta(t - t')$$

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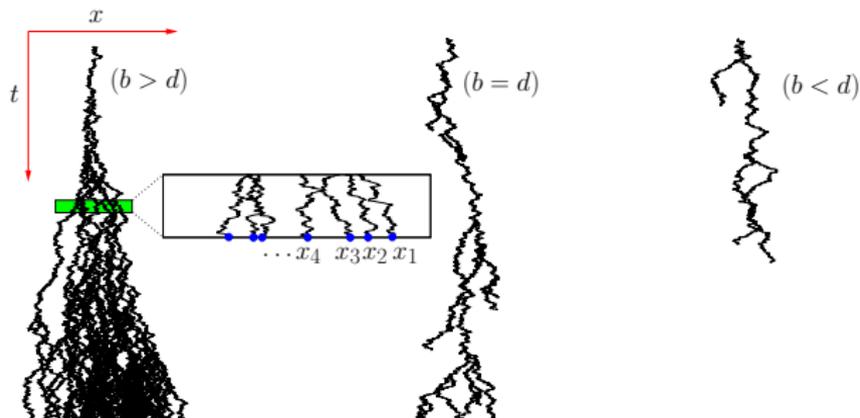
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BBM \Rightarrow prototype model

- Evolutionary system (biology)
- Genealogy
- Cascade model (nuclear physics)
- Directed polymer (statistical physics)
- Epidemic spread,etc.

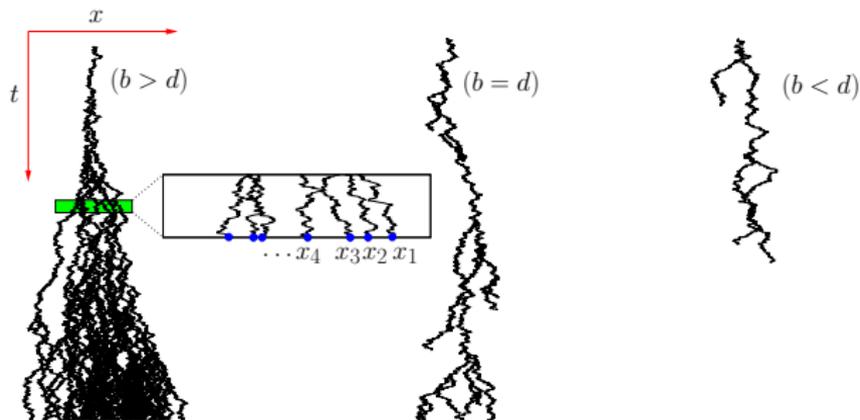
Motivation: extreme and gap statistics

K. Ramola, S. N. Majumdar, G. S., PRL 2014



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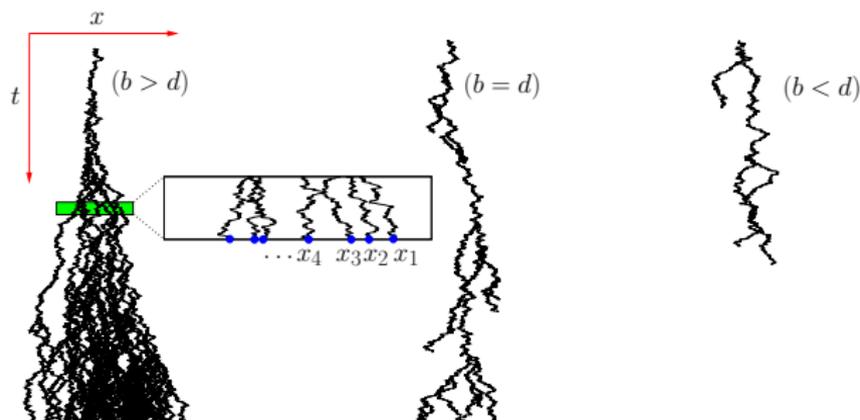
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Q: extreme and gap statistics of BBM ?

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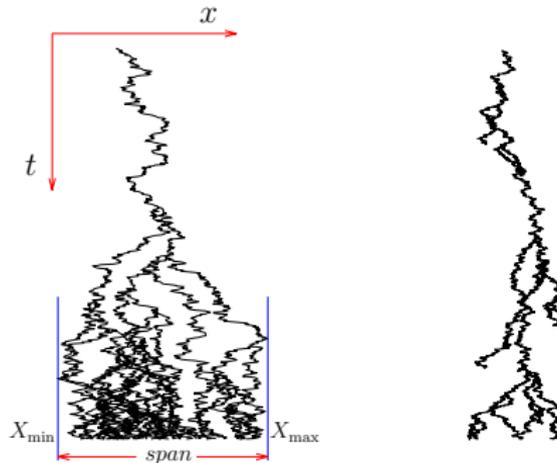


Q: extreme and gap statistics of BBM ?

\Rightarrow Difficult problem because $x_1(t) > x_2(t) > x_3(t) > \dots$ are **strongly correlated**

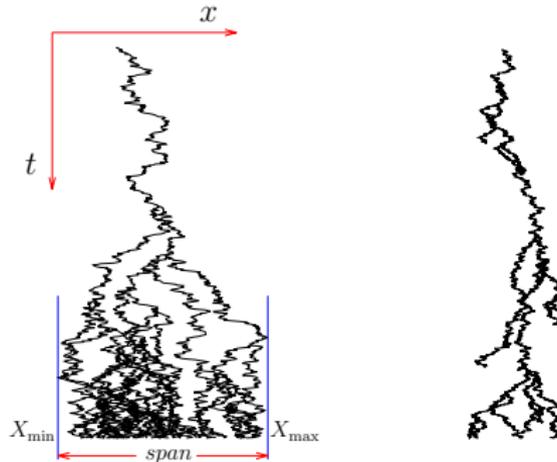
Application of extreme statistics: the **span** of BBM

K. Ramola, S. N. Majumdar, G. S., PRE 2015



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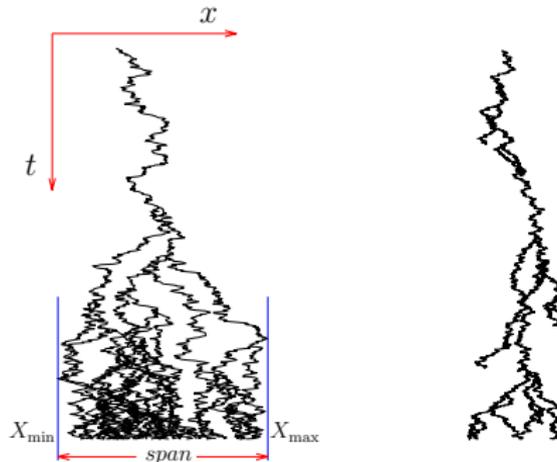
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\implies Application to the **spatial extent** of epidemic spreads

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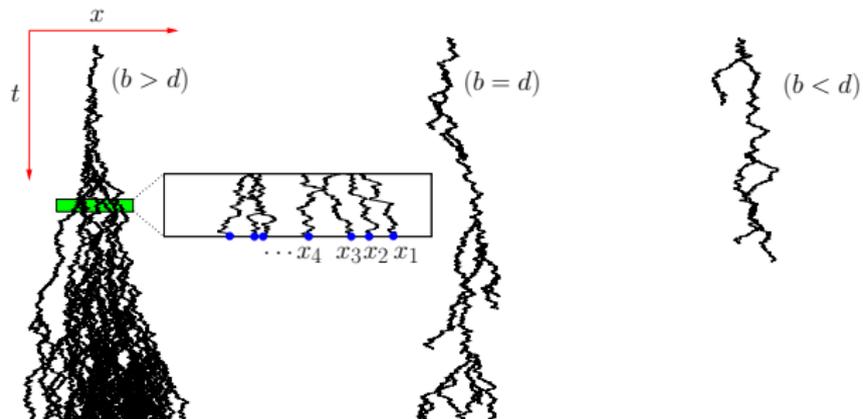
\implies Application to the **spatial extent** of epidemic spreads

see also A. Kundu, S. N. Majumdar, G. S., PRL 2013 (N independent BMs)

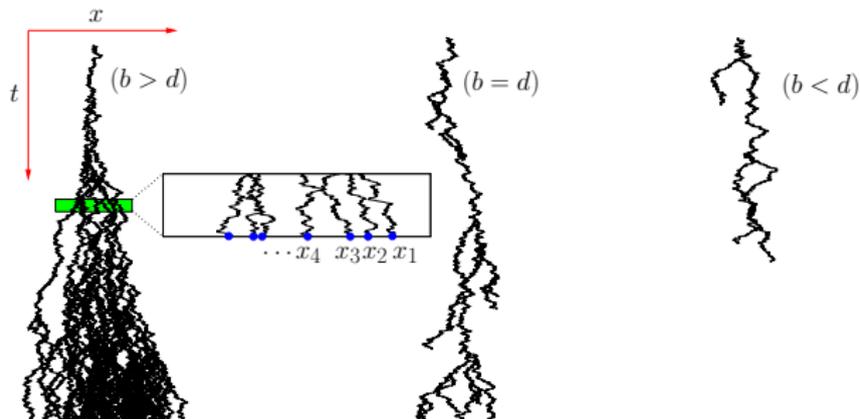
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Evolution of Population Size

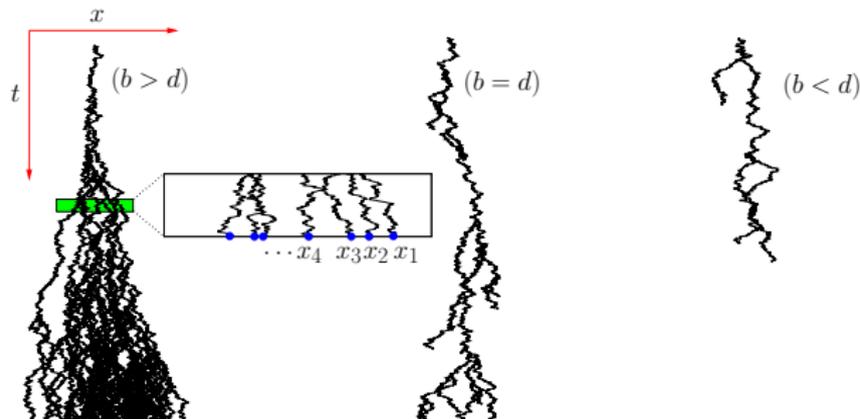


Evolution of Population Size



population size $n(t) \Rightarrow$ random variable

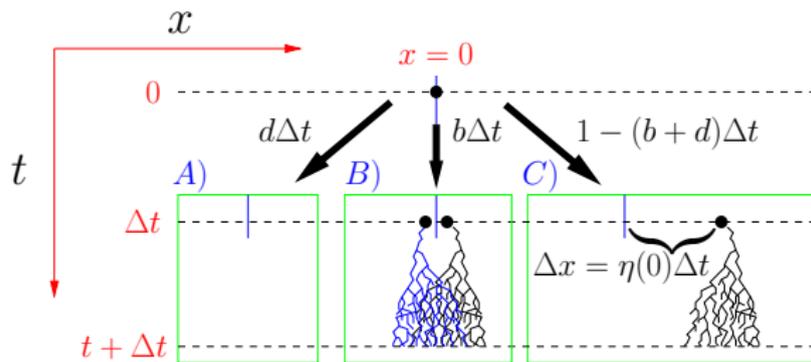
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population size $n(t) \Rightarrow$ random variable

$P(n, t) = \text{Prob.}[n(t) = n]$ satisfies a **backward** equation

Evolution of Population Size: **backward** approach

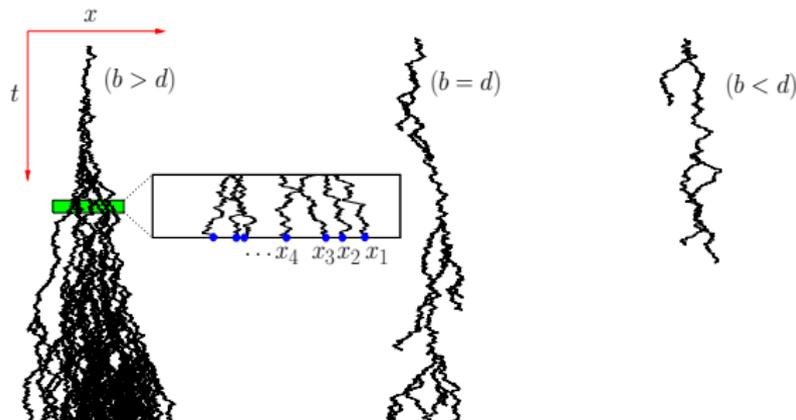


$$P(n, t + \Delta t) = d\Delta t \delta_{n,0} + b\Delta t \sum_{n'=0}^n P(n', t)P(n - n', t) + (1 - (b+d)\Delta t)P(n, t)$$

$$\frac{dP(n, t)}{dt} = -(b+d)P(n, t) + d\delta_{n,0} + b \sum_{n'=0}^n P(n', t)P(n - n', t)$$

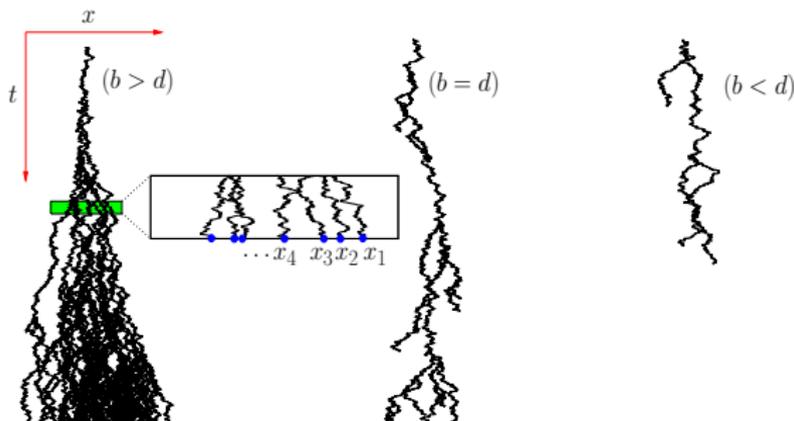
Exact solution via generating function: $\tilde{P}(z, t) = \sum_{n=0}^{\infty} z^n P(n, t)$

Phase transition at $b = d$



$$P(n, t) = (b - d)^2 e^{(b+d)t} \frac{(be^{bt} - de^{dt})^{n-1}}{(be^{bt} - de^{dt})^{n+1}} \quad \text{for } n \geq 1$$
$$= d \frac{(e^{bt} - e^{dt})}{(be^{bt} - de^{dt})} \quad \text{for } n = 0$$

Phase transition at $b = d$

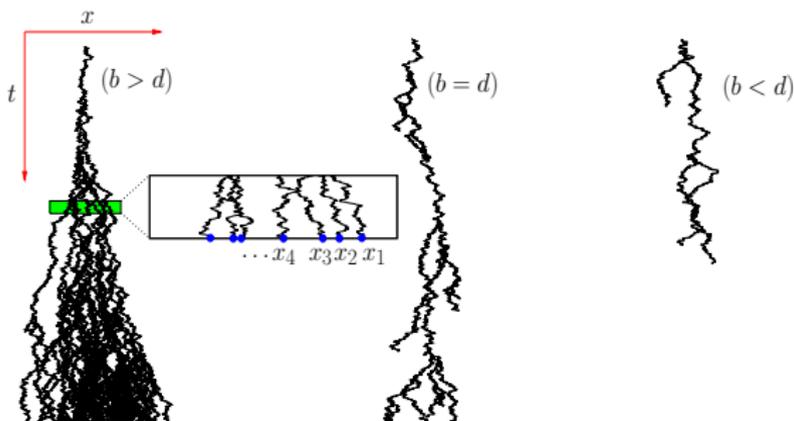


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Consequently:

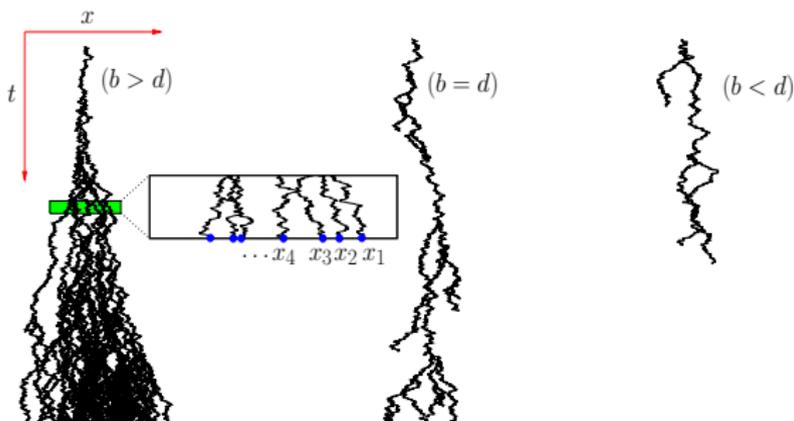
$$\langle n(t) \rangle = e^{(b-d)t} \xrightarrow{t \rightarrow \infty} \begin{cases} +\infty, & b > d \\ 1, & b = d \\ 0, & b < d \end{cases} \quad \begin{array}{l} \text{supercritical} \\ \text{critical} \\ \text{subcritical} \end{array}$$

Typical vs. Average Population Size



- Supercritical phase ($b > d$): $P(n, t) \sim \exp\left[-\frac{n}{n^*(t)}\right]$

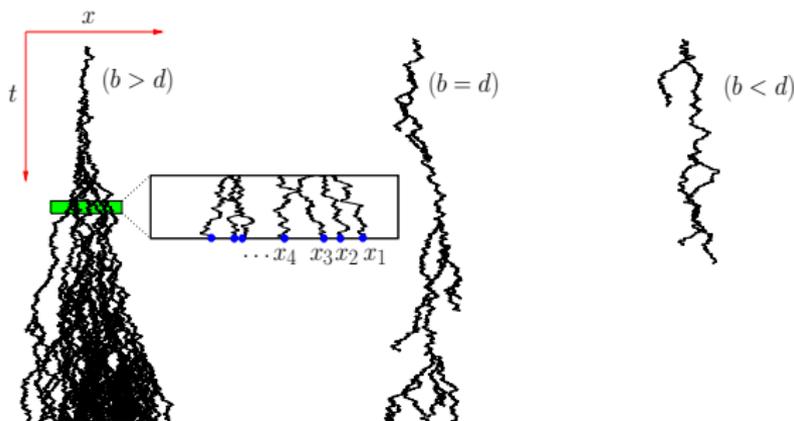
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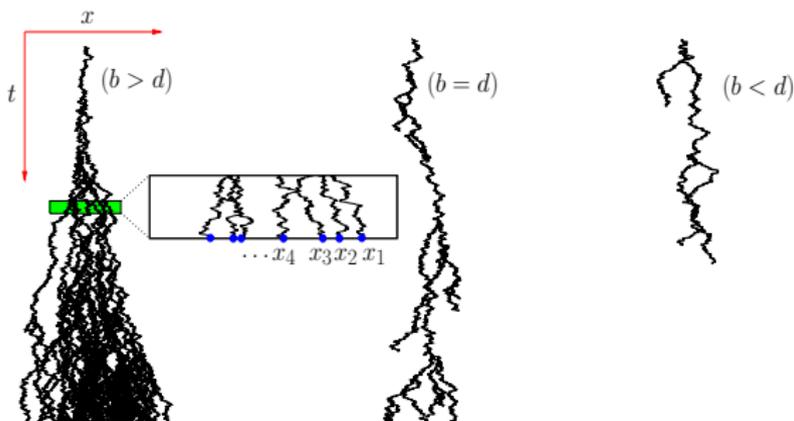


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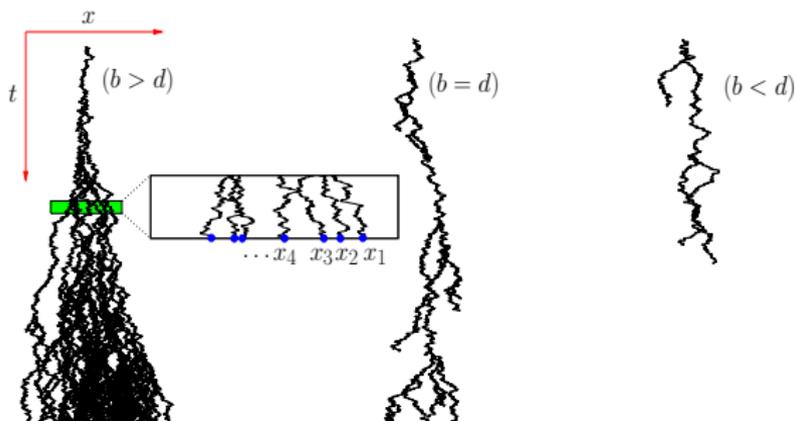
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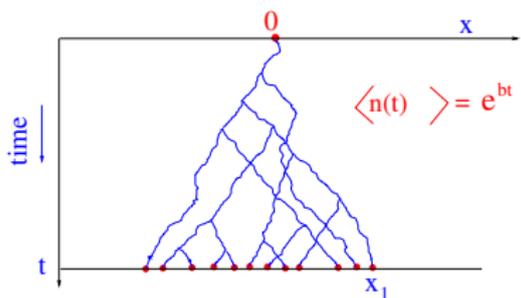
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\Rightarrow Strong fluctuations at the critical point

Outline

- 1 Fluctuations of the particle number
- 2 Order and gap statistics of BBM with no death (a reminder)**
- 3 Extreme and order statistics with death

Extremal statistics with no death $d = 0$

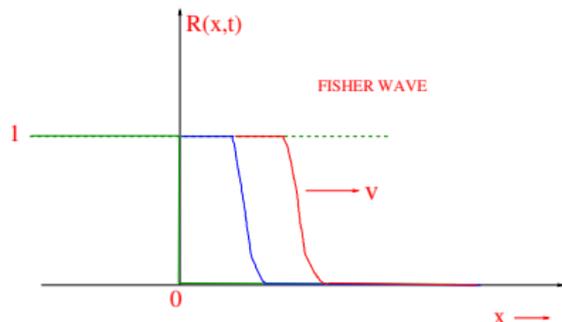


ordered positions (right to left)

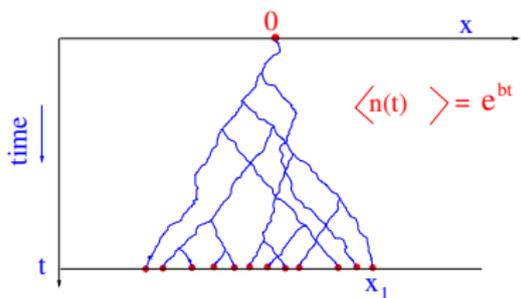
$$x_1(t) > x_2(t) > x_3(t) > \dots$$

$R(x, t) = \text{Prob.}[x_1(t) > x]$ satisfies Fisher-KPP equation:

$$\frac{\partial R}{\partial t} = D \frac{\partial^2 R}{\partial x^2} + bR - bR^2 \quad \text{starting from } R(x, 0) = \theta(-x)$$



Extremal statistics with no death $d = 0$



$$\langle n(t) \rangle = e^{bt}$$

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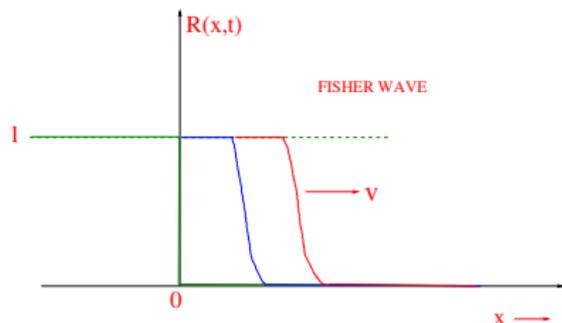
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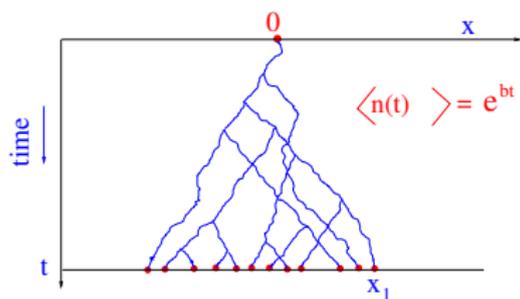
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travelling front solution: $R(x, t) \rightarrow f(x - vt)$ with speed $v = 2\sqrt{bD}$

$$x_1(t) \rightarrow 2\sqrt{bD}t + O(\ln t) \quad \text{at late times}$$



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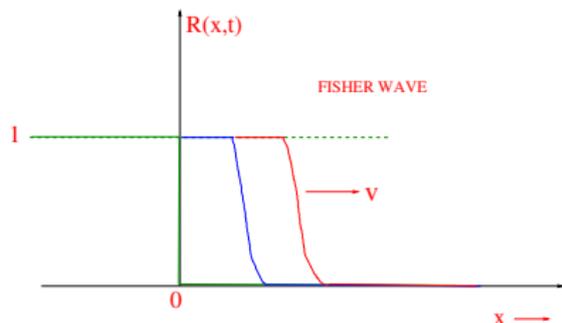
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Mckean '75, Bramson '78, Lalley & Selke '87, Kessler et. al. '97,..., Brunet & Derrida '09,...

see also B. Derrida's talk for large deviations

Order and Gap Statistics for $d = 0$

Brunet & Derrida, 2009-2010

$$x_k(t) \xrightarrow{t \rightarrow \infty} 2\sqrt{bD}t$$

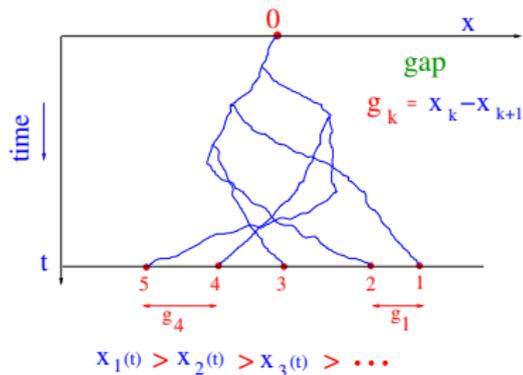
$$\text{Gap: } g_k(t) = x_k(t) - x_{k+1}(t)$$

→ stationary random variable
as $t \rightarrow \infty$

In units of $\sqrt{D/b}$:

$$\langle g_1 \rangle \approx 0.496, \langle g_2 \rangle \approx 0.303, \dots$$

$$\langle g_k \rangle \xrightarrow{k \rightarrow \infty} \frac{1}{k} - \frac{1}{k \ln k} + \dots$$



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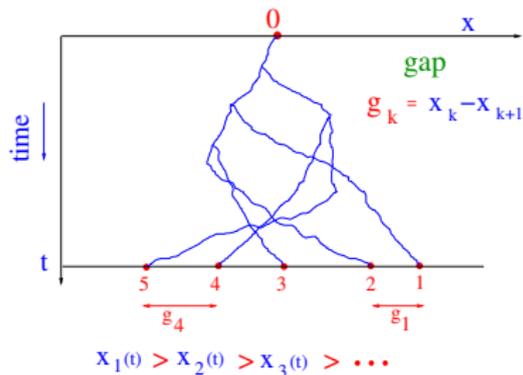
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Heuristic argument for the first gap distribution (Brunet & Derrida, '10):

$$P(g_1, t \rightarrow \infty) \approx \exp \left[-(1 + \sqrt{2})\sqrt{b/D} g_1 \right]$$

Order and Gap Statistics for $d > 0 \rightarrow$ this talk

- **Question:** What happens to **order & gap** statistics when the death rate $d > 0$ is switched on? In particular, at the critical point $d = b$

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- Naturally leads one to study the **Conditioned Ensemble** where the system is conditioned to have a fixed n particles at time t .
- Prob. distr. of any observable \hat{O} in the full problem

$$P(\hat{O}, t) = \frac{\sum_{n>0} Q(\hat{O}, t|n) P(n, t)}{\sum_{n>0} P(n, t)} \approx Q(\hat{O}, t|n^*(t))$$

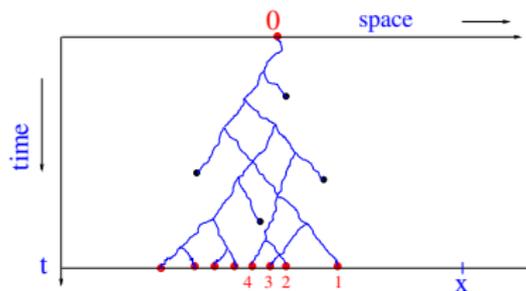
where $Q(\hat{O}, t|n) \rightarrow$ proba. of \hat{O} in the **Conditioned Ensemble**
 $n^*(t) \rightarrow$ **typical** population size

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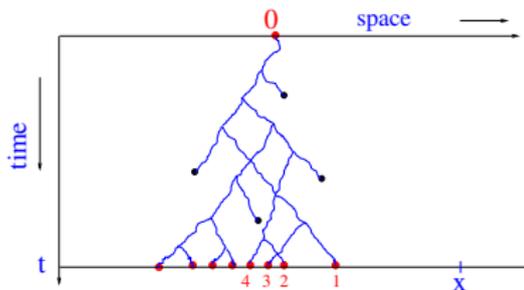
Extremal statistics for $d > 0$ with fixed nber of part.

$C(n, x, t)$ → proba. of having n particles at t with all of them to the left of x



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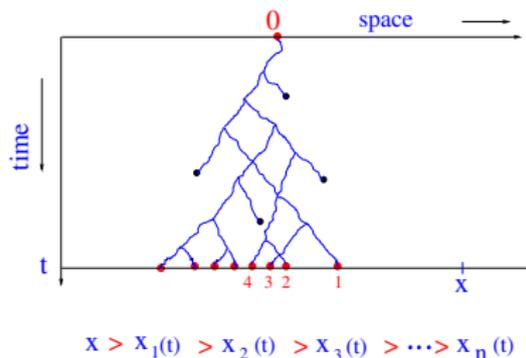


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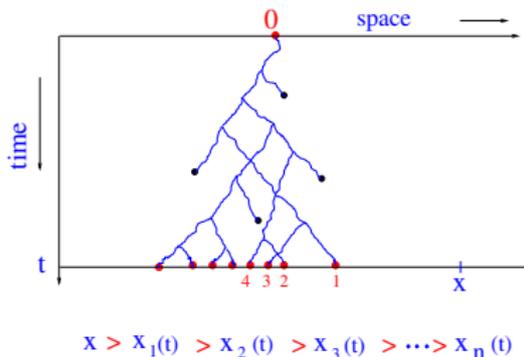
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Conditional prob.

$Q(x, t|n) = \frac{C(n, x, t)}{P(n, t)}$ → given that there are n particles at t , the prob. that all of them are to the left of x

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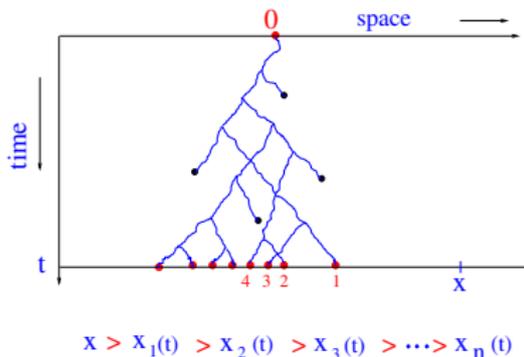
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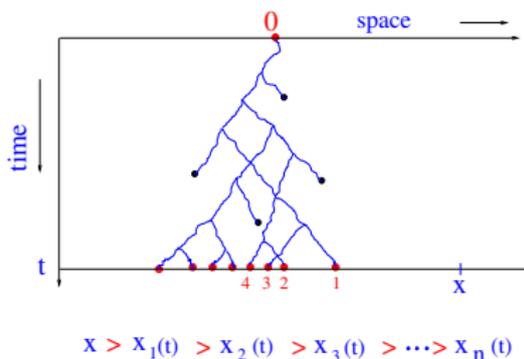
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$$\frac{\partial Q(x, t|n)}{\partial t} = D \frac{\partial^2 Q(x, t|n)}{\partial x^2} + a(t) \sum_{r=1}^{n-1} [Q(x, t|r)Q(x, t|n-r) - Q(x, t|n)]$$

where $a(t) = \frac{(b-d)^2 e^{(b+d)t}}{(e^{bt} - e^{dt})(be^{bt} - de^{dt})}$ and one starts with $Q(x, t|0) = 1$

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⇒ **linear** in $Q(x, t|n)$ and can be solved **recursively**

Order Statistics at the Critical Point $b = d$

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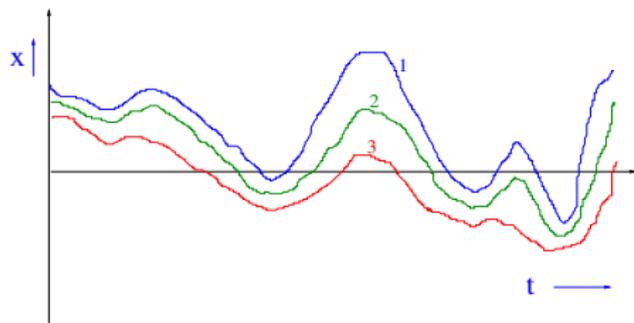
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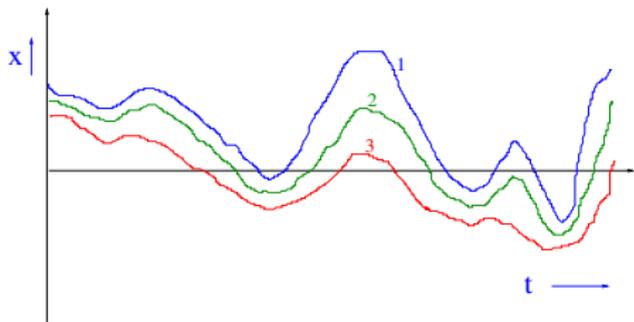
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- Similarly the **second**, the **third**...all **diffuse**: $P(x_k, t|n) \rightarrow \frac{1}{\sqrt{4\pi Dt}} \exp \left[-\frac{x_k^2}{4Dt} \right]$

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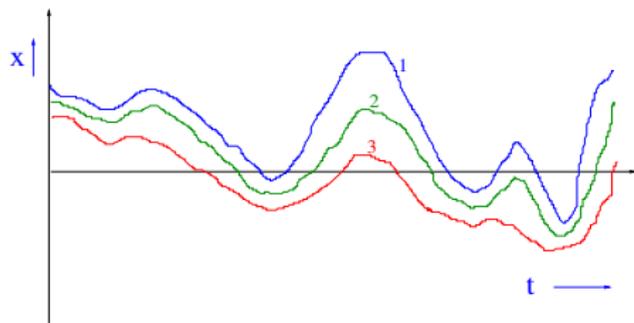


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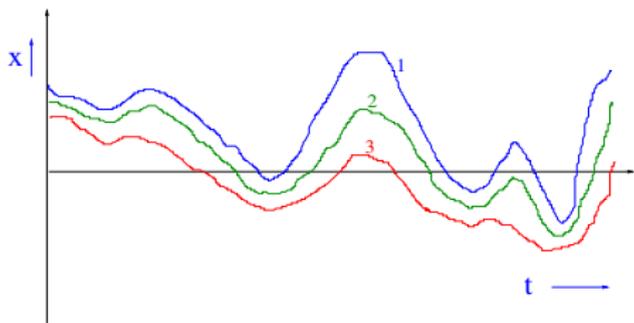


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Exact computation of the gap distribution in the **Conditioned Ensemble**
(fixed n)

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[K. Ramola, S. N. Majumdar & G. S., PRL **112**, 210602 (2014)]

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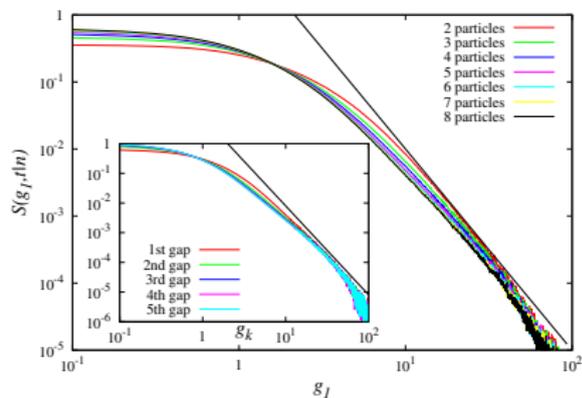
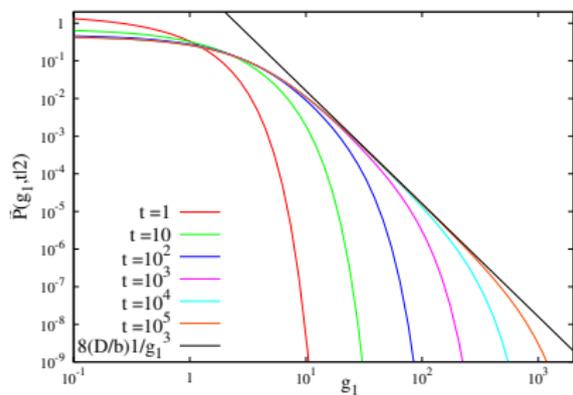
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Numerical simulations at the critical point $b = d$



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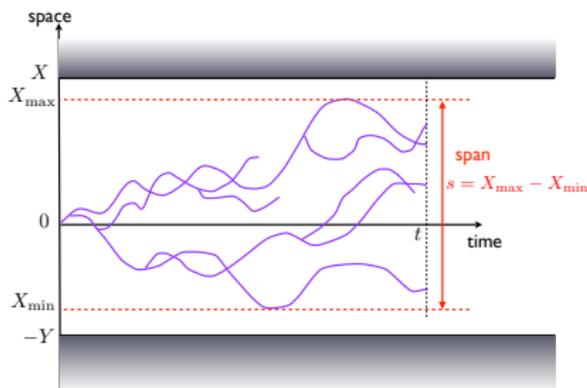
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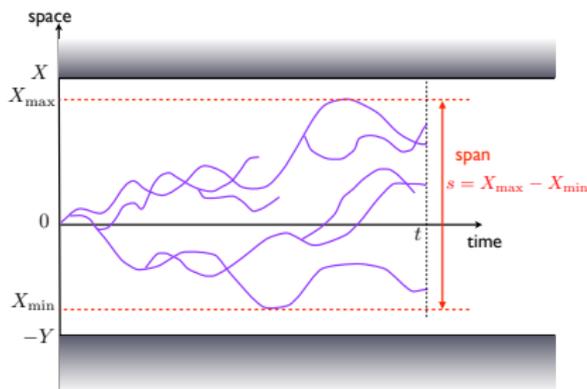


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[K. Ramola, S. N. Majumdar & G. S., PRE 2015]