



Theory of robust many-body quantum scars

In long-range interacting systems

arXiv:2309.12504, in press on PRX

Started in January 2020:

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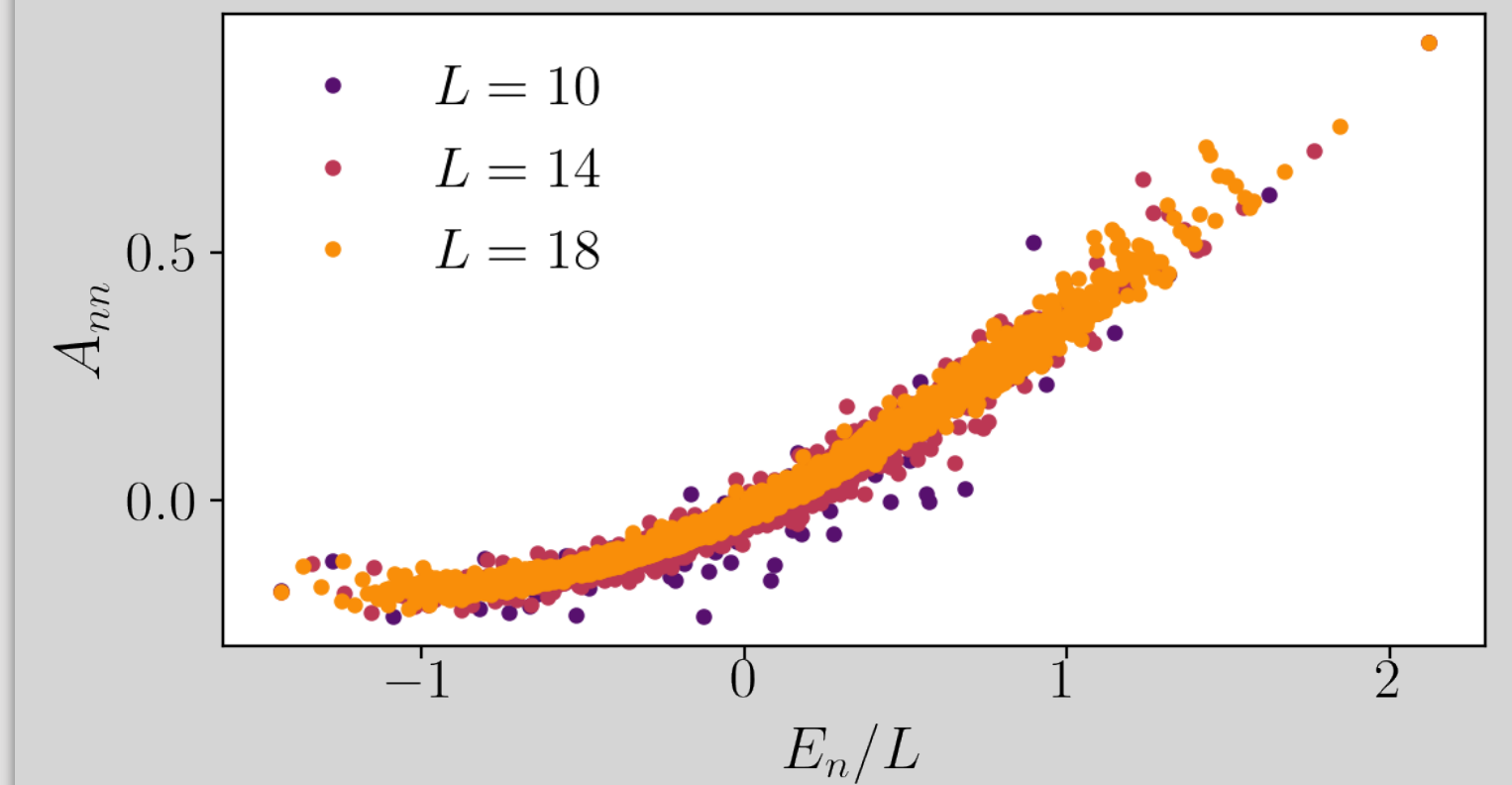
Dmitry Abanin
Princeton

Non-equilibrium states of matter

Fast thermalization

Quantum Many-Body chaos
(Eigenstate Thermalization Hypothesis)

Observables $\langle E_n | A | E_n \rangle$ smooth functions of E_n/L



Prethermalization
Fragmentation

Many-Body quantum Scars

This talk

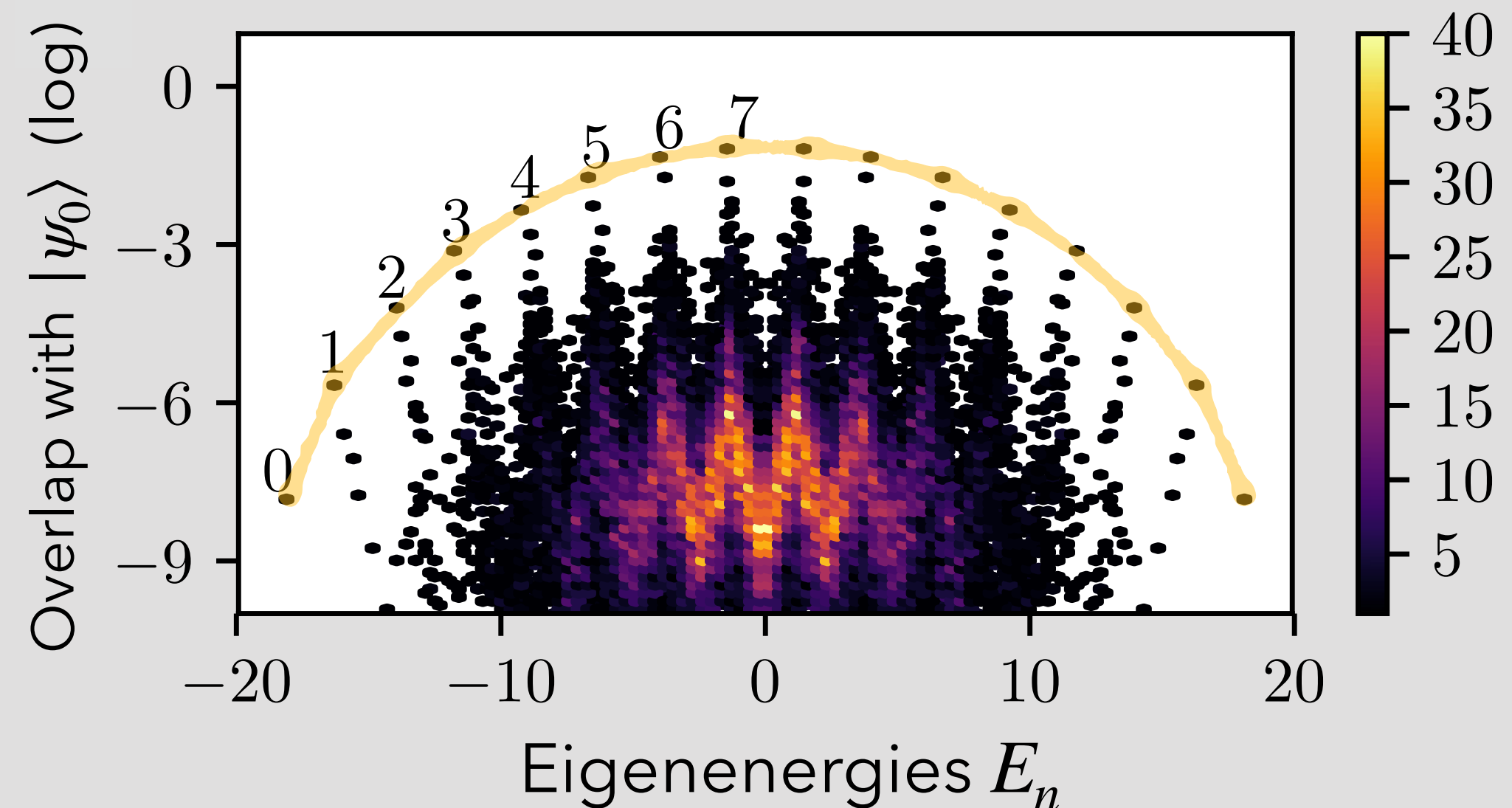
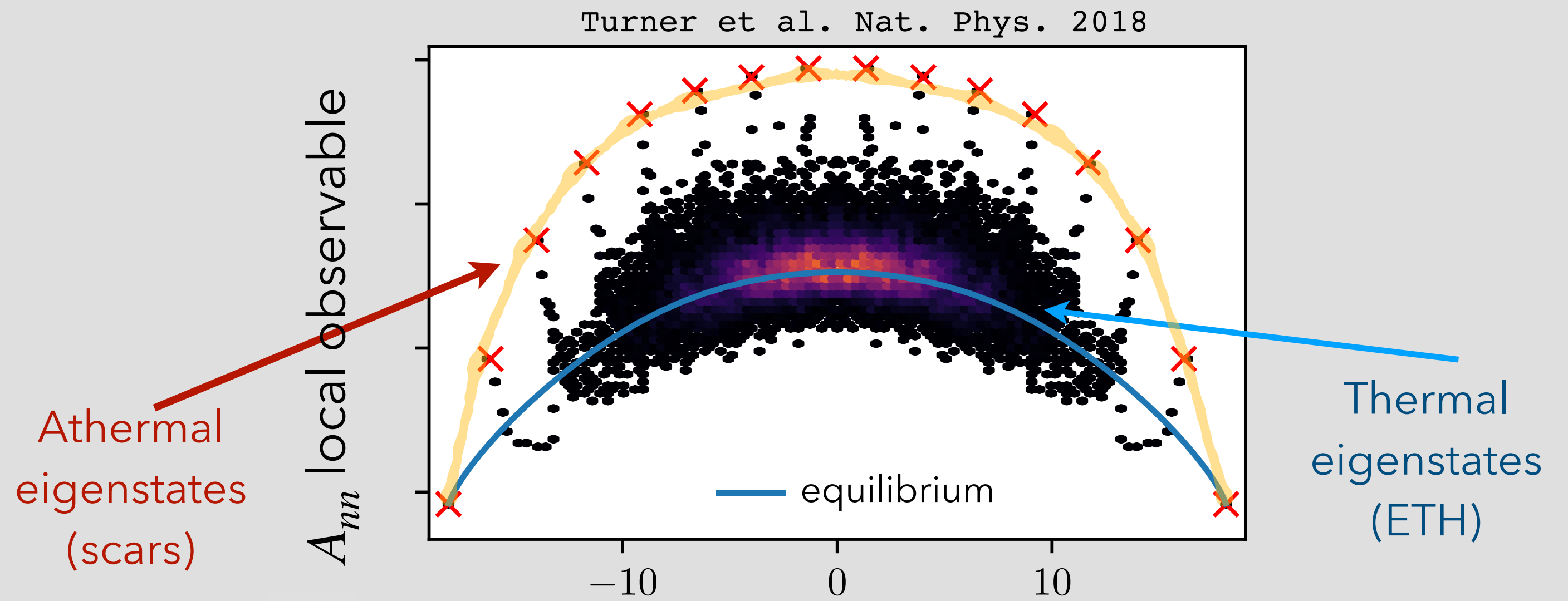
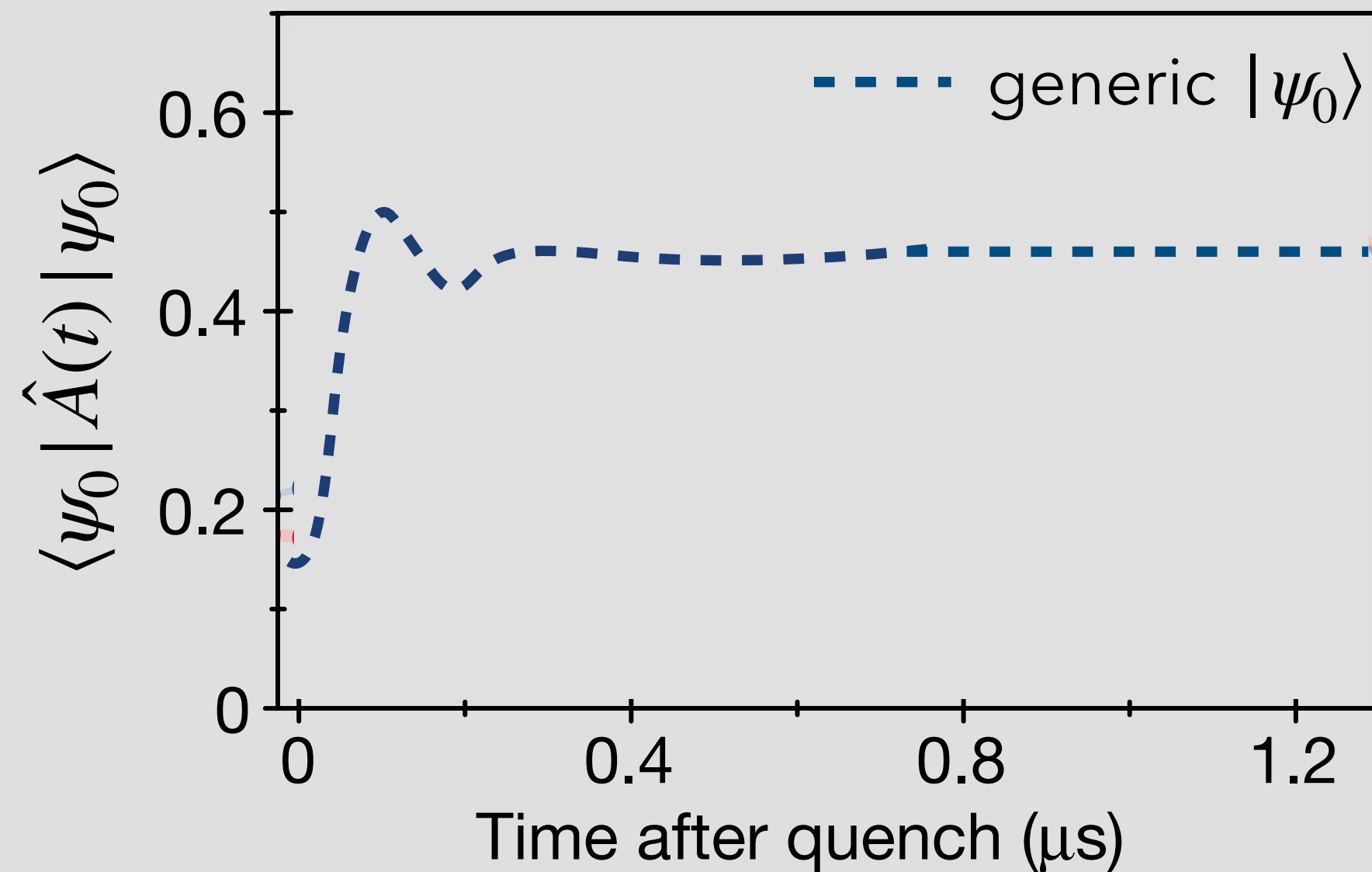
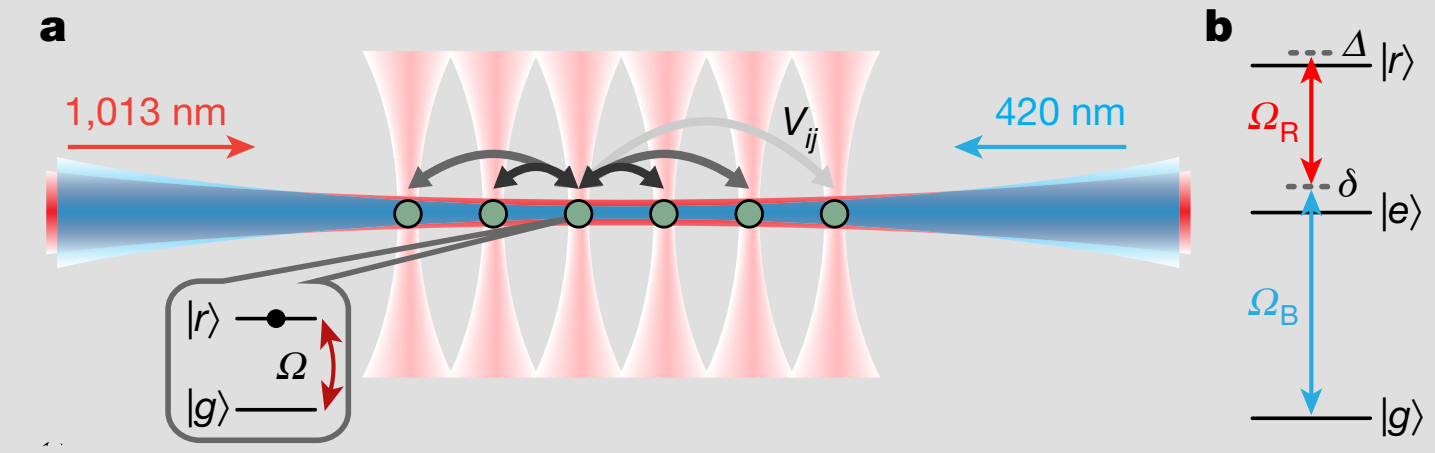
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Weak ergodicity breaking

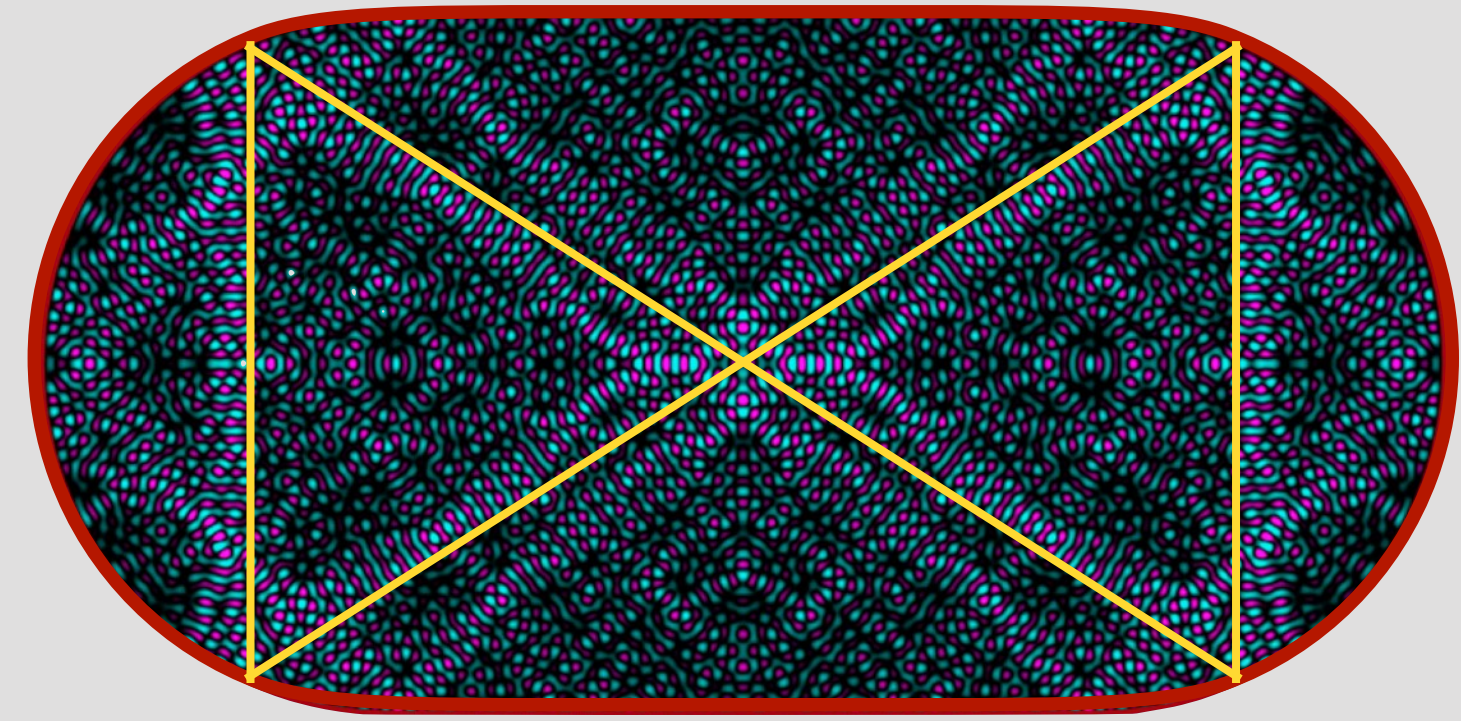
Many-body localization,
Quantum glasses

Non-ergodic dynamics

Many-body quantum scars

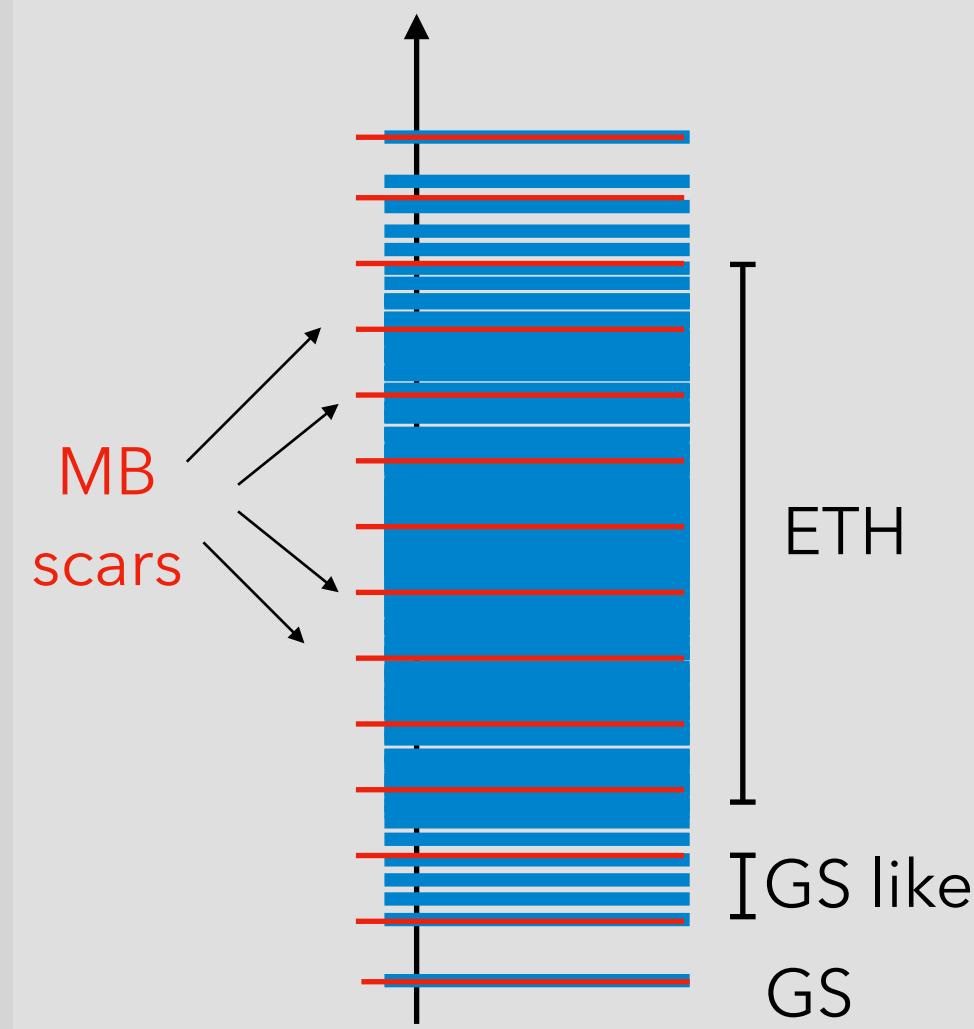


Inspired by "quantum scars" in billiards $\hbar \ll 1$ Heller 1974



Theory of many-body quantum scars

Turner et al. Nat. Phys. 2018 + Papic, Moudgalya, Chandran, Iadecola, Müller, ...

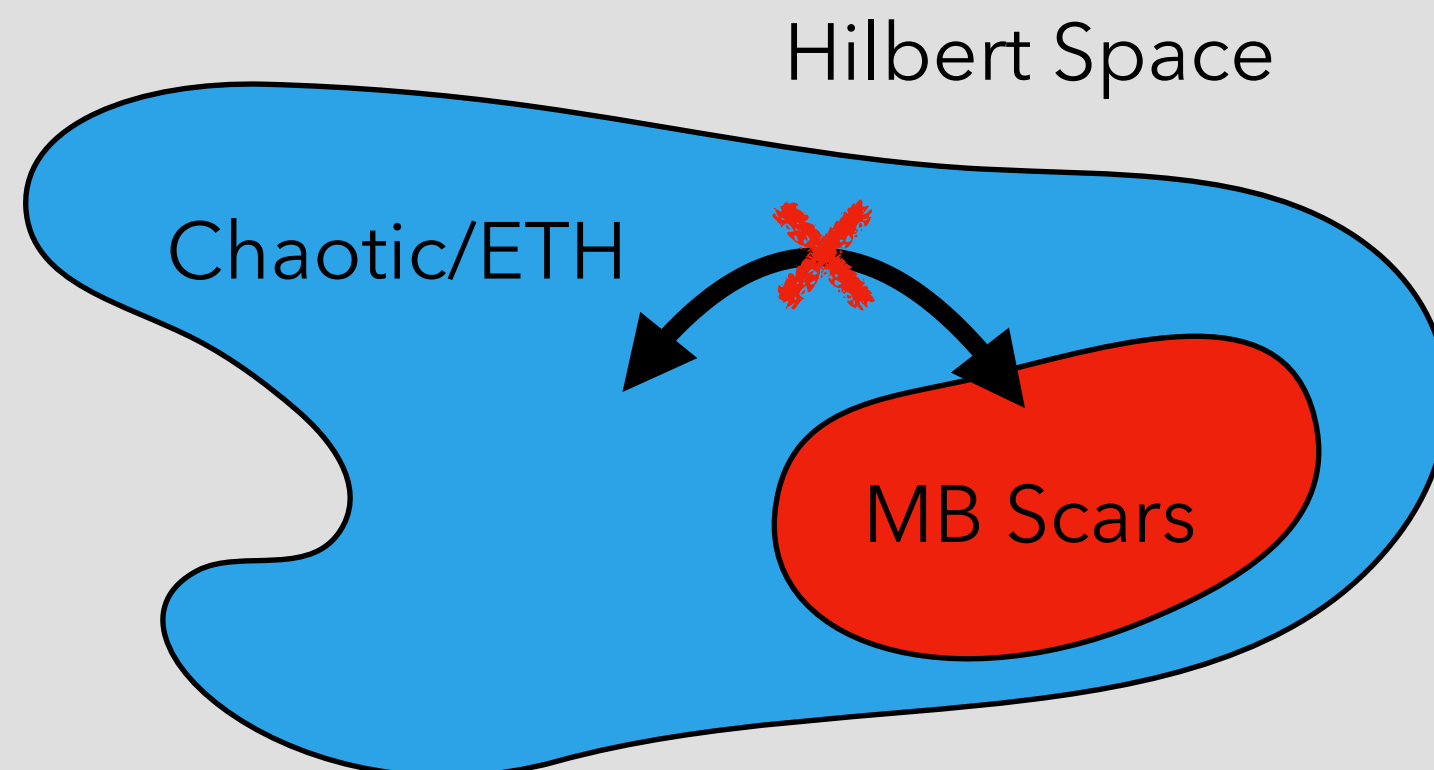
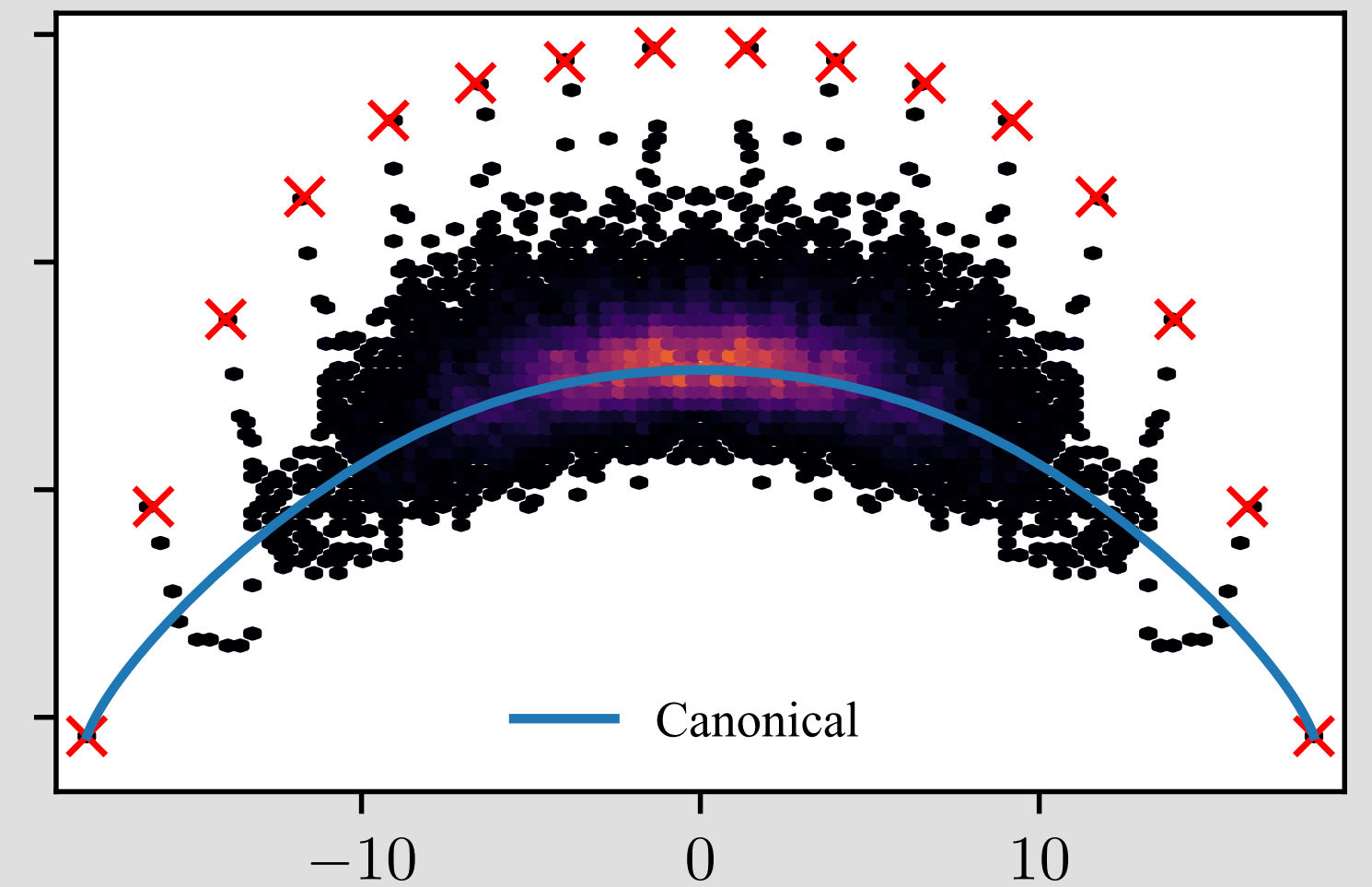


- construction of local \hat{H} with exact scars

Shiraishi & Mori - PRL (2017)
Motrunich, Moudgalya, Iadecola, Surace, ...

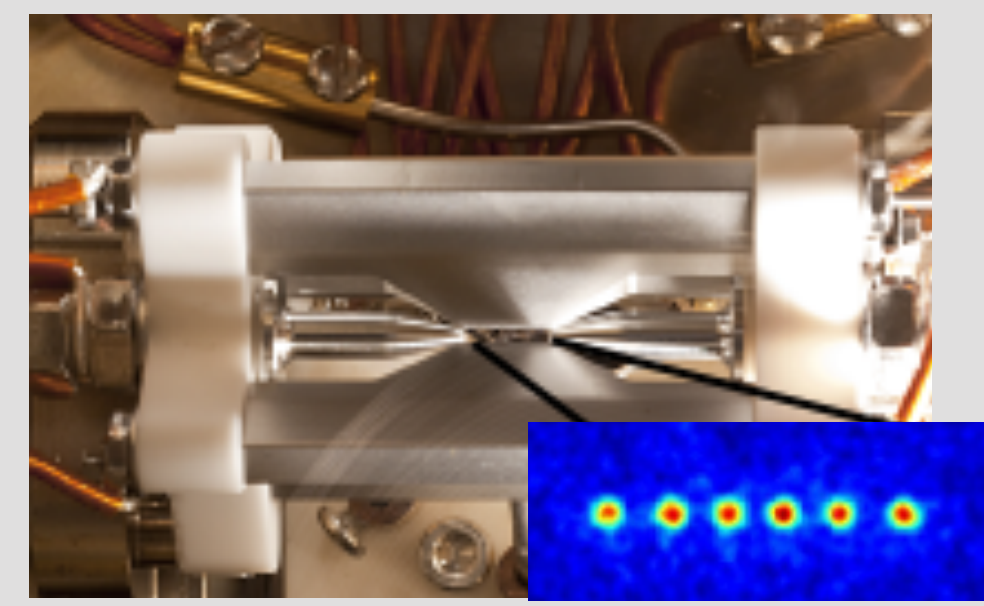
- **unstable** to generic perturbations

Lin et al. - PRR (2020), Surace et al. - PRB (2021)



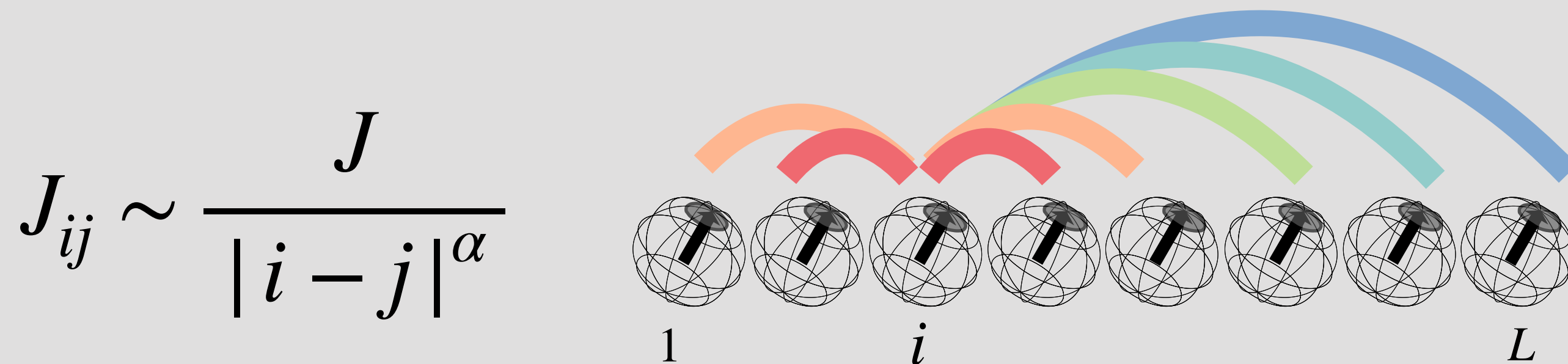
Is there a class of systems with *robust* scars?

Long-range interacting systems



$$0.5 < \alpha < 1.8$$

- Two level systems (spins 1/2 or “qubits”)
- Interactions mediated by spatially delocalized degrees of freedom



$$J_{ij} \sim \frac{J}{|i-j|^\alpha}$$

Variable-range quantum Ising chain

$$H = - \sum_{i < j}^L J_{ij} s_i^x s_j^x - h \sum_i^L s_i^z$$

Quantum experiments in AMO physics

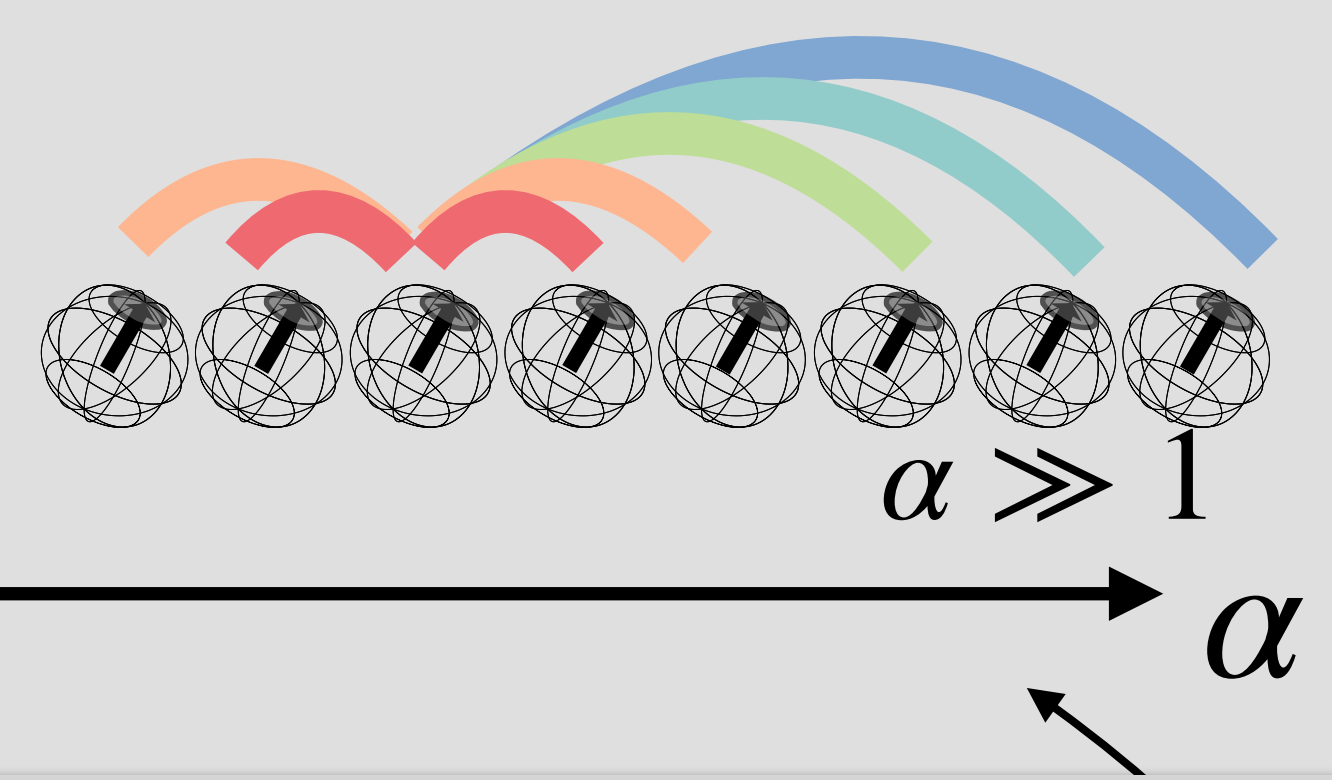
Unusual quantum dynamics:

Defenu et al. Rev. Mod. Phys 2023
Defenu, Lerose, **SP** - Phys. Rep. 2024

Dynamics in long-range systems

$\alpha = 0$

$\alpha = d$

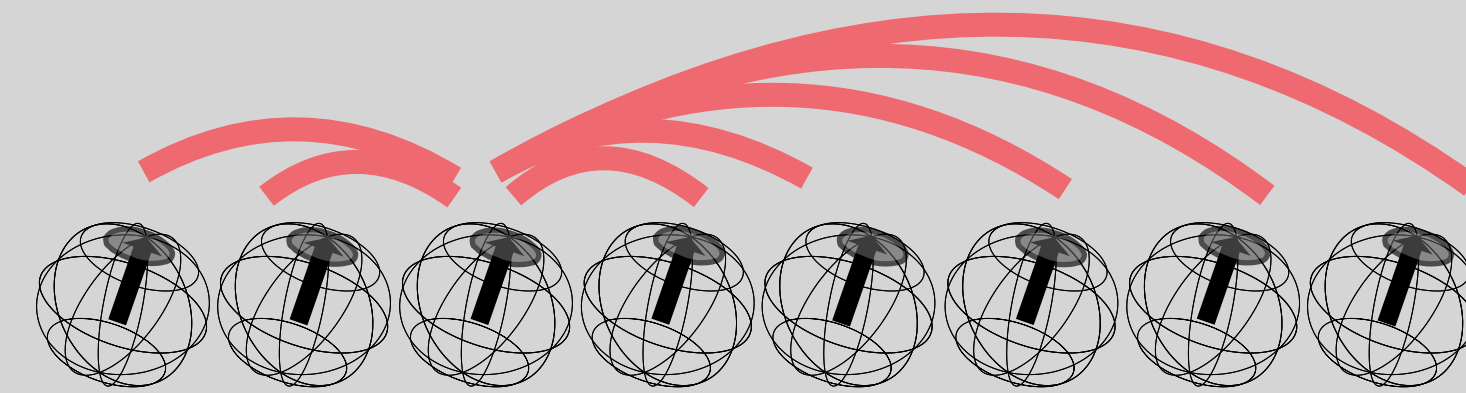


Mean field Dynamics (Solvable) $J = J_0/N_\alpha$

$$\hat{H}_{\alpha=0} = -J \left(\sum_{ij} s_i^x s_j^x \right) - h \sum_i s_i^z$$

$(\hat{S}_x)^2$

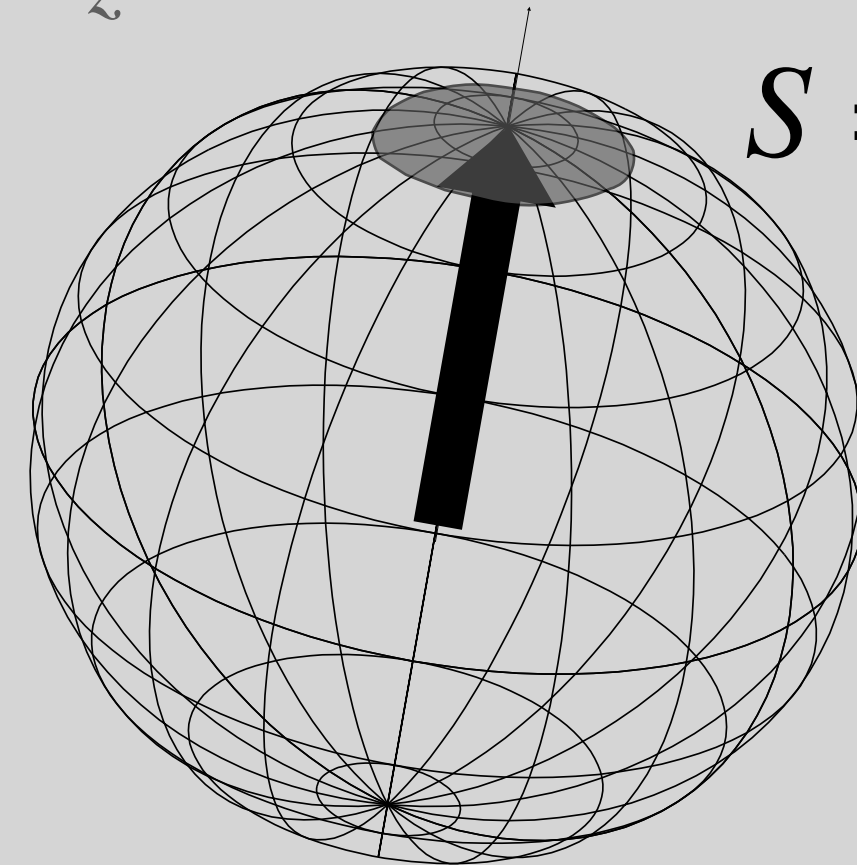
\hat{S}_z



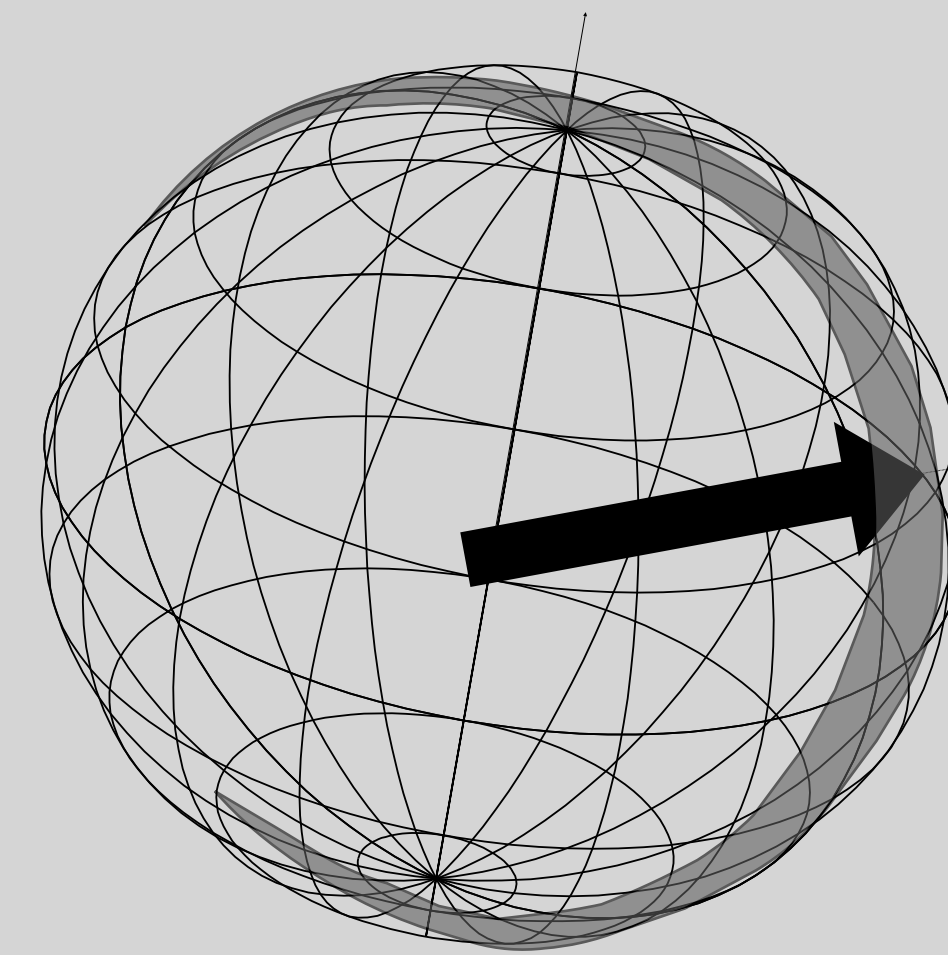
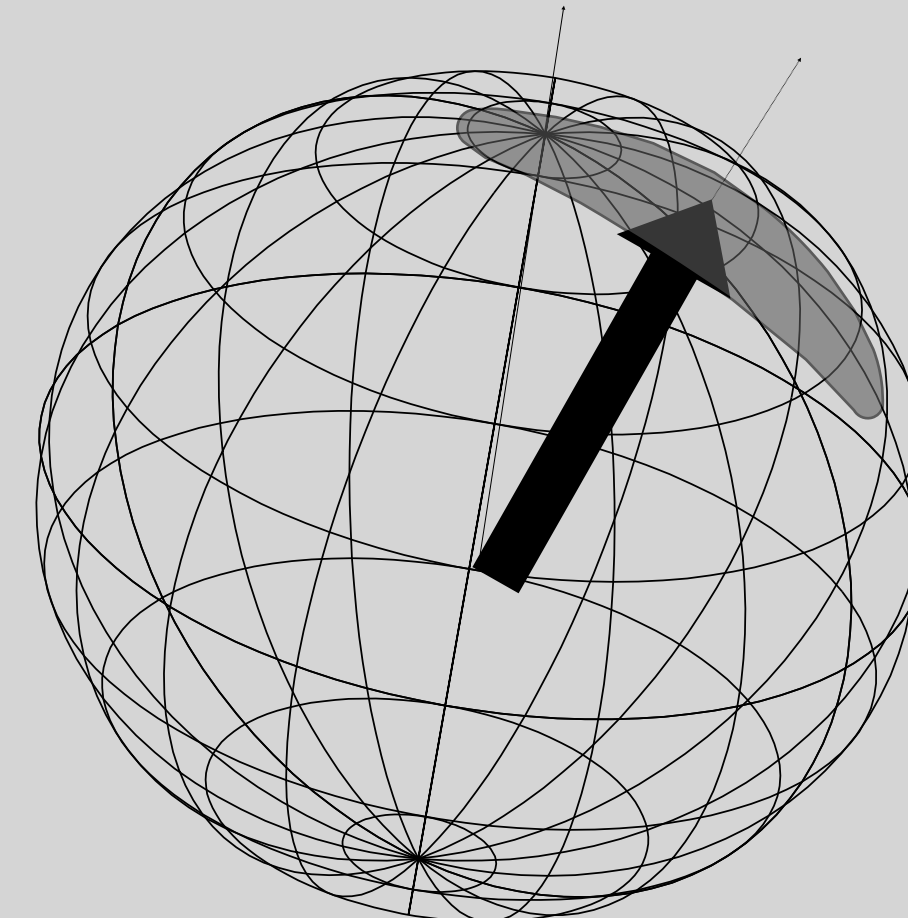
Semiclassical picture of collective spin squeezing

GS and fully polarized states

$$\hbar_{\text{eff}} \sim \hbar/L$$

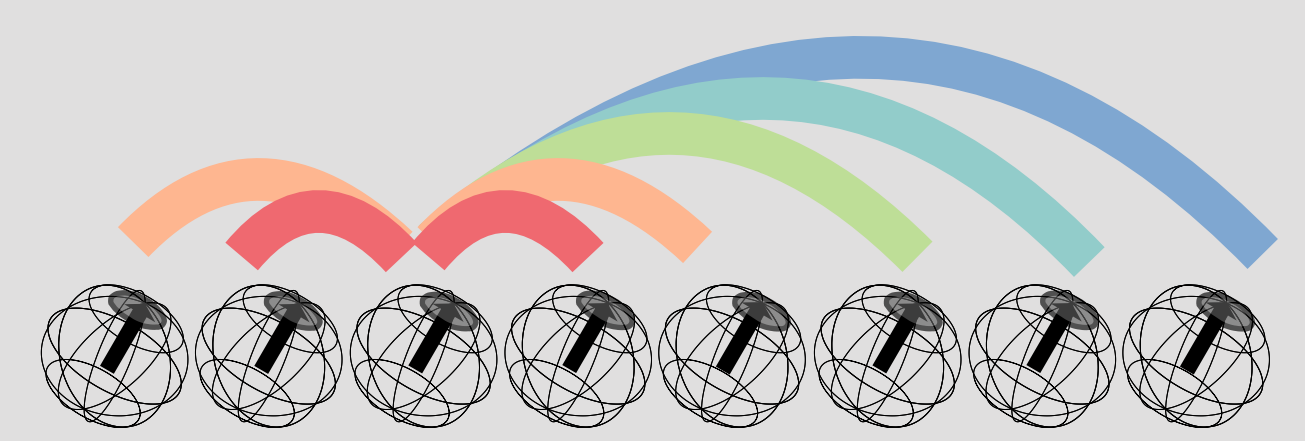


$S = L/2$



\Rightarrow **Thermalization impossible** (spin size conserved) $\vec{S}(t) \propto L$

Dynamics in long-range systems



$\alpha = 0$

$\alpha = d$

α

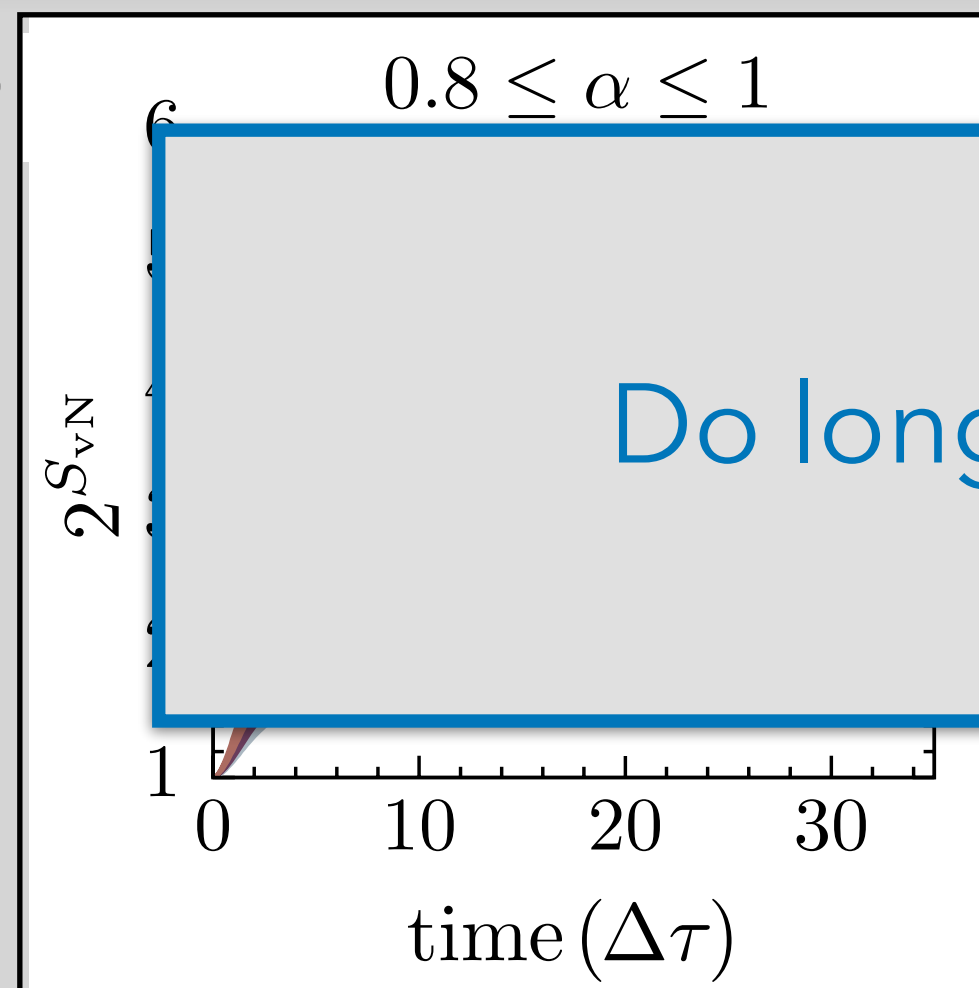
2013

PHYSICAL REVIEW X 3, 031015 (2013)

Entanglement Growth in Quench Dynamics with Variable Range Interactions

J. Schachenmayer,¹ B. P. Lanyon,² C. F. Roos,² and A. J. Daley¹

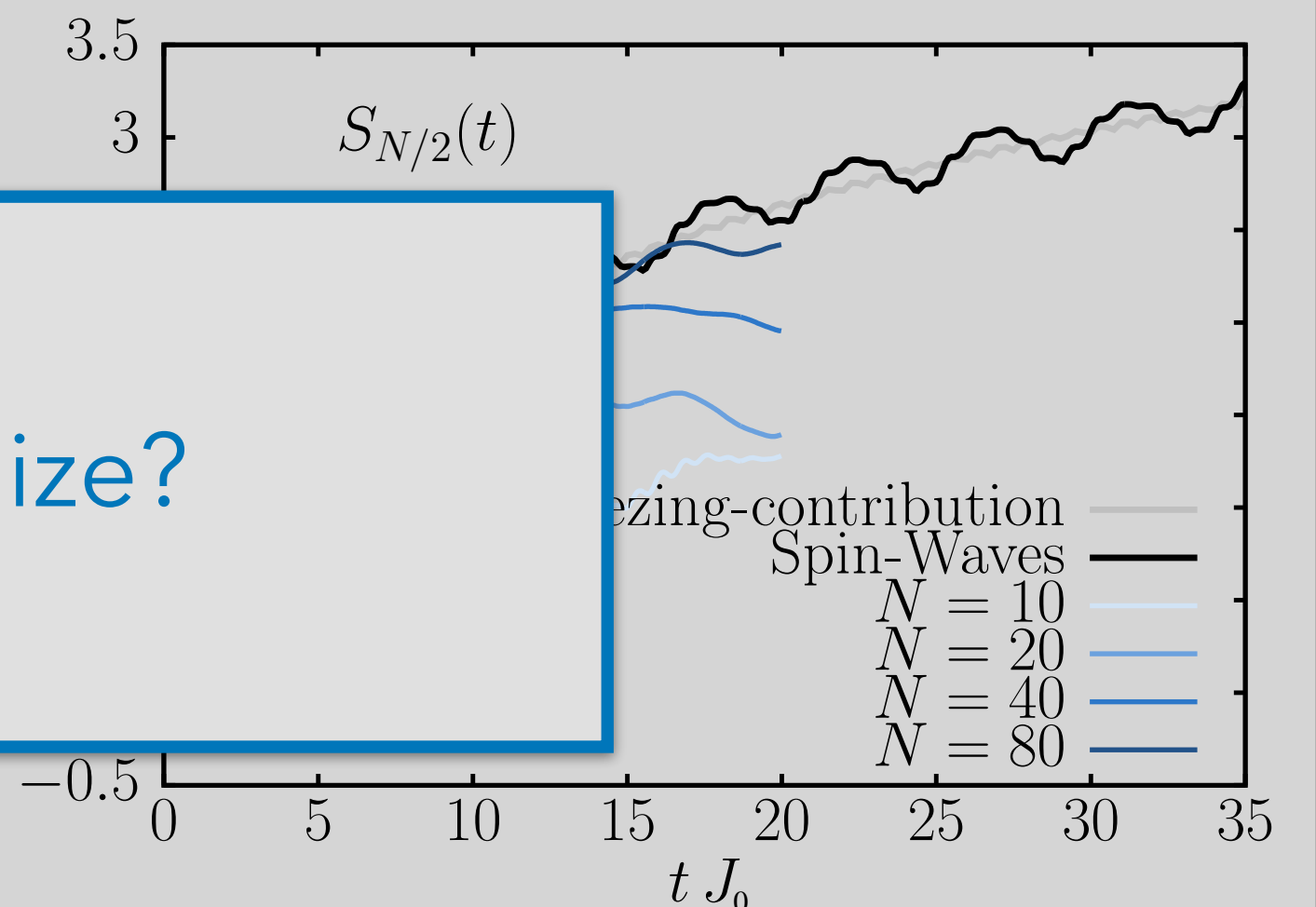
$$S_{\text{vN}} = -\text{Tr}(\rho_{N/2} \ln \rho_{N/2})$$



Do long-range systems ultimately thermalize?

2020

Theory of entanglement entropy dynamics for $0 \leq \alpha \leq d$



different behavior. **Counterintuitively**, quenches above the critical point for these **long-range interactions lead only to a logarithmic increase of bipartite entanglement in time**, so that in this regime, long-range interactions produce a **slower growth of entanglement than short-range interactions**.

NO Semiclassical chaos

\implies Slow (log) growth for $0 < t \ll N^\beta$

Semiclassical chaos

\implies Fast (linear) growth for $0 < t \ll \log N$

In this talk:

Long-range interacting quantum spin systems



Robust quantum many-body scars

1. Integrability of the classical limit $\alpha = 0$ (KAM-like)
2. Sufficiently long-range interactions $0 < \alpha < d$

- I. Numerical analysis
- II. Analytical theory
- III. New theory predictions

Numerics: level spacing statistics

Std metric of quantum chaos/ergodicity

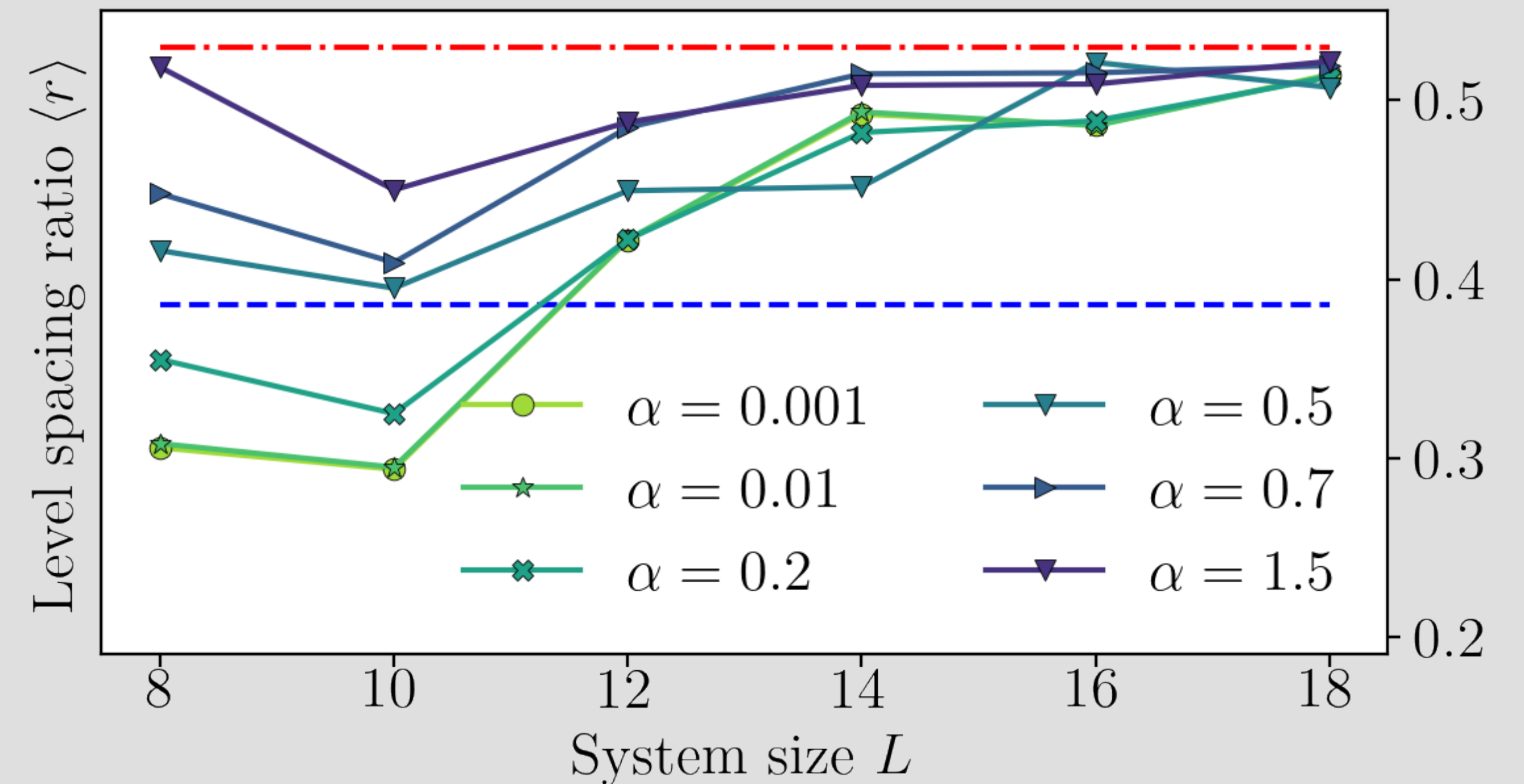
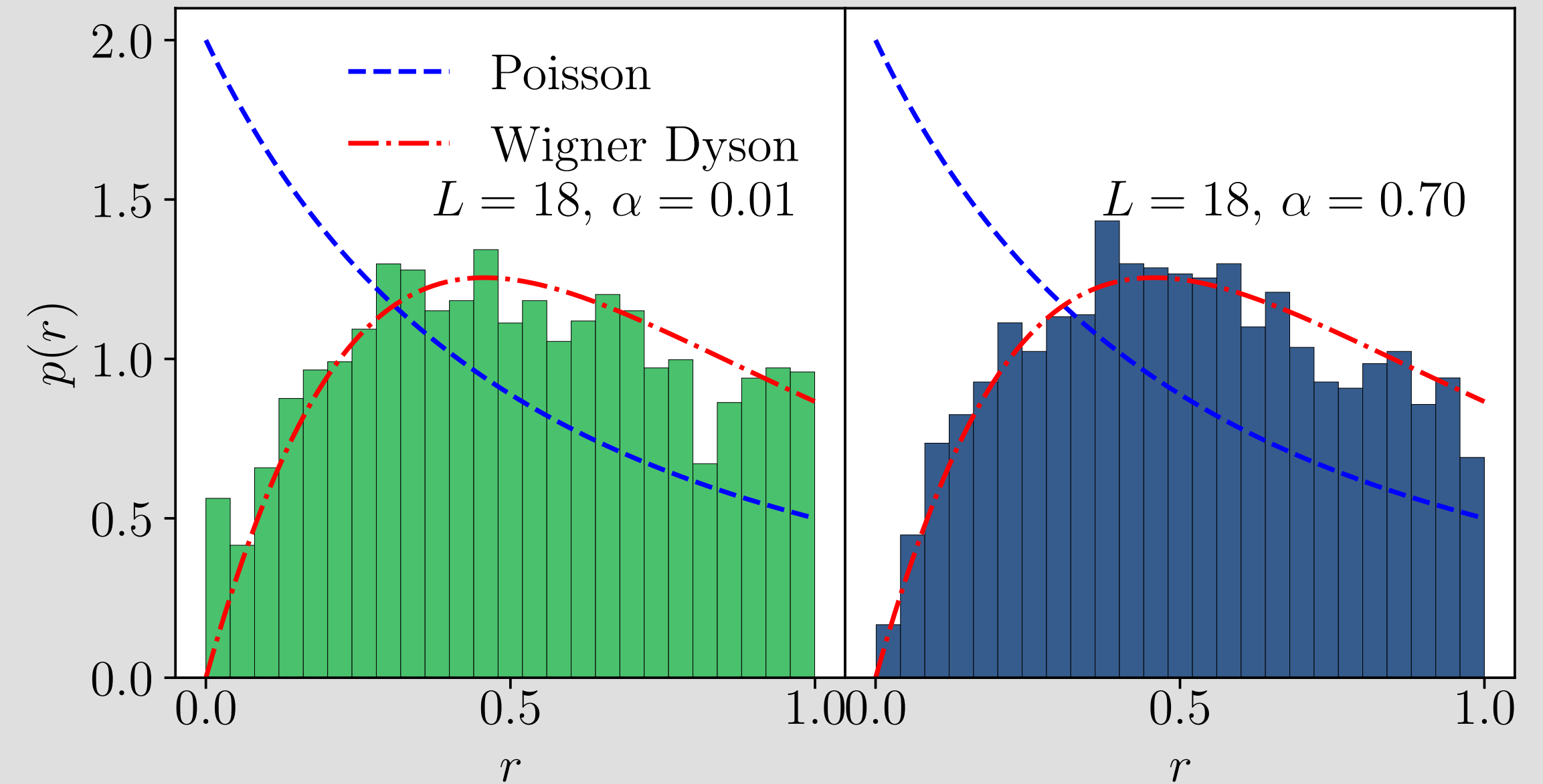
Level spacing ratio

$$r_n = \frac{\min(\Delta E_{n+1}, \Delta E_n)}{\max(\Delta E_{n+1}, \Delta E_n)}$$

- chaos: Wigner Dyson
- Interability: Poisson

Level repulsion for infinitesimal $\alpha > 0 \implies$
quantum ergodicity?

$$h = 2J_0$$



Numerics: individual eigenstates

Collective spin size depletion

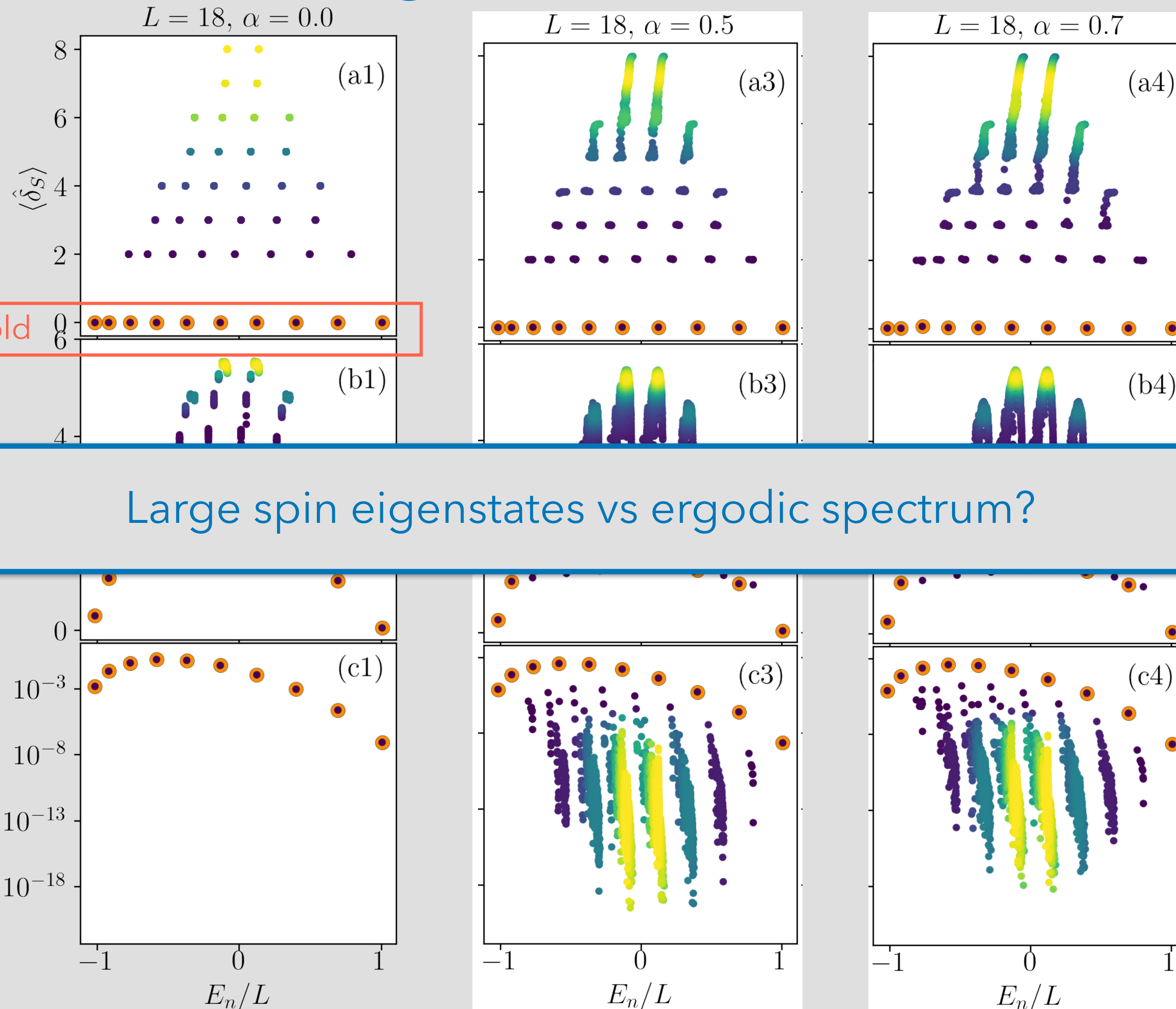
$$\delta_S = L/2 - S$$

$\delta_S = 0$: Dicke manifold

Half-chain entanglement entropy

$$S_{L/2} = -\text{Tr}(\rho_{L/2} \ln \rho_{L/2})$$

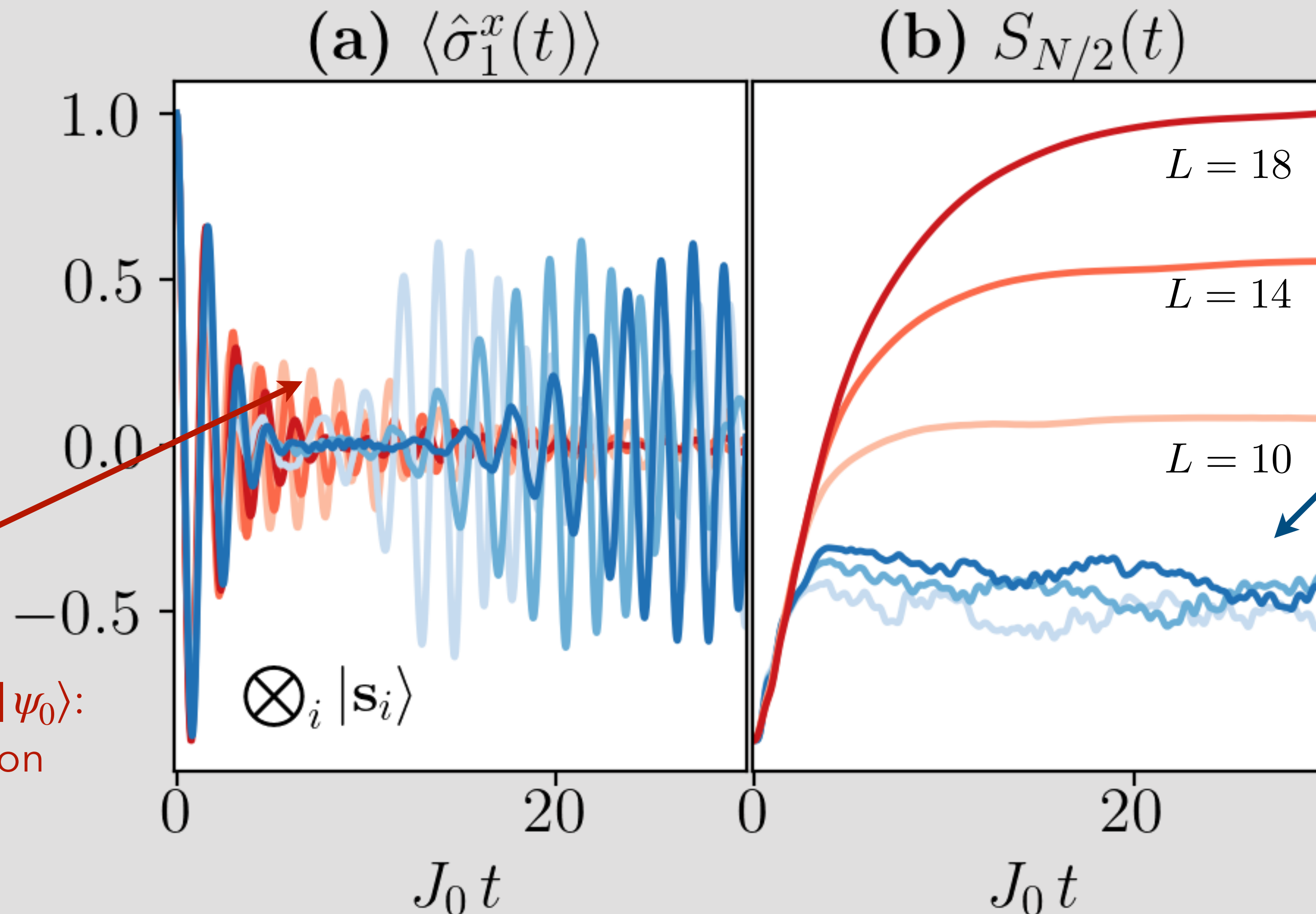
Overlap with spin-coherent state



Numerics: Dynamics at $\alpha = 0.7$

Is it a persistent effect?

Dynamics which strongly depends on the initial state:

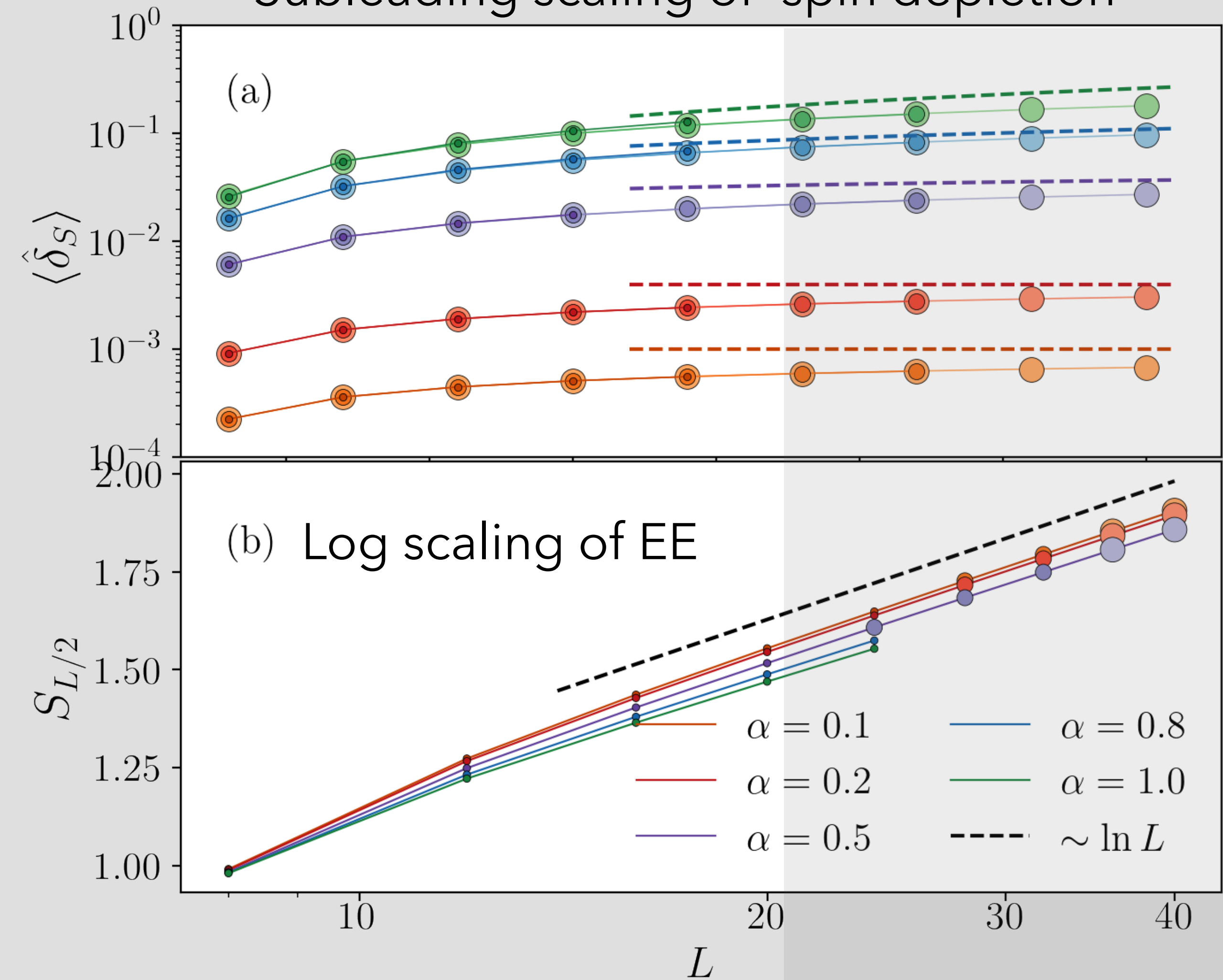


Polarized (coherent)
 $|\psi_0\rangle$: recurrences and
log EE

Random product $|\psi_0\rangle$:
fast thermalization

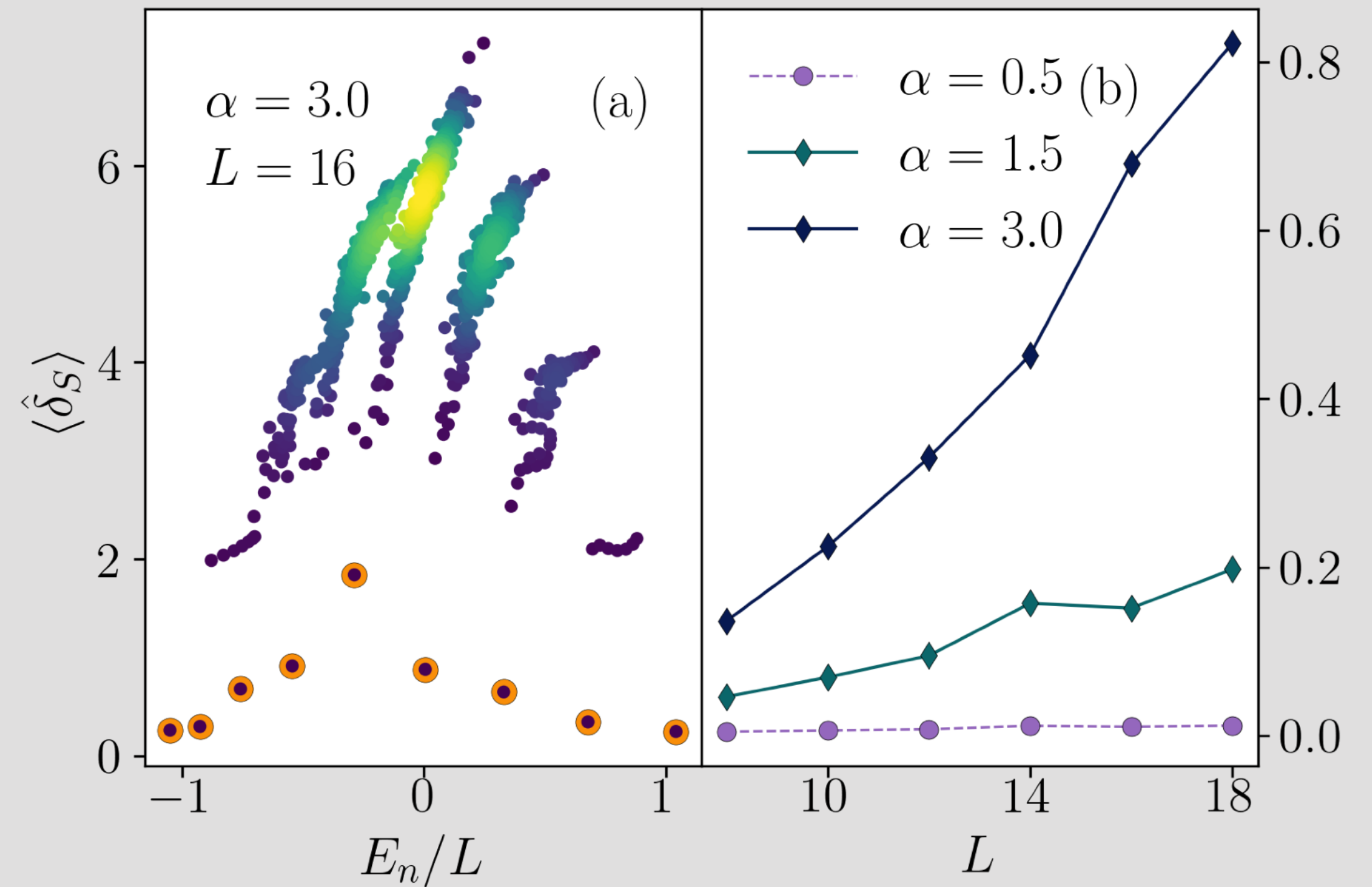
Numerics: scaling with the system size?

Subleading scaling of spin depletion



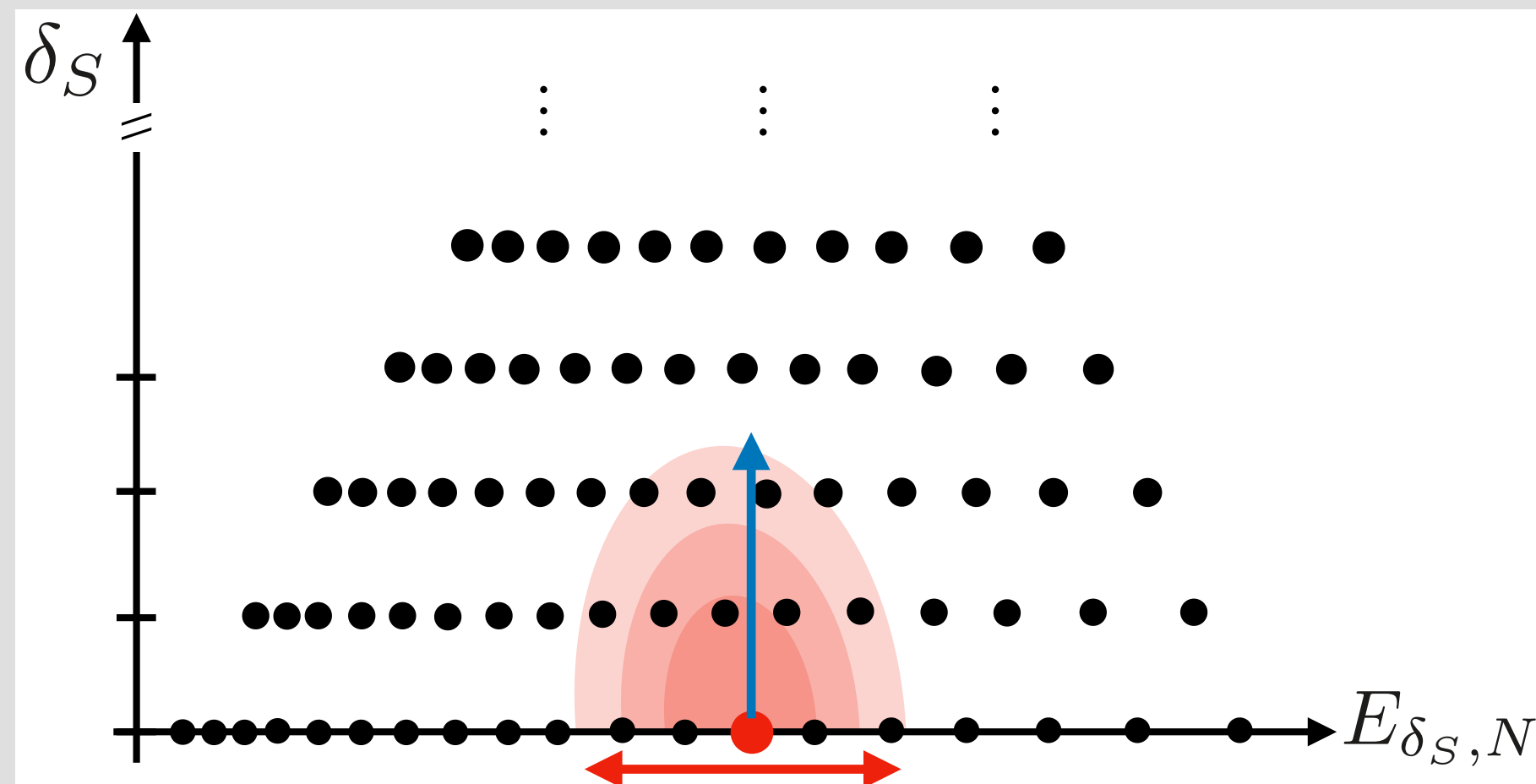
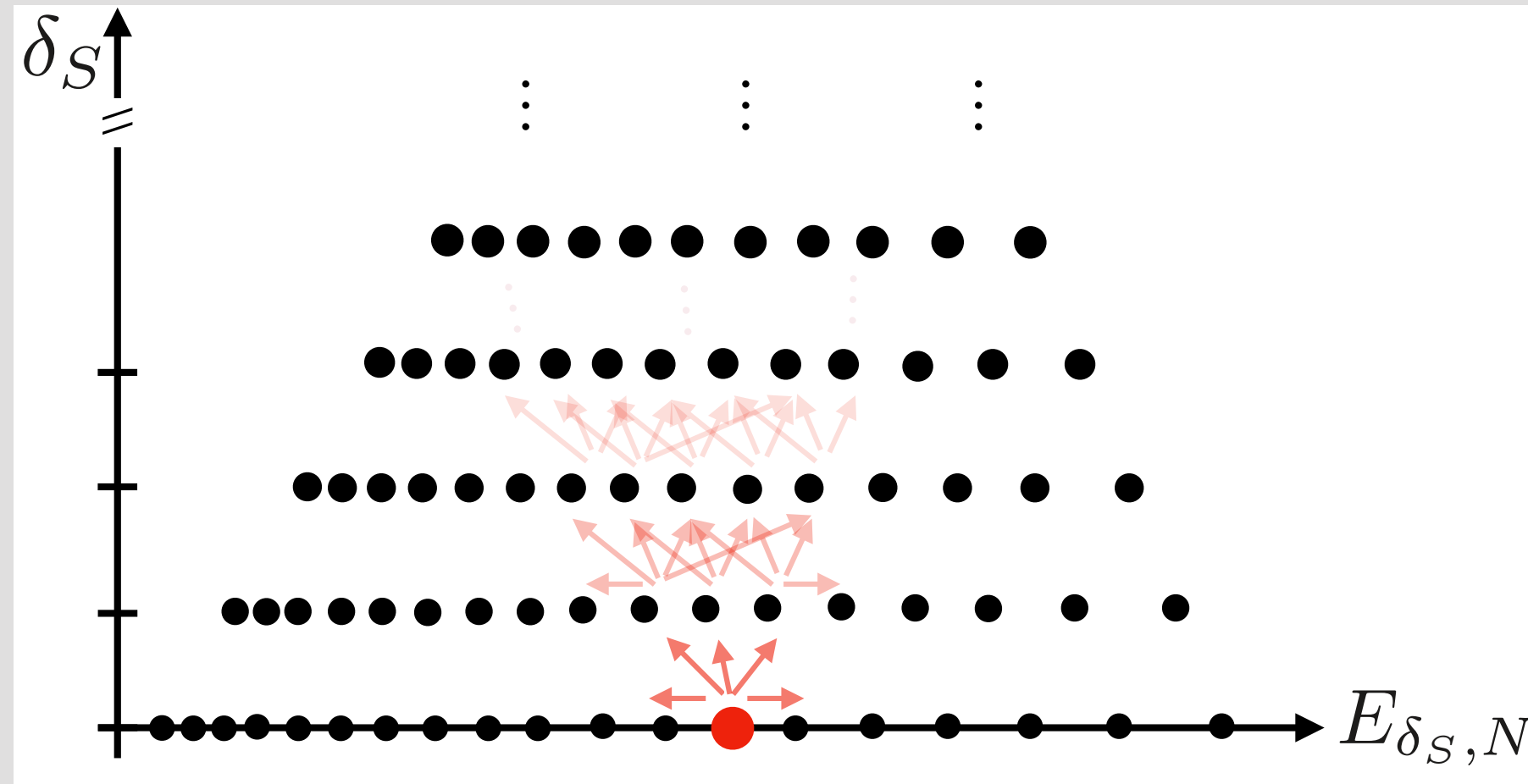
New numerical method

Larger α :



Stability of eigenstates for $0 < \alpha \lesssim 1$?

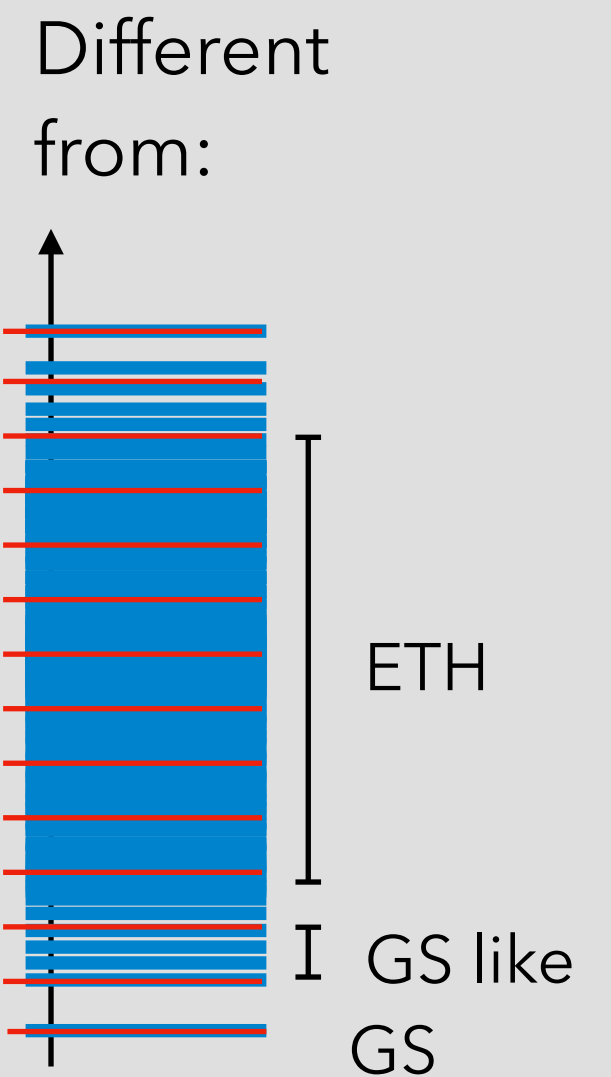
Intuition for eigenstate localization



$$H_{\alpha \neq 0} :$$

off-resonant creation/destruction of magnon pairs

1. small α : $\sum_{k \neq 0} |\tilde{f}_k(\alpha)|^2 \sim L^0$
2. Absence of resonances
(integrable mean-field limit)



eigenstate localization in subspace

$$\delta_S \ll L/2, \quad |\delta_N| \ll L$$

Full calculation

Self consistent many-body scars for $0 < \alpha \lesssim d$

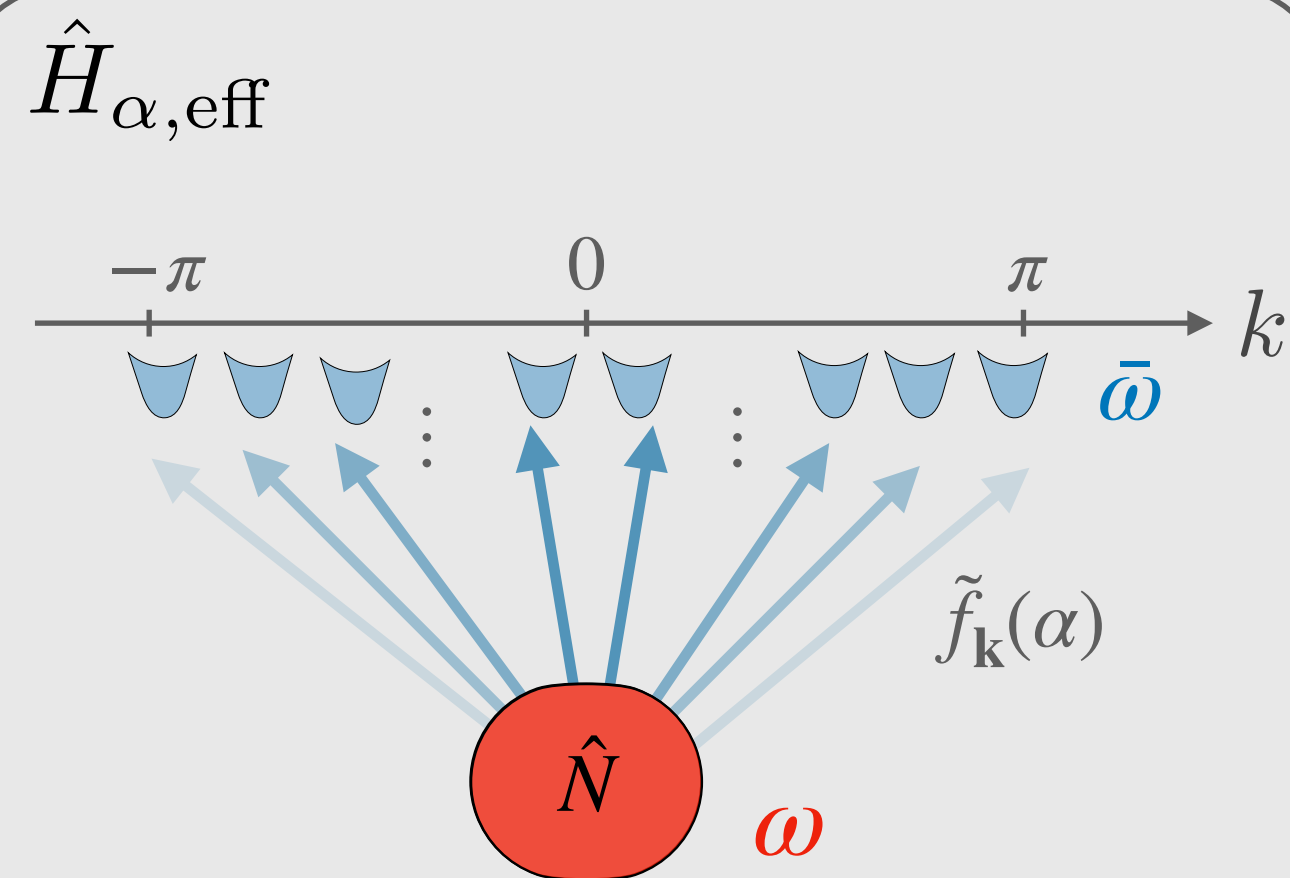
1. Write matrix elements around a collective eigenstate

2. Exactly solvable away from resonances $\omega \neq p\bar{\omega}$ (spectrum + eigenvectors)

3. Compute $\langle \hat{\delta}_S \rangle \sim \sqrt{\langle \hat{\delta}_N^2 \rangle} \sim \sum_{k \neq 0} |f_k(\alpha)|^2 \sim \begin{cases} \text{finite} & \text{for } 0 < \alpha < 1/2, \\ \log L & \text{for } \alpha = 1/2, \\ L^{2\alpha-1} & \text{for } 1/2 < \alpha < 1. \\ c(\alpha) \cdot L & \text{for } \alpha > 1 \end{cases}$

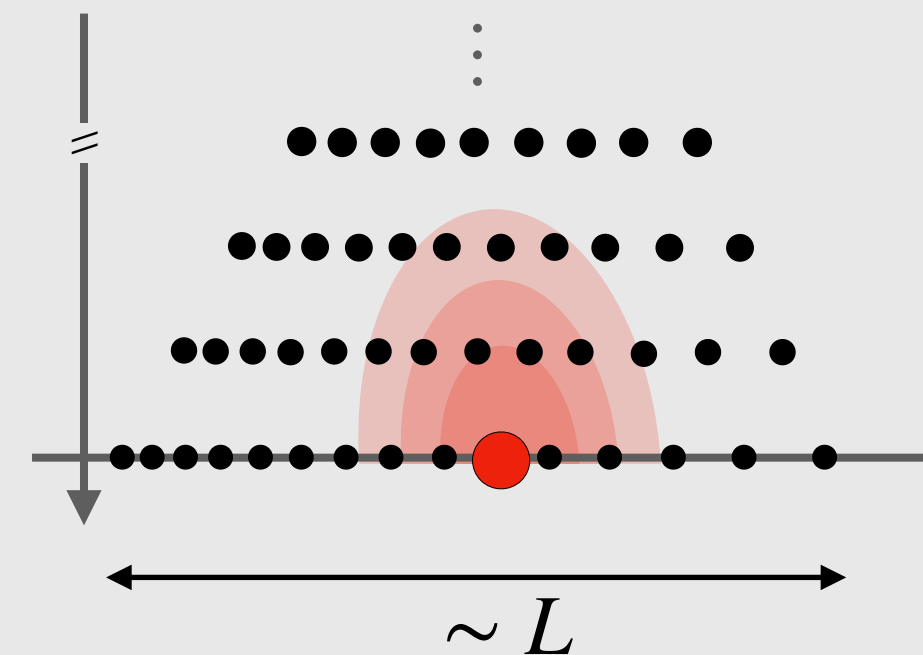
4. Self-consistency

1. $\alpha \neq 0$: Rotor-magnon Hamiltonian interacting quantum impurity model



4. $0 < \alpha < d$: Self-consistent eigenstate localization

2.
$$e^{i\hat{F}} e^{i\hat{S}} \hat{H}_{\alpha,\text{eff}} e^{-i\hat{S}} e^{-i\hat{F}} = \omega \hat{N} + \sum_{k \neq 0} \omega_k(\alpha) \hat{n}_k$$



3.
$$\delta_S \sim L^\beta$$

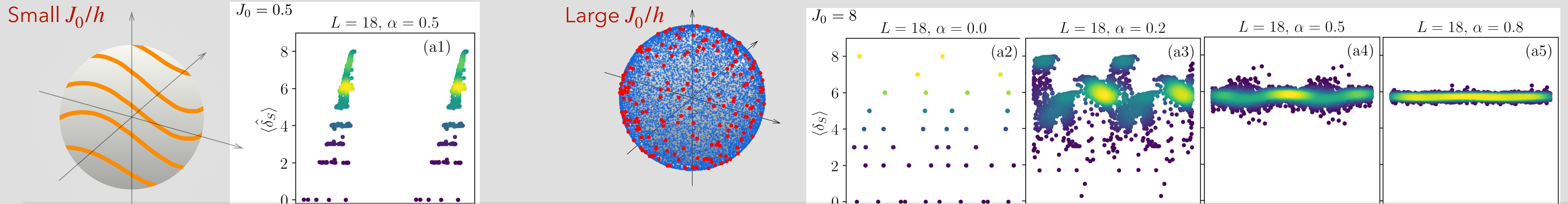
$$\delta_N \sim L^\beta$$

$$\beta = \begin{cases} 0 & \text{for } 0 < \alpha < d/2 \\ 2\alpha - d & \text{for } d/2 < \alpha < d \end{cases}$$

Prediction: Instability from mean-field chaos

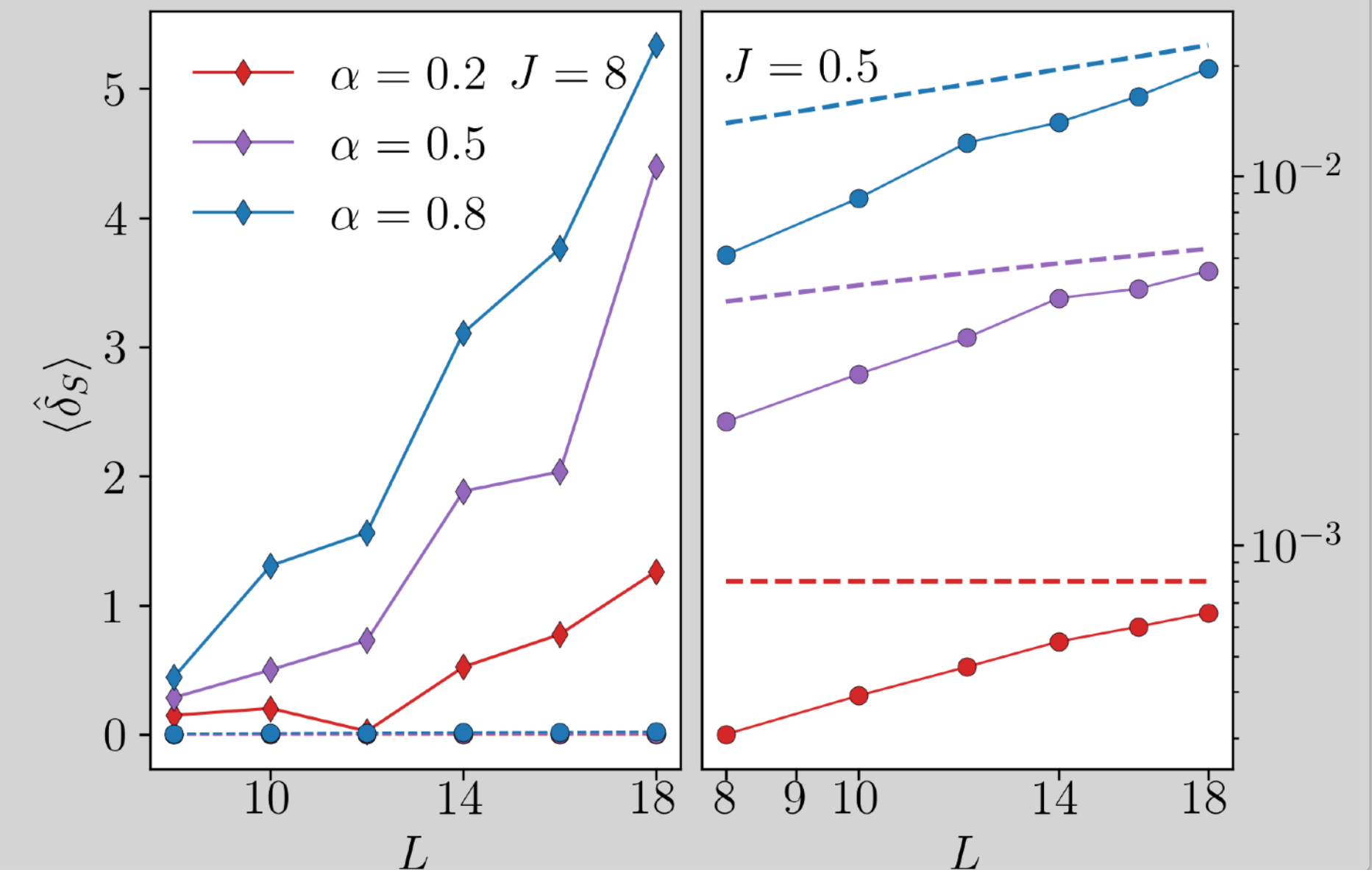
Quantum many-body kicked top:
$$\hat{H}_\alpha(t) = \begin{cases} -\frac{J_0}{\mathcal{N}_{\alpha,L}} \sum_{j=1}^L \sum_{r=1}^{L/2} \frac{\hat{\sigma}_j^x \hat{\sigma}_{j+r}^x}{r^\alpha} & t \in \left[-\frac{T}{4}, \frac{T}{4}\right) \bmod T \\ -h \sum_{j=1}^L \hat{\sigma}_j^z & t \in \left[\frac{T}{4}, \frac{3}{4}T\right) \bmod T \end{cases}$$

$\alpha = 0$: semiclassical Haake, ...
integrability-chaos crossover



Scaling with system size

- Absence of heating for $0 < \alpha < d$ in classical KAM regime
- Semiclassical chaos destroys quantum many-body scarring

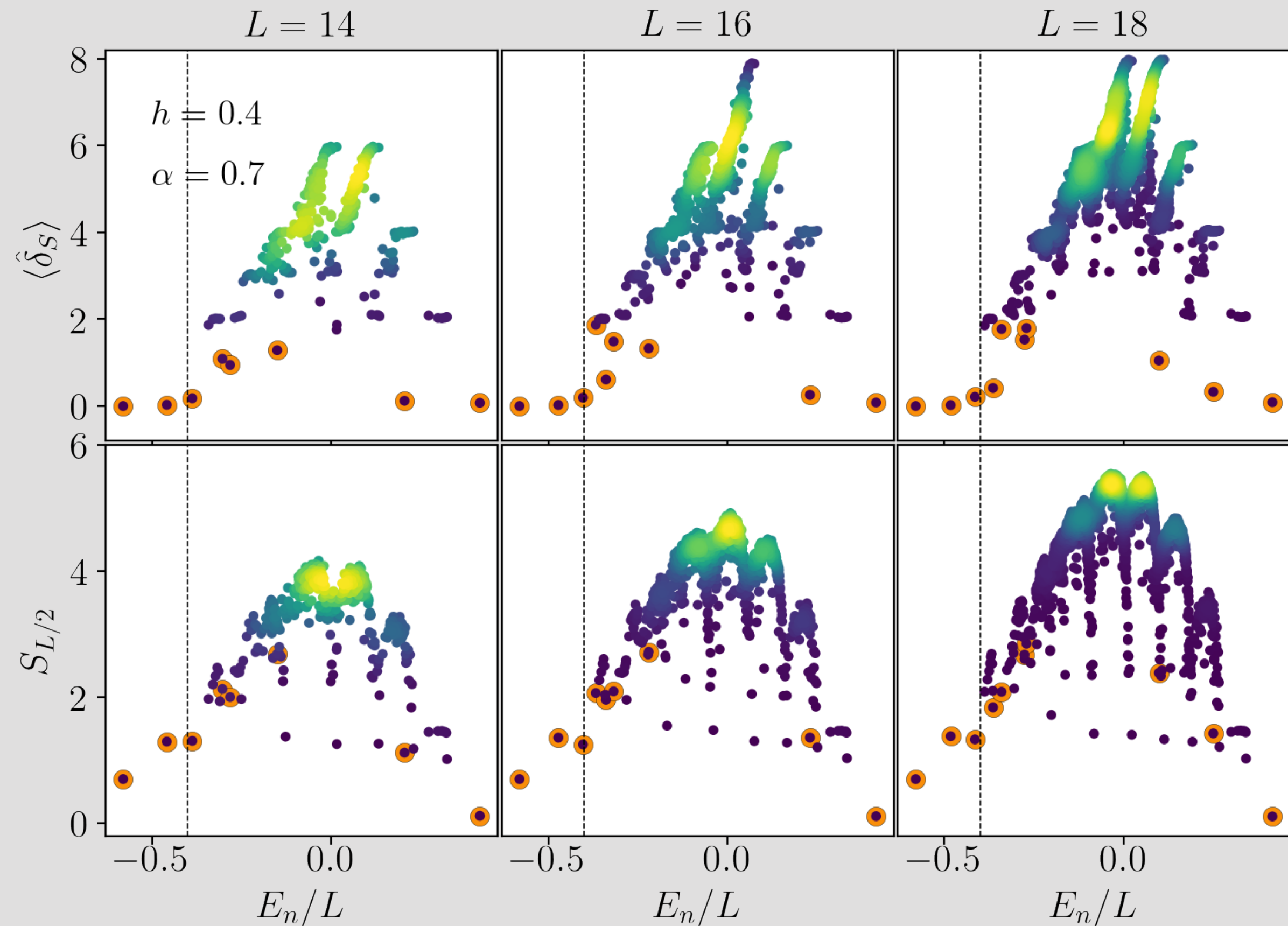
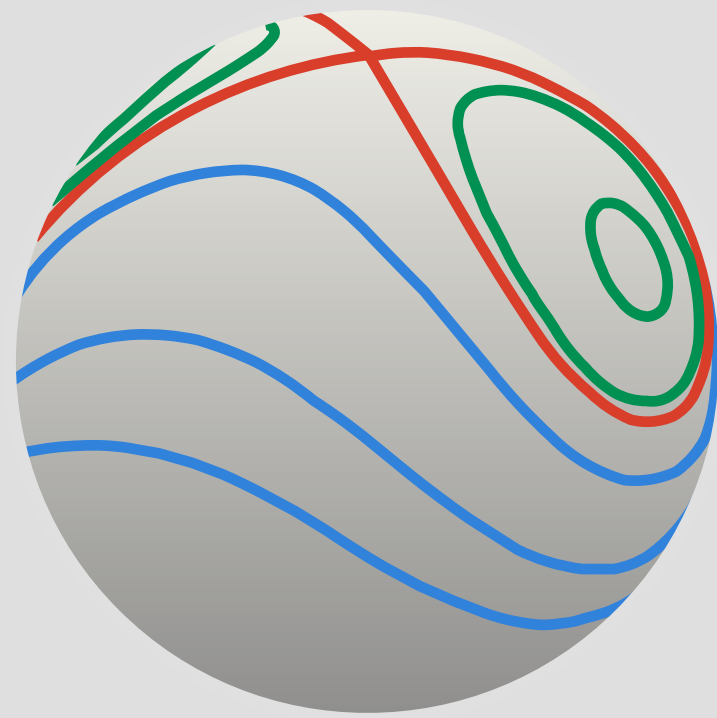


Prediction: instability at the Excited State Phase Transition

Mean-field criticality at finite energy densities

Rev: Cejnar, Strànsky Macek Kloc, J.phys.A 2021

For $\alpha = 0$ associated to separatrix trajectories

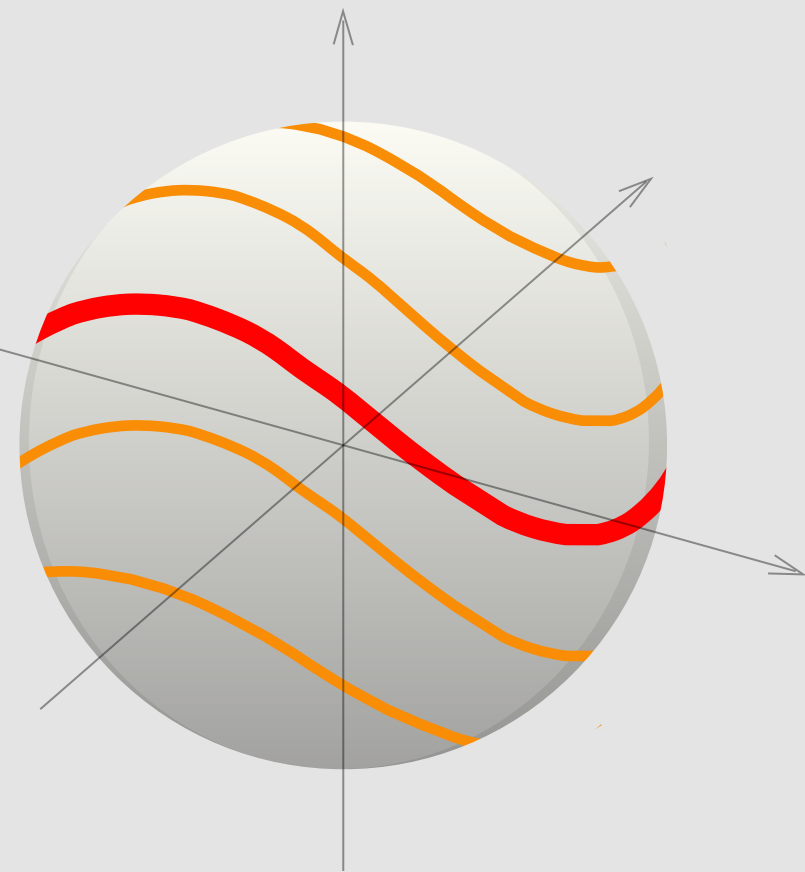


In summary...

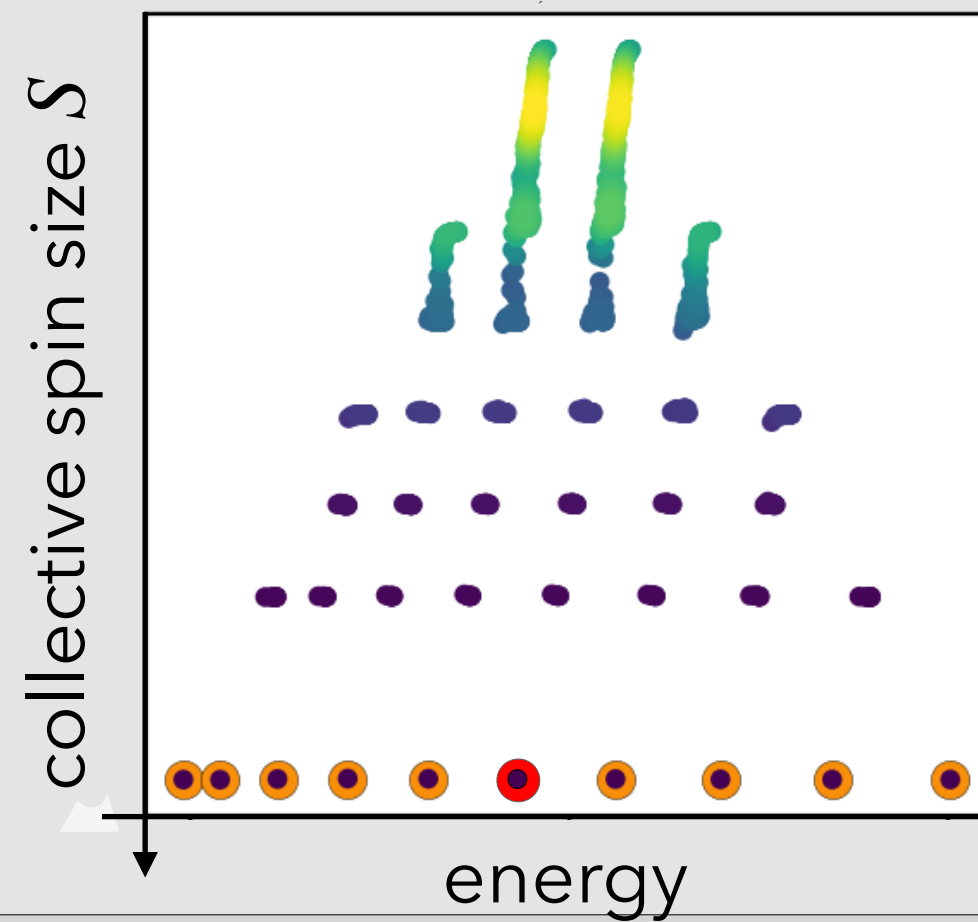
1) Integrability of the underlying classical limit

Classical limit $\alpha = 0$

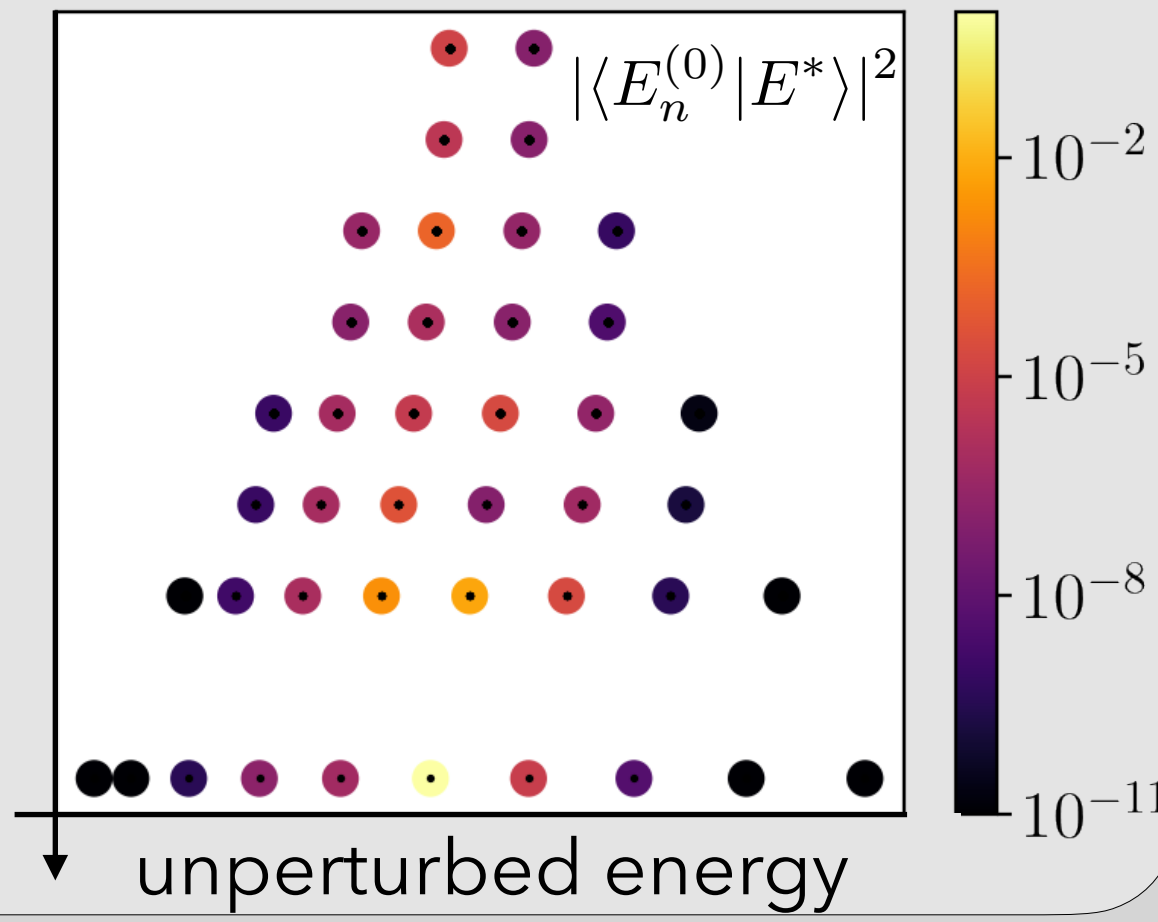
Ising model



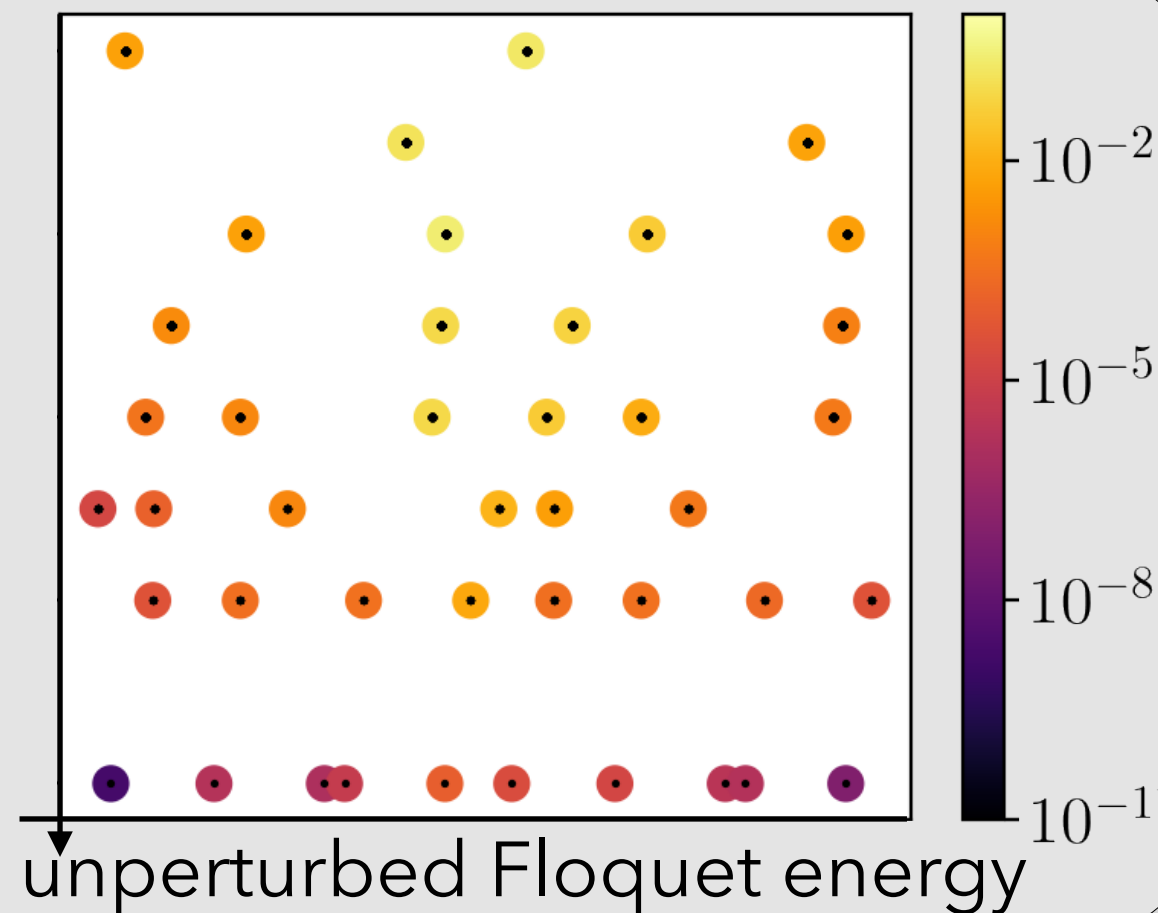
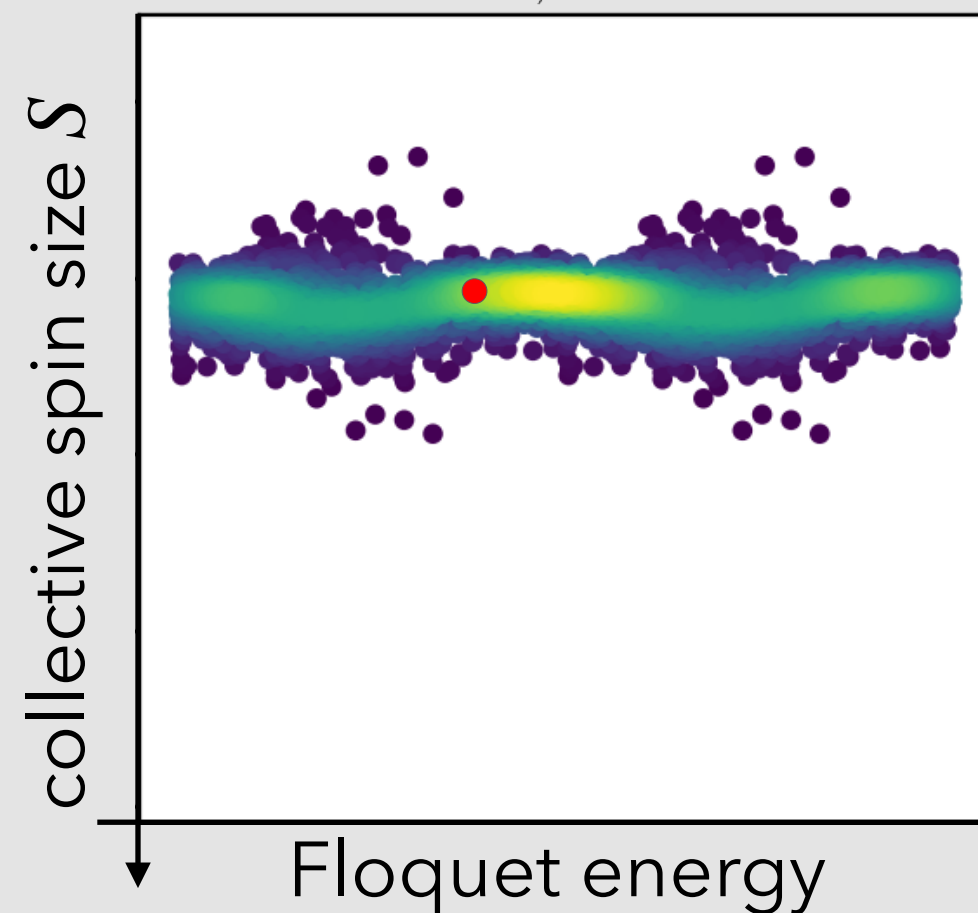
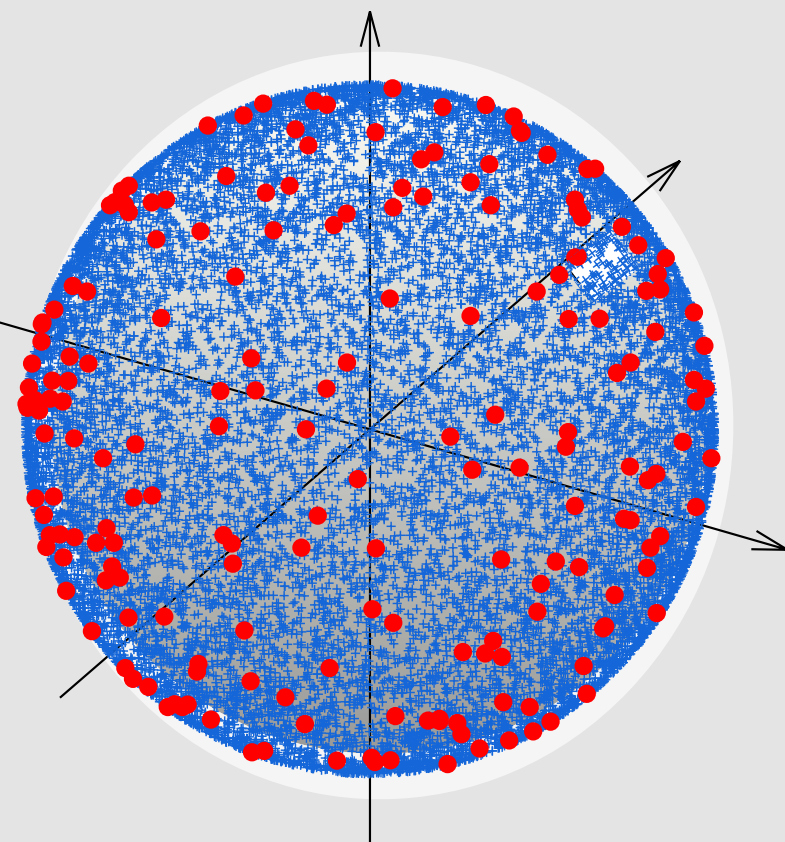
Exact spectrum
 $0 < \alpha < d$



Eigenstate (de)localization
 $0 < \alpha < d$



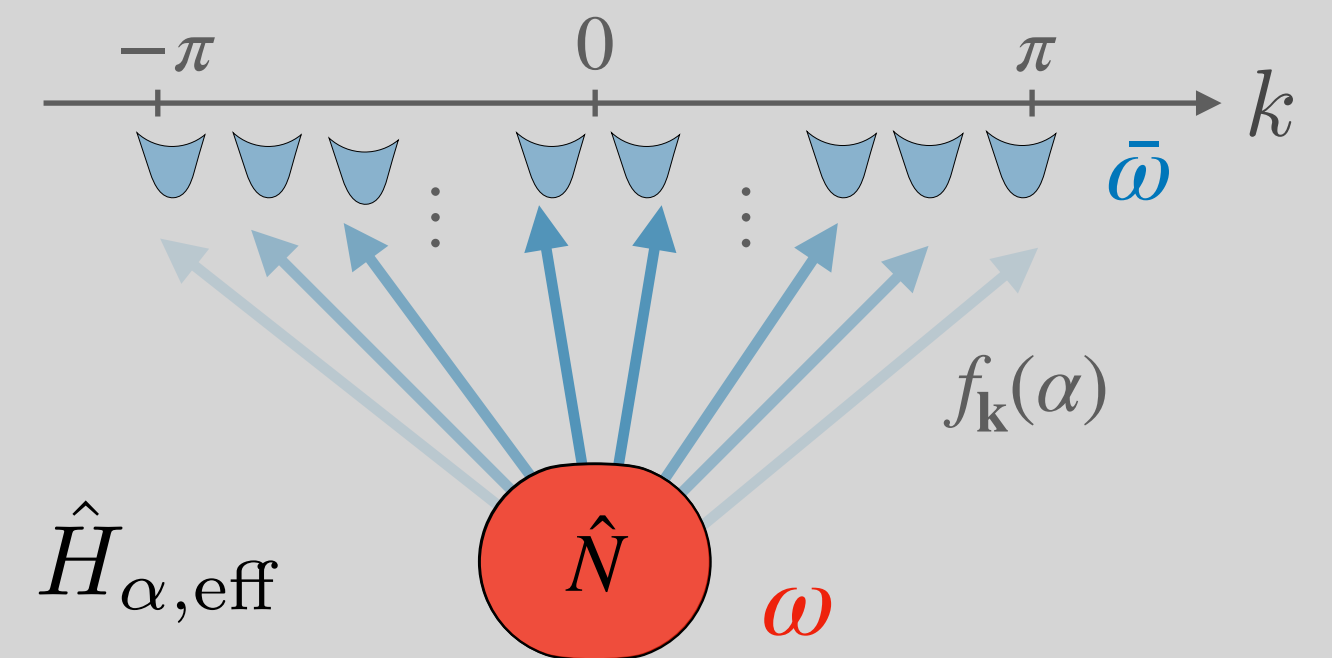
Kicked Ising model



2) Long-range interactions

$0 < \alpha < d$

Shown by a self-consistent magnon/rotor projection

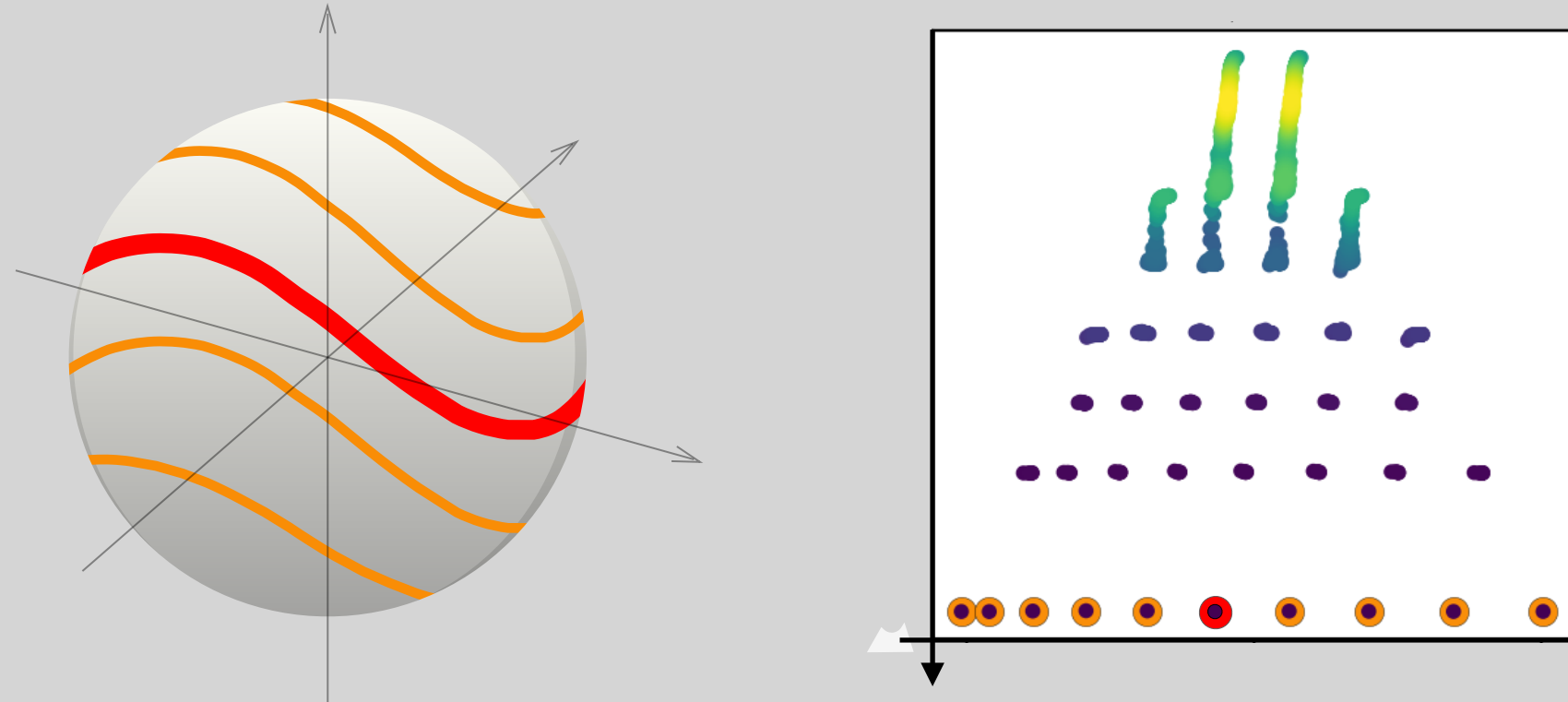


these systems effectively interpolates between few body and many-body physics

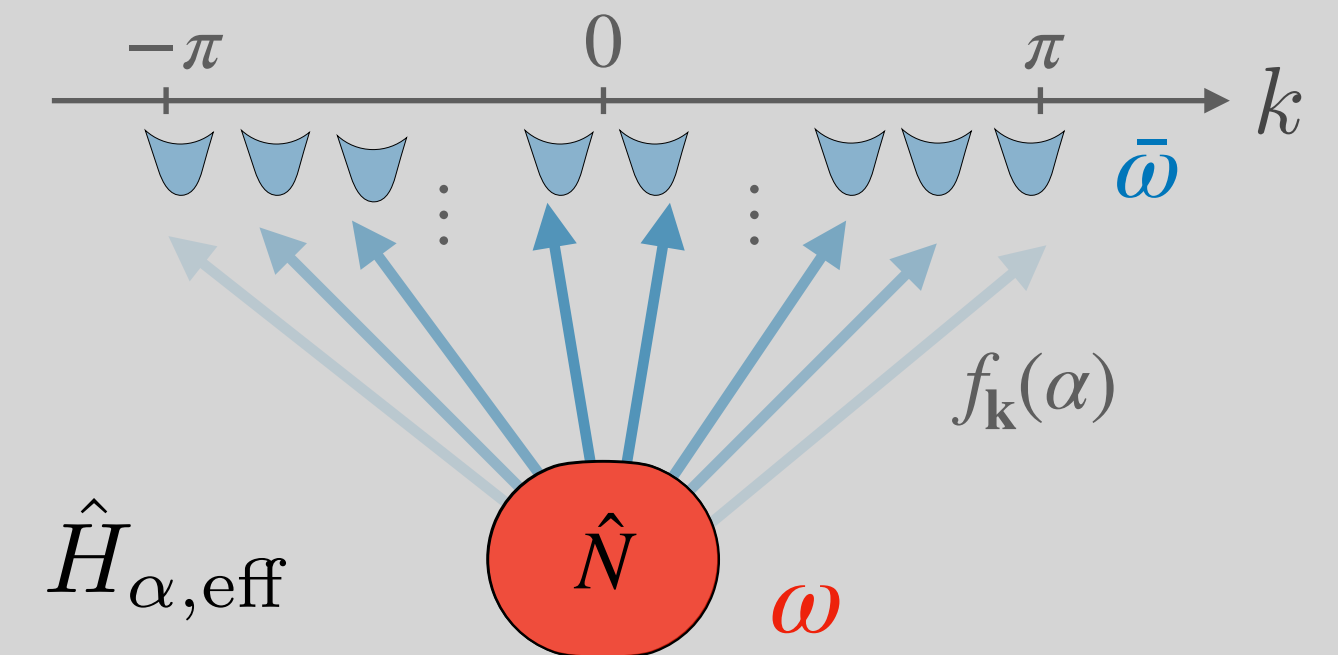
... many open perspectives

Lerose, Parolini, Fazio, Abanin and **SP**, Arxiv 2309.12504

1) Integrability of the underlying classical limit



2) Long-range interactions



- Quantum KAM? Instability to many-body perturbations?
- suppressed heating for $0 < \alpha < d$: time-crystal?
- *Robust* Dicke-like states with $\alpha \neq 0 \rightarrow$ Useful metrological applications?
- generalize mechanism to be applied to other systems (PXP)?
- novel analytical method \rightarrow new numerical approach?

Thanks!