# Non Gaussian Log Correlated fields in RMT 

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Based on joint works with Elliot Paquette, Nick Cook, Fanny Augeri and Raphael Butez

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## Asymptotically Gaussian fields in random matrix theory

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Central limit theorem

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Central limit theorem $f: \mathbb{R} \rightarrow \mathbb{R}$ compactly supported, smooth. Consider

$$
W_{f, N}=\sum_{i=1}^{N} f\left(\lambda_{i}\right)-N \int f d \sigma
$$

## CLT

Theorem (Johansson '98; $\beta$ ensembles)
$W_{f, N}$ satisfies CLT, mean $(2 / \beta-1) \int f d \nu$, variance

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\frac{(2 / \beta)}{4 \pi^{2}} \iint_{-2}^{2} f(t) f^{\prime}(s) \frac{\sqrt{4-s^{2}}}{(t-s) \sqrt{4-t^{2}}} d s d t
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CLT's of this type go back at least to CLT of Jonsson for moments ('82), Pastur and co-workers, Bai-Silverstein, Shcherbina, ....

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If it is log-correlated, what about the extrema?

## CUE char poly



Figure 1: Realizations of $\log \left|\mathrm{P}_{N}\left(e^{\mathrm{i} h}\right)\right|, 0 \leq h<2 \pi$, for $N=50$ and $N=1024$. At microscopic scales, the field is smooth away from the eigenvalues, in contrast with the rugged landscape at mesoscopic and macroscopic scales.
(From Arguin, Belius, Bourgade '17)

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where $W$ has the law of the sum of two independent Gumbels.

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Both use in essential way CUE (aka $\beta=2$ ), where joint distribution of eigenvalues is

$$
\prod_{i<j}\left|\lambda_{i}-\lambda_{j}\right|^{2}
$$

for which Gaussianity of traces follows from Diaconis-Shashahani and moments ofdeterminant (=exponential moments of $M_{N}(z)$ ) are Toeplitz determinants.

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The clincher:

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\end{array}\right)\binom{\Phi_{k}(z)}{\Phi_{k}^{*}(z)}, \Phi_{k}^{*}(z)=z^{k} \overline{\Phi_{k}\left(\bar{z}^{-1}\right)} .
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$\alpha_{k}=B_{k} e^{2 \pi i \theta_{k}}, E B_{k}^{2} \sim 2 / \beta k$, beta variable. $\alpha_{k} \sim g_{k}+i g_{k}^{\prime}$, Gaussian.

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$\alpha_{k}=B_{k} e^{2 \pi i \theta_{k}}, E B_{k}^{2} \sim 2 / \beta k$, beta variable. $\alpha_{k} \sim g_{k}+i g_{k}^{\prime}$, Gaussian.
In addition, $\sup _{|z|=1}|\log | M_{N}(z)|-\log | \Phi_{k}^{*}(z)| |$ is tight.

## Recursions in the lab

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\begin{gathered}
\log \Phi_{k}^{*}\left(e^{i \theta}\right)-\log \Phi_{k-1}^{*}\left(e^{i \theta}\right)=\log \left(1-\alpha_{j} e^{i \Psi_{k-1}}\right) \sim-\alpha_{j} e^{i \Psi_{k-1}(\theta)} \\
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Work in progress: Paquette-Z ('20?) Convergence in law of max $\log \left|\Phi_{N}^{*}\left(e^{i \theta}\right)\right|$ to Gumbel shifted by (unknown) r.v..

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where $C_{\ell}=$ number of cycles of length $\ell$, essentially Poisson. Almost independent additive structure. But there are arithmetic issues.
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Field is still log-correlated.

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## Random permutation char poly



(A) $N=100, I=(0,1)$
(B) $N=10^{4}, I=(0,1)$

(C) $N=10^{4}, I=(0.1,0.11)$
(D) $N=10^{9}, I=(0,0.6)$

Figure 1. Simulations of the field $X_{N}$ on subintervals $I \subset \mathbb{T}$, computed from the cycle structures for the permutations $P_{N}$ using the formula (2.2)). The cycle structures are random partitions of $[N]$ generated using the Chinese restaurant process. The respective partitions are:
(A): $\{56,22,9,9,4\}$,
(B/C): $\{6310,1914,909,668,79,47,33,19,12,5,3,1\}$,
(D): $\{892060223,78087020,19479718,9152317,630684,352623,114502,104059$, $8973,8193,1641,33,5,3,2,2,1,1\}$.
In (D) there are noticeable dips in the field near the rationals $0,1 / 2$, and $1 / 3$.
(From Cook, Zeitouni '17)

## Towards the wild: $\mathbf{G} \beta \mathbf{E}$

We take $X_{N} \sim \mathrm{G} \beta \mathrm{E}$, ie joint distribution of eigenvalues on $\mathbb{R}^{N}$

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\prod_{i<j}\left|\lambda_{i}-\lambda_{j}\right|^{\beta} e^{-\beta \frac{N}{4} \sum \lambda_{i}^{2}}
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Also, connection to GMC for $\beta=2$ : Berestycki-Webb-Wong '18 ( $L^{2}$ phase) For general $\beta$ : even CLT of log-det not clear!

## log-det trajectory

Beta=2, Matrice de taille $\mathbf{n = 1 0 0 0}$ Pas de temps= $\mathbf{0 . 0 0 1}$


## Empirical facts



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Skewed?

## Reason for skewness in simulations



## CLT for log determinant $\mathrm{G} \beta \mathrm{E}$

The case $z=0$ is special.
Theorem (Tao-Vu '11)
$\left(M_{N}(0)-N \int \log |z-x| \sigma(d x)-a_{\beta} \log N\right) / \sqrt{\log N}$ converges (for Wigner matrices, 4 matching moments) to standard Gaussian.

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Bourgade-Mody '19: extends w/out matching 4 moments. By replacement principle, the key step in the TV proof is the result for $\mathrm{G} \beta \mathrm{E}, \beta=1,2$. Their proof extends to general $\beta>0$, and is based on recursions.

## The Dumitriu-Edelman representation

## Theorem (Dumitriu-Edelman '05)

$X_{N}$ from $G \beta E$ is unitarily equivalent to the following 3-diagonal Jacobi matrix

$$
\frac{1}{\sqrt{N}} X_{N}=\frac{1}{\sqrt{N}}\left(\begin{array}{lllll}
b_{1} & a_{1} & 0 & \ldots & 0 \\
a_{1} & b_{2} & a_{2} & 0 & \cdots \\
0 & a_{2} & b_{3} & a_{3} & 0 \\
\cdots & \cdots & \cdots & \cdots & \cdots \\
0 & 0 & 0 & a_{N-1} & b_{N}
\end{array}\right)
$$

where $b_{i} \sim N(0, \sqrt{2 / \beta}), a_{i} \sim \chi_{i \beta} / \sqrt{\beta}$.
Here $a_{i} \sim \chi_{i \beta} / \sqrt{\beta}$ means $a_{i}^{2}$ has chi-square distribution with $i \beta$ degrees of freedom, ie $\chi_{i \beta} \sim \sqrt{i}+\sqrt{1 / 2 \beta} G+O(1 / i)$.

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We set

$$
\Psi_{k}(z)=\phi_{k}(z \sqrt{N}) \frac{e^{-k \beta U(z \sqrt{N / k})+c k}}{\sqrt{k!}}
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and then

$$
\Psi_{k}(z)=\left(z \sqrt{N}-b_{k}\right) \Delta_{k} \Psi_{k-1}(z)-a_{k-1}^{2} \Delta_{k} \Delta_{k-1} \Psi_{k-2}(z)
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Here, $\Delta_{k}=1 / \sqrt{k} \alpha\left(t_{k}\right), t_{k}=z \sqrt{N / k}$, and $\alpha(t)=1$ for $t<2$ and $\alpha(t)=\sqrt{t^{2} / 4-1}+t / 2$ for $t \geq 2$.

## Recursions

Set $k_{0}$ so that $t_{k_{0}}=2$ (if $z=0$ then $k_{0}=1$ ). In matrix form, for $k \geq k_{0}$,

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\begin{aligned}
& \binom{\Psi_{k+1}(z)}{\Psi_{k}(z)} \\
& \sim\left(\begin{array}{ll}
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\end{array}\right)\binom{\Psi_{k}(z)}{\Psi_{k-1}(z)}+\left(\begin{array}{ll}
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Tao-Vu show that $\Psi_{k-1}(z)^{2}+\Psi_{k-1}(z)^{2}$ (essentially) forms a martingale with quadratic variation process of increment $\sim 1 / k$. This gives the CLT.

## Recursions - general $z$

The following is joint work with Fanny Augeri and Raphael Butez, in progress. We have a CLT for log-characteristic polynomial, and work on the log-correlated structure.

## Recursions - general $z$

There are several regimes to consider. Fix $\epsilon>0$, recall that $k_{0}=z^{2} N / 4 k$.

- $k<\epsilon k_{0}$ : one easily checks that $\Psi_{k}(z) \sim 1$.
- $k \in\left[\epsilon k_{0}, k_{0}\right]$ : write

$$
X_{k}=\Psi_{k} / \Psi_{k-1}=1+\delta_{k}, \quad X_{k}=A_{k}+B_{k} / X_{k-1}
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for appropriate $A_{k}, B_{k}$.

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for appropriate $A_{k}, B_{k}$. In this regime, $\delta_{k} \sim 0$ and one obtains a recursion

$$
\delta_{k} \sim u_{k}+v_{k} \delta_{k-1}
$$

where $u_{k} \sim b_{k} / \sqrt{k \alpha_{k}^{2}}+1 / 2 k \alpha_{k}^{2}-g_{k} / \sqrt{k \alpha_{k}^{4}}$,
$v_{k}=\left(1-1 / 2 k+g_{k} / \sqrt{k}\right) / \alpha_{k}^{2}$, which one solves.

- $k>k_{0}$ : Oscillatory regime, most interesting.


## Recursions - general $z$ - the scalar regime

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\delta_{k}=\sum_{j=2}^{k} u_{j} \prod_{\ell=j+1}^{k} v_{\ell}
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is a martingale, and small. We need to compute $\sum \delta_{k}$, and $\delta_{k}$ are correlated!.

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## Recursions - general $z$ - the oscilatory regime

$$
X_{k}=\binom{\Psi_{k+1}}{\Psi_{k}}, k>k_{0}
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We have

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X_{k+1}=\left(A_{k}+W_{k}\right) X_{k},
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where,

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A_{k}=\left(\begin{array}{cc}
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$z_{k}=z \sqrt{\frac{n}{k}}=2-\frac{1}{k_{0}}$ and $b_{k} \sim \mathcal{N}(0,2 / \beta)$ and $g_{k} \sim \mathcal{N}(0,2 / \beta)$.

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$z_{k}=z \sqrt{\frac{n}{\hbar}}=2-\frac{1}{k_{0}}$ and $b_{k} \sim \mathcal{N}(0,2 / \beta)$ and $g_{k} \sim \mathcal{N}(0,2 / \beta)$.
Eigenvalues of $A_{k}$ for $k>k_{0}$ are complex of (essentially) unit norm. Change basis to eigenvector basis, get

$$
\hat{X}_{k}=Q_{k} \prod_{i=k_{0}}^{k-1} Q_{i+1}^{-1} Q_{i}\left(R_{i}+\hat{W}_{i}\right) Q_{k_{0}}^{-1} \hat{X}_{k_{0}}
$$

where $R_{i}$ are rotation matrices of angle $\theta_{k} \sim \sqrt{k / k_{0-1}}$.

## Recursions - general $z$ - the oscilatory regime

First order approximation: divide to blocks of length $\ell_{i}=\left(k_{0} / i\right)^{1 / 3}$, linearize in each block, and get contribution to variance of order $1 / i$.

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Caveat: unlike the case of $z=0$, the quadratic variation of the (log) of the norm is not a function of the norm!

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Caveat: unlike the case of $z=0$, the quadratic variation of the (log) of the norm is not a function of the norm!
Solution: along block we have $\prod R_{i}=I$, but the vector $(1,0)^{T}$ is not mapped to $\rho_{i}(1,0)$ due to the noise. So instead, stop (at random time) where

$$
\prod_{i=l}^{\ell_{i+1}} Q_{i+1}^{-1} Q_{i}\left(R_{i}+\hat{W}_{i}\right)(0,1)^{T} \sim \rho_{i}(0,1)^{T}
$$

## Recursions - general $z$ - the oscilatory regime

First order approximation: divide to blocks of length $\ell_{i}=\left(k_{0} / i\right)^{1 / 3}$, linearize in each block, and get contribution to variance of order $1 / i$.
Caveat: unlike the case of $z=0$, the quadratic variation of the (log) of the norm is not a function of the norm!
Solution: along block we have $\prod R_{i}=I$, but the vector $(1,0)^{T}$ is not mapped to $\rho_{i}(1,0)$ due to the noise. So instead, stop (at random time) where

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\prod_{i=\ell_{i}}^{\ell_{j+1}} Q_{i+1}^{-1} Q_{i}\left(R_{i}+\hat{W}_{i}\right)(0,1)^{T} \sim \rho_{i}(0,1)^{T}
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We have $\ell_{j+1}-\ell_{j} \sim\left(k_{0} / j\right)^{1 / 3}$, and variance computation as in sketch.

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Complication when blocks get too small - cannot ensure the approximation;
But variance is small there, so can combine blocks!
Computing correlation between different $z s$ is complicated in the regime $\left|z-z^{\prime}\right|<N^{-2 / 3}$ because of block structure.

